

## INFRARED SPECTRA OF CRYSTALLINE PHASE ICES CONDENSED ON SILICATE SMOKES AT $T < 20$ K

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### ABSTRACT

Infrared spectra of  $\text{H}_2\text{O}$ ,  $\text{CH}_3\text{OH}$ , and  $\text{NH}_3$  condensed at  $T < 20$  K on amorphous silicate smokes reveal that predominantly crystalline phase ice forms directly on deposit. Spectra of these molecules condensed on aluminum substrates at  $T < 20$  K indicate that amorphous phase ice forms. On aluminum, crystalline phase  $\text{H}_2\text{O}$  and  $\text{CH}_3\text{OH}$  are formed by annealing amorphous deposits to 155 K and 130 K, respectively (or by direct deposit at these temperatures); crystalline  $\text{NH}_3$  is formed by direct deposit at 88 K. Silicate smokes are deposited onto aluminum substrates by evaporation of  $\text{SiO}$  solid or by combustion of  $\text{SiH}_4$  with  $\text{O}_2$  in flowing  $\text{H}_2$  followed by vapor phase nucleation and growth. Silicate smokes which are oxygen-deficient may contain active surface sites which facilitate the amorphous-to-crystalline phase transition during condensation. Detailed experiments to understand the mechanism are currently in progress. The assumption that amorphous phase ice forms routinely on grains at  $T < 80$  K is often used in models describing the volatile content of comets or in interpretations of interstellar cloud temperatures. This assumption needs to be reexamined in view of these results.

*Subject headings:* comets: general — dust, extinction — infrared: general — ISM: molecules —  
 line: identification — techniques: spectroscopic

Comets are believed to have formed from grains and amorphous ices located in the outer part of the presolar nebula  $4.6 \times 10^9$  yr ago or in interstellar clouds (for a review see Donn 1991). Based on the interstellar dust model of Greenberg (1982), individual grains are likely to contain a silicate core, a mantle of refractory material, and an outer layer of  $\text{H}_2\text{O}$  dominated volatile ices. The signature of these ice/dust materials is often detected in comets and in interstellar objects. For example, the  $9.7 \mu\text{m}$  feature attributed to silicates is observed in a variety of interstellar sources (e.g., Nuth & Hecht 1990) and has been observed in several comets (e.g., Bregman et al. 1987). The mid-infrared features of  $\text{H}_2\text{O}$  have been detected in more than 50 sources. Frequently the  $3 \mu\text{m}$  region of interstellar sources is fitted with laboratory spectra of amorphous phase water ice or water-rich icy mixtures (e.g., Smith, Sellgren, & Tokunaga 1989). However, fitting the band shape sometimes requires both the amorphous and crystalline phase of  $\text{H}_2\text{O}$  even when a mixture of other components is used.

The formation of amorphous phase deposits of  $\text{H}_2\text{O}$  and many other volatile species at  $T < 50$  K is documented in numerous laboratory studies using X-ray techniques (e.g., Narten, Venkatesh, & Rice 1976). Amorphous ice, once formed, is unstable against temperature cycling. It converts irreversibly to a crystalline phase on a timescale which decreases with increasing temperature. The conversion is completed on the order of minutes, for example, for  $\text{H}_2\text{O}$  at 155 K and for  $\text{CH}_3\text{OH}$  at 130 K.

Laboratory measurements of various ice/dust properties are crucial for our understanding of these complex comet mixtures (Boice, Naegeli, & Hubner 1989) and interstellar analogs. Zhao

(1990) compared a residue formed from a photolyzed icy mixture on a pressed silicate powder (ground from bulk materials to a grain size near 100s of nanometers) with a residue from a photolyzed icy mixture deposited on a metal substrate. He detected some changes in the subfeatures in the  $3.4 \mu\text{m}$  absorption region and in the rate of formation of  $-\text{NH}_2$  groups. The idea that very reactive surface sites can form in high surface area  $\text{SiO}_2$  powders has been studied using gas phase reactions of interest to the microelectronics industry; for example, the dissociation of  $\text{NH}_3$  and  $\text{H}_2\text{O}$  during chemisorption has been documented (Morrow & Cody 1976). Similar studies of ice/silicate mixtures have not been done.

In this *Letter* we report on laboratory studies showing that predominantly crystalline phase ices form at  $T < 20$  K on laboratory-produced silicate smokes. Infrared spectra of  $\text{H}_2\text{O}$ ,  $\text{CH}_3\text{OH}$ , and  $\text{NH}_3$  are discussed;  $\text{SO}_2$  and  $\text{H}_2\text{CO}$  have also been studied with similar results.  $\text{H}_2\text{O}$  and  $\text{CH}_3\text{OH}$  are abundant interstellar molecules and play important roles in the physical chemistry of comets;  $\text{NH}_3$  may play a less important role.

Silicate smokes were made either by evaporation of  $\text{SiO}$  solid (Nuth & Donn 1982) or by combustion of  $\text{SiH}_4$  with  $\text{O}_2$  followed by vapor phase nucleation and growth in an  $\text{H}_2$  atmosphere (Nelson et al. 1989). Smokes were deposited onto polished aluminum substrates (area =  $5 \text{ cm}^2$ ); thicknesses were typically 0.1–0.5 mm. Smoke colors were shades of beige-gold indicating an oxygen-deficient silicate; completely oxidized silicate ( $\text{SiO}_2$ ) is white. Smoke particles are typically 5–10 nm in diameter (Rietmeijer & Nuth 1991) and are generally amorphous in composition and in morphology (Rietmeijer, Nuth, & MacKinnon 1986); individual smoke grains consolidate during deposition resulting in a smoke porosity we estimate to be 97%. Bonding to the aluminum was not a problem, although the smokes were fragile and could easily be compressed or rubbed off. All samples were stored in a vacuum desiccator until used for an experiment. Preliminary measure-

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ments yield a typical smoke density of  $0.08 \text{ g cm}^{-3}$ ; the surface area of these smokes, determined from  $\text{N}_2$  adsorption isotherms at 77 K, is  $\sim 125 \text{ m}^2 \text{ g}^{-1}$ .

The absorbance spectrum of a typical oxygen-deficient silicate smoke is shown in Figure 1 from  $1600 \text{ cm}^{-1}$  ( $6.2 \mu\text{m}$ ) to  $400 \text{ cm}^{-1}$  ( $25 \mu\text{m}$ ), a region which contains the strongest absorption bands (absorptions in other regions are very weak or nonexistent). Characteristic absorptions of both amorphous  $\text{SiO}_2$  and  $\text{Si}_2\text{O}_3$  are present (Nuth & Donn 1982). The so-called  $10 \mu\text{m}$  silicate feature peaks at  $1087 \text{ cm}^{-1}$  ( $9.2 \mu\text{m}$ ); less intense bands occur at  $880 \text{ cm}^{-1}$  ( $11.4 \mu\text{m}$ ) and  $457 \text{ cm}^{-1}$  ( $21.9 \mu\text{m}$ ). In contrast to this spectral profile, highly oxidized silicate smokes have strong absorptions at  $808 \text{ cm}^{-1}$  ( $12.4 \mu\text{m}$ ) and near  $460 \text{ cm}^{-1}$  ( $21.6 \mu\text{m}$ ).

Each ice/silicate composite was formed in vacuum by slow condensation of gas onto a silicate smoke which was cooled to  $T < 20 \text{ K}$ . The rate of deposition was the order of  $10^{20}$  molecules  $\text{hr}^{-1}$  which is equivalent to  $\sim 10$ s of microns  $\text{hr}^{-1}$  on an aluminum substrate, but is estimated to be only  $\sim 10$ s of monolayers  $\text{hr}^{-1}$  on each grain of a  $0.5 \text{ mm}$  thick smoke. Since ice/silicate spectra were compared to ice spectra on aluminum substrates, typical minimum absorbances for a  $\text{CH}_3\text{OH/silicate}$  in the  $3 \mu\text{m}$  region were 0.2. These deposits routinely formed crystalline ices. Ultrathin  $\text{CH}_3\text{OH}$  deposits (absorbance  $< 0.2$ ) on silicates were amorphous but evolved, with the addition of more  $\text{CH}_3\text{OH}$  at  $T < 20 \text{ K}$ , into the crystalline signatures we originally observed. We estimate the transition to the crystalline ices occurs after the equivalent of  $\approx 10$  monolayers of condensate have formed. Spectra of the  $45 \mu\text{m}$  band of  $\text{H}_2\text{O}$ , the  $3 \mu\text{m}$  region of  $\text{CH}_3\text{OH}$ , and the  $10 \mu\text{m}$  region of  $\text{NH}_3$  were recorded using a Mattson (Polaris) FTIR. These spectral regions were chosen because of the large differences observed between the amorphous and crystalline spectral profiles. Spectra consist of at least 60 coadded scans and have a resolution of  $4 \text{ cm}^{-1}$ . In our configuration, the infrared beam is reflected from the aluminum substrate producing an

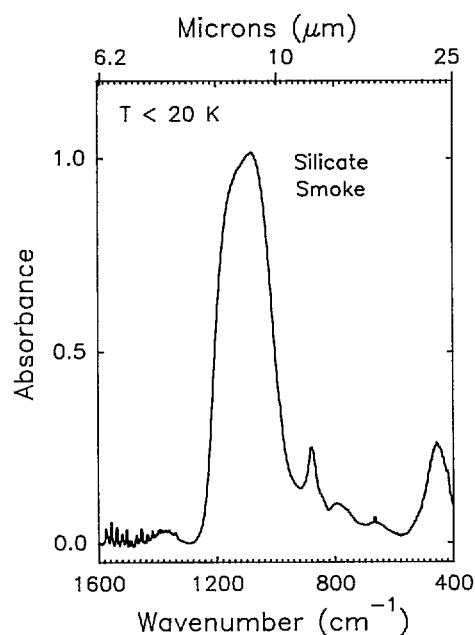


FIG. 1.—Infrared spectrum, typical of laboratory-produced silicate smokes on aluminum substrate, is a mixture of  $\text{SiO}_2$  and  $\text{Si}_2\text{O}_3$ .

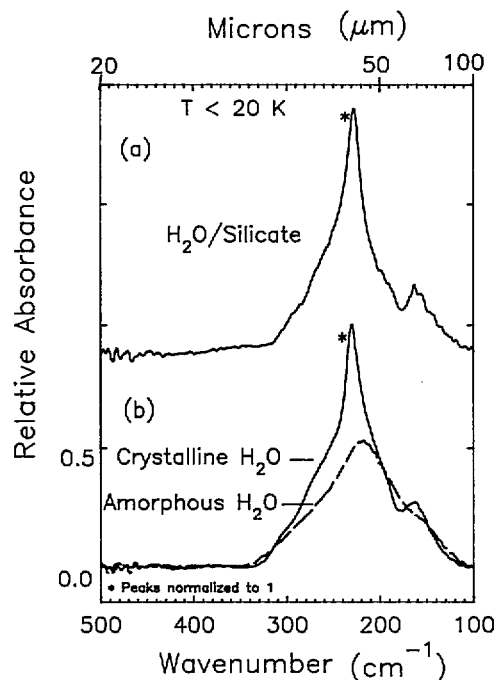


FIG. 2.—Far-infrared spectra of  $\text{H}_2\text{O}$ .  $\text{H}_2\text{O}$  deposited directly onto a silicate smoke (a) is compared with spectra of  $\text{H}_2\text{O}$  in the amorphous and crystalline phase on an aluminum substrate (b).

absorbance spectrum after two passes through the ice/silicate film. Each ice/silicate spectrum was ratioed with the spectrum of the pure silicate to remove the silicate features from the spectrum (for more details, see Moore & Hudson 1992).

Figure 2 summarizes the far-infrared spectral results for  $\text{H}_2\text{O}$  deposits at  $T < 20 \text{ K}$ .  $\text{H}_2\text{O}$  condensed onto silicate smoke results in the  $\text{H}_2\text{O/silicate}$  spectrum shown in Figure 2a. This spectrum is compared with Figure 2b showing the single broad symmetrical band of amorphous ice deposited on an aluminum substrate at  $T < 20 \text{ K}$ <sup>5</sup> and the spectrum of that same ice converted to the crystalline form by warming to 155 K and then recooling to 13 K. The ice/silicate bands at  $230 \text{ cm}^{-1}$  ( $43.5 \mu\text{m}$ ) and  $162 \text{ cm}^{-1}$  ( $61.7 \mu\text{m}$ ) correspond to the assigned positions for the transverse optical and longitudinal acoustic modes of crystalline  $\text{H}_2\text{O}$ , respectively (Bertie, Labbé, & Whalley 1969). The  $\text{H}_2\text{O/silicate}$  spectrum is best fitted with the spectrum of crystalline  $\text{H}_2\text{O}$  on aluminum. The addition of an amorphous ice component broadens the spectrum resulting in a poorer fit.

Figure 3 shows the results for  $\text{CH}_3\text{OH}$  deposits in the O—H stretch region at  $T < 20 \text{ K}$ . Figure 3a is the spectrum of  $\text{CH}_3\text{OH/silicate}$ . This is compared in Figure 3b with the broad band typical of amorphous phase  $\text{CH}_3\text{OH}$  deposited on an aluminum substrate at  $T < 20 \text{ K}$ , as well as the double-peaked spectrum of crystalline phase  $\text{CH}_3\text{OH}$  formed by warming amorphous ice to 130 K followed by recooling to  $T < 20 \text{ K}$ . The number of peaks and relative intensities of the  $\text{CH}_3\text{OH/silicate}$  spectrum correspond closely to the spectrum of crystalline  $\text{CH}_3\text{OH}$ . The two most intense peaks of the  $\text{CH}_3\text{OH/silicate}$  ice occur at  $3297 \text{ cm}^{-1}$  and  $3180 \text{ cm}^{-1}$ ; these

<sup>5</sup> The actual intensity of each amorphous ice deposit has been maintained in Figs. 2b, 3b, and 4b relative to the intensity of each annealed crystalline ice which has been normalized to 1.

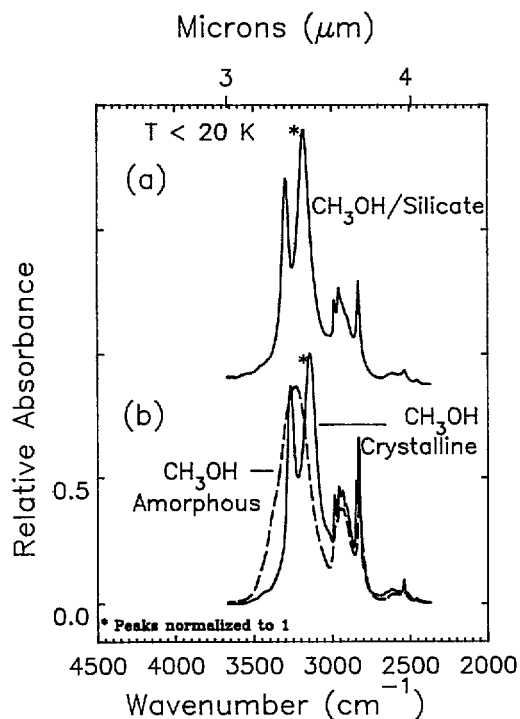


FIG. 3.—Infrared spectra of the O—H stretch region of  $\text{CH}_3\text{OH}$ .  $\text{CH}_3\text{OH}$  deposited directly onto a silicate smoke (a) is compared with spectra of  $\text{CH}_3\text{OH}$  in the amorphous and crystalline phase on an aluminum substrate (b).

are blueshifted more than  $30\text{ cm}^{-1}$  from the corresponding  $3264\text{ cm}^{-1}$  and  $3140\text{ cm}^{-1}$  features in the crystalline phase formed on the aluminum substrate (Fig. 3b). Curve fitting the  $\text{CH}_3\text{OH}/\text{silicate}$  spectrum gives a best fit for 100% crystalline.<sup>6</sup>

The mid-infrared spectrum of the  $\nu_2$  region of  $\text{NH}_3$  near  $10\text{ }\mu\text{m}$  is shown in Figure 4. Figure 4a is the spectrum of the  $\text{NH}_3/\text{silicate}$ . This spectrum is compared with Figure 4b showing the single broad band of amorphous  $\text{NH}_3$  deposited on an aluminum substrate at  $T < 20\text{ K}$ . This deposit was annealed by warming it to near  $110\text{ K}$  and then returning to  $T < 20\text{ K}$ , resulting in the development of two overlapping peaks at  $1097\text{ cm}^{-1}$  ( $9.1\text{ }\mu\text{m}$ ) and  $1069\text{ cm}^{-1}$  ( $9.4\text{ }\mu\text{m}$ ); most likely these are related to the metastable phase of  $\text{NH}_3$  (Sill, Fink, & Ferraro 1980). Complete annealing of this metastable phase to produce a pure crystalline phase was never accomplished before the sample was lost due to rapid sublimation. A spectrum typical of cubic crystalline  $\text{NH}_3$  (also shown in Fig. 4b) was obtained by depositing  $\text{NH}_3$  on an aluminum substrate at  $88\text{ K}$  and then cooling to  $T < 20\text{ K}$ . The single-peaked spectrum of  $\text{NH}_3/\text{silicate}$  has a  $\text{FWHM} = 13\text{ cm}^{-1}$  which is comparable to cubic  $\text{NH}_3$  ( $\text{FWHM} = 14\text{ cm}^{-1}$ ); in contrast, amorphous  $\text{NH}_3$  has a  $\text{FWHM} = 65\text{ cm}^{-1}$ . The peak of the  $\text{NH}_3/\text{silicate}$  ice occurs at  $1065\text{ cm}^{-1}$  ( $9.4\text{ }\mu\text{m}$ ); it is blueshifted from the  $1052\text{ cm}^{-1}$  ( $9.5\text{ }\mu\text{m}$ ) peak of cubic crystalline  $\text{NH}_3$  on aluminum. Curve fitting  $\text{NH}_3/\text{silicate}$  spectra was not straightforward and will require further examination.

Our laboratory results demonstrate that the formation of predominantly crystalline phase ice occurs during direct

<sup>6</sup> Lorentzian curves were fitted to the features of pure amorphous and pure crystalline  $\text{CH}_3\text{OH}$  on aluminum. These template curves were fitted to the  $\text{CH}_3\text{OH}/\text{silicate}$  spectra. The percent crystalline, for example, was assigned as the fraction of the total peak area of the  $\text{CH}_3\text{OH}/\text{silicate}$  represented by the area of the pure crystalline template.

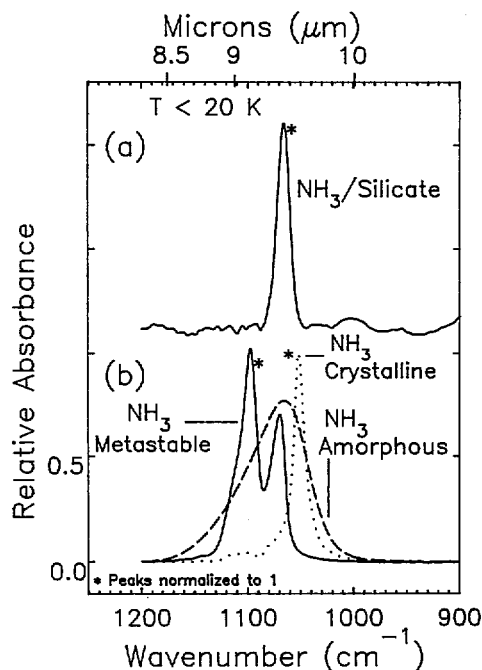


FIG. 4.—Infrared spectra of the  $\nu_2$  region of  $\text{NH}_3$ . The  $\text{NH}_3/\text{silicate}$  spectrum (a) is compared with spectra of  $\text{NH}_3$  deposited on an aluminum substrate in the amorphous phase and after annealing that ice to  $110\text{ K}$  forming a metastable phase (b). The single peak of cubic crystalline  $\text{NH}_3$  (b) was formed by deposition at  $88\text{ K}$ .

deposit of  $\text{H}_2\text{O}$ ,  $\text{CH}_3\text{OH}$ , or  $\text{NH}_3$  at  $T < 20\text{ K}$  on amorphous silicate smoke substrates. Some intersample variations occur with silicate smokes prepared at different times (i.e., there is variation in the  $10/20\text{ }\mu\text{m}$  ratio and the strength of the  $11.4\text{ }\mu\text{m}$  feature), but all amorphous silicate substrates resulted in the low-temperature crystallization (LTC) of condensed ices.

The LTC effect was *not* observed when the ice was formed either on a loose powder made from grains ( $3\text{--}100\text{ }\mu\text{m}$  in size) ground from bulk materials (e.g., Si,  $\text{SiO}$ ,  $\text{SiO}_2$ ), or on  $\text{Si}_2\text{O}_3$  smoke scraped from the inside of our laboratory nucleation chamber. Electron microscope images showed that the scraped off  $\text{Si}_2\text{O}_3$  had a clump size near several microns. These results suggest that the size of the grains may be an important variable.

It is possible that the LTC occurs due to molecular interactions between the condensing gas and reactive silicate surface sites. Silicate smokes are thought to be a disordered network structure with numerous imperfections and dangling bonds (our preliminary EPR measurements show the presence of unpaired electrons). The hypothesis that active sites (perhaps oxygen vacancy defect sites) may facilitate crystallization is supported by the observation that a lower percent crystalline ice/silicate component formed on smoke substrates partly oxidized by heating to  $325^\circ\text{C}$  in air (reducing the number of defect sites). It seems unlikely that any surface reactivity responsible for the LTC would extend beyond an ice layer, yet we find that a second layer of ice is also crystalline in phase. This raises the question of whether the LTC effect is connected to the silicate's low thermal conductivity (TC) or whether some of the active sites were left exposed and in turn facilitated the LTC in the second ice layer.

We have completed a variety of preliminary experiments to examine the possibility that the LTC of ice results from a large

thermal gradient within the silicate smoke. A summary of results are as follows:

1. LTC results on all similarly oxidized silicates regardless of smoke thickness with essentially the same percentage of crystalline component.
2. CO and CH<sub>4</sub> condense on smokes, and CH<sub>4</sub> condenses as a second layer on ultrathin amorphous and on a crystalline CH<sub>3</sub>OH/silicate. As temperature probes, CO and CH<sub>4</sub> deposits imply temperatures below 50 K.
3. The temperature-dependent vaporization of ices from ice/silicate smokes appears to be the same as ice films from aluminum substrates.

These results are supported by the following calculations: We assumed that TC for a silicate smoke was  $22 \mu\text{W cm}^{-1} \text{K}^{-1}$ , the measured value for Cab-O-Sil® (Cabot Corporation 1990) at 100 K (Kuhn et al. 1960). Cab-O-Sil is fumed, amorphous SiO<sub>2</sub>, with a porosity and grain size similar to our smokes. The estimated heat flux from the 300 K surfaces seen by the silicate is  $\sim 100$  times less than the heat flux conducted through a 0.5 mm smoke layer at 120 K. Therefore, the surface temperature of a smoke on a 20 K substrate would remain too cold to crystallize ices during direct deposit. The heat of condensation released during deposition was estimated to be an insignificant heat input using the deposition rates typical of these experiments. It has been reported that the TC of slowly deposited

H<sub>2</sub>O at 135 K is  $4.5 \mu\text{W cm}^{-1} \text{K}^{-1}$  (Sack 1992), a value about one-fifth the TC of Cab-O-Sil. Even if the ice/silicate TC drops by one-fifth after coating with ice, the heat conducted is still estimated to be an order of magnitude above the incoming heat flux.

These experimental results impact ideas concerning the nature of ices in both cometary and interstellar environments. The hypothesis that most ices are accreted in the amorphous phase is central to theories proposing a relationship between cometary outbursts and the exothermic amorphous-to-crystalline phase change of an outer ice layer, and to the idea that volatile species are trapped directly in amorphous deposits and in any clathrate structures formed during warming. The accretion of amorphous phase water ice is the basis for deducing differences in interstellar cloud environments, particularly the temperatures of such clouds. Changes in the 3.1  $\mu\text{m}$  water band have been interpreted in terms of cloud evolution (Graham & Chen 1991; Pendleton, Tielens, & Werner 1990).

The assumption that amorphous phase deposits routinely form on grains at  $T < 80$  K needs to be reexamined in light of our laboratory results. We have considered different possibilities for the mechanisms responsible for the LTC, but there are a large number of parameters. Studies are underway to analyze in detail the composition and morphology of the silicate smokes and the ice/silicate mixtures. However, the astrophysical implications of these results remain unchanged regardless of the mechanism.

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