

## Radiation chemistry of H<sub>2</sub>O + O<sub>2</sub> ices

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### Abstract

The chemistry and spectroscopy of proton-irradiated H<sub>2</sub>O + O<sub>2</sub> ices have been investigated in relation to the production of oxidants in icy satellite surfaces. Hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>), ozone (O<sub>3</sub>), and the hydroperoxy (HO<sub>2</sub>) and hydrogen trioxide (HO<sub>3</sub>) radicals have all been observed, and their temperature and dose dependent production trends have been measured. We find that O<sub>2</sub> aggregates form during the growth of H<sub>2</sub>O + O<sub>2</sub> ice films, and the presence of these aggregates greatly affects the HO<sub>2</sub> and H<sub>2</sub>O<sub>2</sub> yields. In addition, we have found that the position of the spectral maximum of the  $\nu_3$  vibration of O<sub>3</sub> shifts with ice composition, giving an indication of the degree of dispersion of O<sub>3</sub> molecules within the ice. We discuss the relevance of these measurements to icy satellite surfaces.

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### 1. Introduction

Energetic radiation can alter the chemical composition of icy satellite surfaces. In the case of Europa, the destruction of water-ice is the likely source of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>) (Carlson et al., 1999) and molecular oxygen (O<sub>2</sub>) (Spencer and Calvin, 2002) that have been detected there. However, O<sub>2</sub> formation is not unique to Europa, and mounting evidence indicates that it may occur on other icy satellites exposed to high-energy radiation. Molecular oxygen was first identified on Ganymede (Spencer et al., 1995), and has also been detected on Callisto (Spencer and Calvin, 2002). In addition, an O atom atmosphere associated with Saturn's rings suggests that O<sub>2</sub> is formed from the water-ice particles found there (Johnson et al., 2006). Other species that may exist in irradiated water-ice include O<sub>3</sub>, which has been tentatively assigned to spectroscopic features of Ganymede (Noll et al., 1996; Hendrix et al., 1999), Rhea and Dione (Noll et al., 1997), the hydroxyl (OH), hydroperoxy (HO<sub>2</sub>), and hydrogen trioxide

(HO<sub>3</sub>) radicals. The latter three species have yet to be detected in outer Solar System ices.

The spectroscopic characteristics of a molecule are often sensitive to the environment in which it is found. For example, the spectra of H<sub>2</sub>O molecules are different in the crystalline and amorphous phases (Bergren et al., 1978; Hudson and Moore, 1992; Hudgins et al., 1993). In addition, the solid-phase spectra of many compounds are not the same in the presence and absence of water-ice (Hudgins et al., 1993; Ehrenfreund et al., 1997), mainly because of the strong hydrogen bonds present in the host water-ice matrix. These shifts are hard to predict from theory, and so experimental data are essential in characterizing them.

We recently have identified an absorption band of HO<sub>2</sub> in irradiated H<sub>2</sub>O + O<sub>2</sub> ice mixtures at 1142 cm<sup>-1</sup> (Cooper et al., 2006). This radical has previously been detected in irradiated ice (Bednarek et al., 1998) via electron paramagnetic resonance spectroscopy, but an infrared spectroscopic identification had until recently been elusive. In the same paper (Cooper et al., 2006), we also identified HO<sub>3</sub> by an absorption band at 1259 cm<sup>-1</sup>. These radicals formed readily in H<sub>2</sub>O + O<sub>2</sub> ices at 9 K and were found to also be formed at 80 K in lesser amounts. In addition, the same absorption features were observed in irradiated H<sub>2</sub>O + O<sub>3</sub> and neat H<sub>2</sub>O<sub>2</sub> ices. The latter measurement

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was recently supported by Loeffler et al. (2006a), who observed the same bands in irradiated  $\text{H}_2\text{O}_2$  ice but did not identify them. Laffon et al. (2006) recently showed the production of  $\text{H}_2\text{O}_2$ , OH,  $\text{HO}_2$ ,  $\text{O}_2$  and O atoms simultaneously using near-edge X-ray absorption fine structure (NEXFAS) spectroscopy on  $\text{H}_2\text{O}$  ice.  $\text{HO}_3$  and  $\text{O}_3$  however, were not detected in their study.

Chemical species such as  $\text{H}_2\text{O}_2$ ,  $\text{O}_2$ ,  $\text{O}_3$ , OH,  $\text{HO}_2$ , and  $\text{HO}_3$  are of interest to planetary scientists because they are all oxidizing agents and may provide a source of chemical energy in outer Solar System ices. Such energy has been proposed as a possibility to fuel extraterrestrial life (Chyba, 2000), although highly reactive OH and  $\text{HO}_2$  are known to have a negative effect on living organisms. If sufficient quantities of these species exist in an ice then they may act as sterilizing agents.

Expanding on our previous work (Cooper et al., 2006), here we provide the dose and temperature dependence of  $\text{HO}_2$  and  $\text{HO}_3$  formation and stability in order to evaluate these radicals' potential relevance to Europa, Ganymede, and other icy outer Solar System bodies. These two radicals may also form in icy grains of the interstellar medium via cosmic ray bombardment. The formation and destruction mechanisms will be discussed, as will some other interesting observations, including a shift in the  $\text{O}_3$  ( $\nu_3$ ) band position under different formation conditions.

## 2. Experimental methods

The experimental apparatus has been described earlier (Hudson and Moore, 1995), but in short, the ice samples were prepared by mixing the desired gas constituents in a vacuum manifold before depositing them onto a polished aluminum target held at  $\sim 9$  K by a closed-cycle helium refrigerator (APD HC-4). The pressure of the gas mixtures were measured with a capacitance manometer and were prepared such that the mixing ratios between experiments were not different by more than 2%. Water was introduced into the vacuum manifold from a bulb containing Millipore-purified  $\text{H}_2\text{O}$ , degassed using several freeze-pump-thaw cycles with liquid nitrogen. Research grade  $^{16}\text{O}_2$  (Matheson Tri-Gas) or  $^{18}\text{O}_2$  (Isotec; purity of  $>97\%$ ) was mixed with  $\text{H}_2\text{O}$  using standard manometric techniques. Ice samples were warmed at a rate of 2–3 K per minute to their irradiation temperature. The ices were then irradiated with 0.8 MeV protons from a Van de Graaff accelerator. Typical beam currents were  $1 \times 10^{-7}$  A over a sample area of  $5.06 \text{ cm}^2$ . Radiation doses were determined by counting the proton fluence ( $\text{p}^+ \text{ cm}^{-2}$ ) at the ice, and converting to a common scale of eV/16-amu molecule. Proton stopping powers and ranges were calculated with Ziegler's SRIM program (Ziegler et al., 1985; www.srim.org).

Sample thickness was measured by counting laser interference fringes and was typically  $3 \mu\text{m}$ , deposited at a rate of  $\sim 0.1 \mu\text{m min}^{-1}$ . Three-micron thick samples were found to be ideal as they provided the best signal-to-noise ratio for the small  $\text{HO}_2$  and  $\text{HO}_3$  absorptions, without too much baseline distortion from the strongly absorbing water bands. A gas-phase  $\text{H}_2\text{O} + \text{O}_2$  ratio of 6:1 was chosen for many experiments, as a compromise between 1:1, which gave poor quality spectral

characteristics, and 10:1, which gave low  $\text{HO}_2$ ,  $\text{HO}_3$ , and  $\text{O}_3$  absorbances.

Infrared (IR) spectra were obtained using a Nexus 670 spectrometer with 100 scans at  $4 \text{ cm}^{-1}$  resolution. In these experiments the IR beam was transmitted through the ice, reflected from the metal mirror substrate, and was re-transmitted through the ice before entering the detector. Each sample spectrum was divided by the reflectance spectrum from the blank aluminum substrate. These ratioed spectra were converted to  $-\log(I/I_0)$  (absorbance), and band areas of interest ( $\text{H}_2\text{O}_2$ ,  $\text{O}_3$ ,  $\text{HO}_2$ , and  $\text{HO}_3$ ) were calculated using an appropriate baseline.

## 3. Results

### 3.1. Irradiation of $\text{H}_2\text{O} + \text{O}_2$ ices

Fig. 1 shows the mid-IR spectral changes of a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice mixture irradiated to a dose of 0.6 eV/16 amu at 9 K. Before irradiation, the forbidden  $1550 \text{ cm}^{-1}$   $\text{O}_2$  vibrational band was also observed along with a small  $\text{CO}_2$  impurity at  $2344 \text{ cm}^{-1}$ . After irradiation new features at  $2850$ ,  $1259$ ,  $1142$ , and  $1039 \text{ cm}^{-1}$  appear in the spectrum. A decrease in intensity of the small absorption band at  $\sim 3673 \text{ cm}^{-1}$ , associated with dangling OH bonds (Rowland et al., 1991), was observed and was consistent with previous work (Palumbo, 2006). The  $2850$  and  $1039 \text{ cm}^{-1}$  absorptions produced by  $\text{H}_2\text{O}_2$  and  $\text{O}_3$ , respectively, have been well characterized in the literature (Moore and Hudson, 2000; Loeffler et al., 2006a). The  $1259$  and  $1142 \text{ cm}^{-1}$  features were recently shown in this laboratory to be from  $\text{HO}_3$  and  $\text{HO}_2$  radicals (Cooper et al., 2006), respectively.

### 3.2. Spectral positions of ozone

Fig. 2 shows the  $\nu_3$  band of  $\text{O}_3$  synthesized in different ices. Spectrum (a) is the absorption of  $\text{O}_3$  produced in a 1:1  $\text{H}_2\text{O} +$

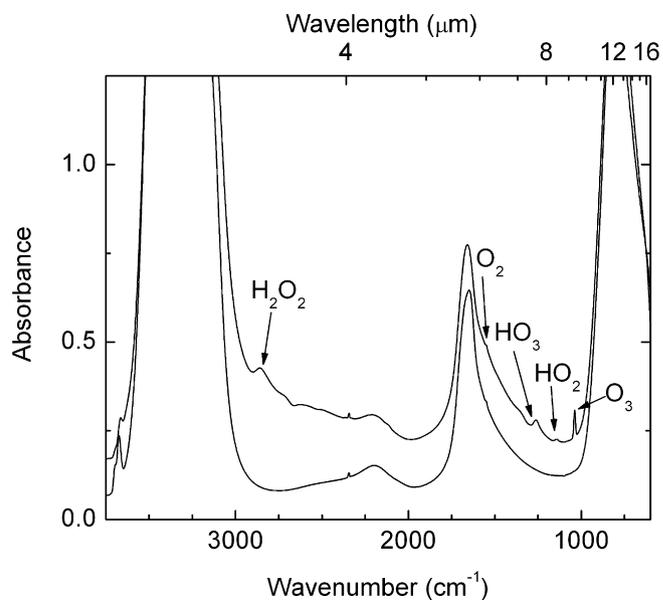


Fig. 1. Spectral changes in a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9 K after a radiation dose of 0.6 eV/16 amu.

$\text{O}_2$  irradiated mixture at 10 K while (b) is from a 6:1  $\text{H}_2\text{O} + \text{O}_2$  sample at the same temperature. Warming the latter ice mixture to 50 K produced (c) and further warming to 125 K gave (d).

### 3.3. $\text{H}_2\text{O}_2$ formation vs temperature

Fig. 3 shows the change in area of the  $\sim 2850 \text{ cm}^{-1}$   $\text{H}_2\text{O}_2$  absorption band with increasing radiation dose for 6:1  $\text{H}_2\text{O} + \text{O}_2$  ices at temperatures between 9 and 100 K. The band appears as a shoulder on the much larger water absorption at  $\sim 3300 \text{ cm}^{-1}$ . The  $\text{H}_2\text{O}_2$  band area was calculated by sub-

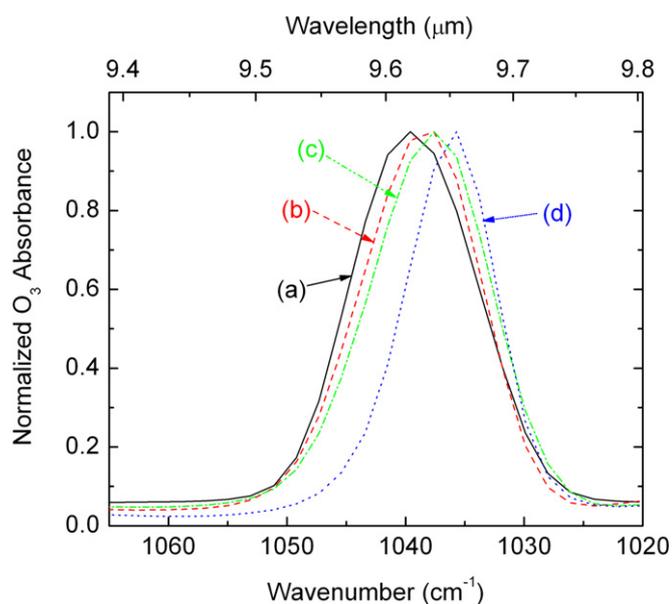


Fig. 2. The  $\nu_3$  band of  $\text{O}_3$  formed in an (a) irradiated 1:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 10 K; (b) irradiated 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9 K; (c, b) warmed to 50 K; (d, b) warmed to 125 K.

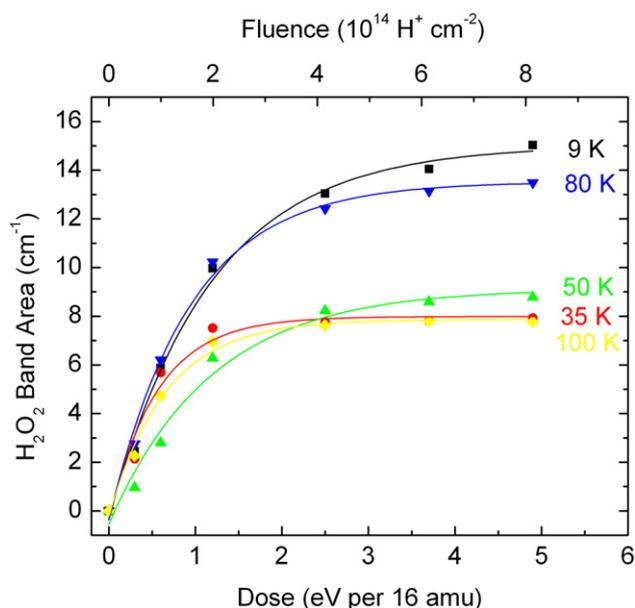


Fig. 3. The dose dependent formation of  $\text{H}_2\text{O}_2$  in a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9, 35, 50, 80, and 100 K.

tracting the spectrum of a pre-irradiated ice mixture from the spectrum of a post-irradiated ice mixture. The band area was then measured from this difference spectrum.

$\text{H}_2\text{O}_2$  is formed slowly at first and then begins to rapidly increase with higher dose, before reaching a steady-state value at  $\sim 5 \text{ eV}/16 \text{ amu}$ . The value of the  $\text{H}_2\text{O}_2$  area at steady-state decreases between 9 and 35 K and is probably associated with the loss of  $\text{O}_2$  from the sample at  $\sim 30 \text{ K}$ . The area then does not change much between 35 and 50 K but rises at 80 K, before decreasing once again at 100 K.

### 3.4. $\text{O}_3$ formation vs temperature

Fig. 4 shows the change in the  $1039 \text{ cm}^{-1}$   $\text{O}_3$  band area as a function of dose and temperature. The saturated peak intensity decreases with increasing irradiation temperature. The large drop in band area between 9 and 35 K is due to the sublimation of some of the  $\text{O}_2$  at  $\sim 30 \text{ K}$ . Loeffler et al. (2006b) report that  $\sim 3\%$   $\text{O}_2$  and  $\text{O}_3$  is lost between 35 and 90 K in similar experiments. This is too small to account for the  $\sim 50\%$  decrease in  $\text{O}_3$  absorbance between 35 and 80 K in the present experiments. However, Chaabouni et al. (2000) report a steady decrease in  $\text{O}_3$  desorption between 60 and 100 K in their non-irradiation experiments and attributed it to a physisorbed state of  $\text{O}_3$  that is identical in spectral properties to pure  $\text{O}_3$ . Other possibilities to explain the lower yield of  $\text{O}_3$  include a temperature dependent formation mechanism or an increased destruction rate at higher temperature.

### 3.5. $\text{HO}_2$ formation vs temperature

Fig. 5 shows the integrated absorbance of the  $\text{HO}_2$  band as a function of dose at temperatures between 9 and 100 K. The band area of the  $\text{HO}_2$  feature decreased with increasing temperature except for the initial change in temperature between 9 and

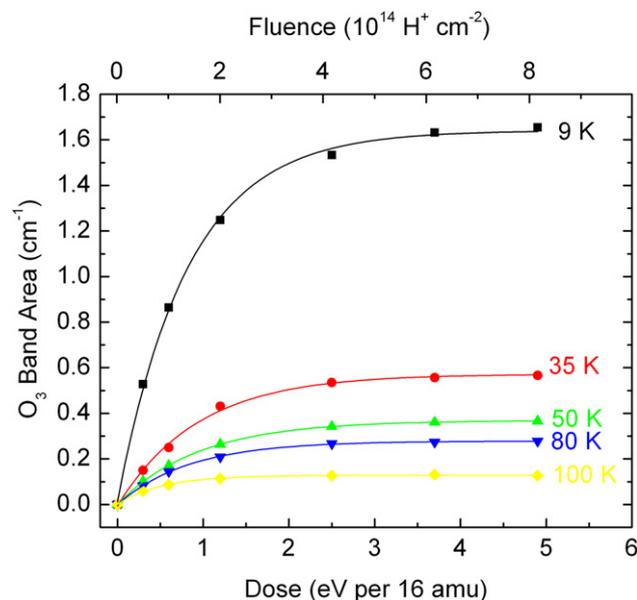


Fig. 4. The dose dependent formation of  $\text{O}_3$  in a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9, 35, 50, 80, and 100 K.

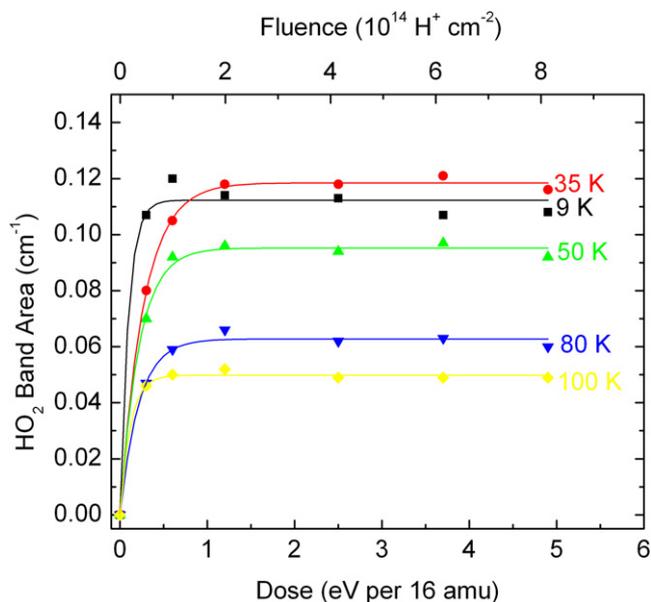


Fig. 5. The dose dependent formation of  $\text{HO}_2$  in a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9, 35, 50, 80, and 100 K.

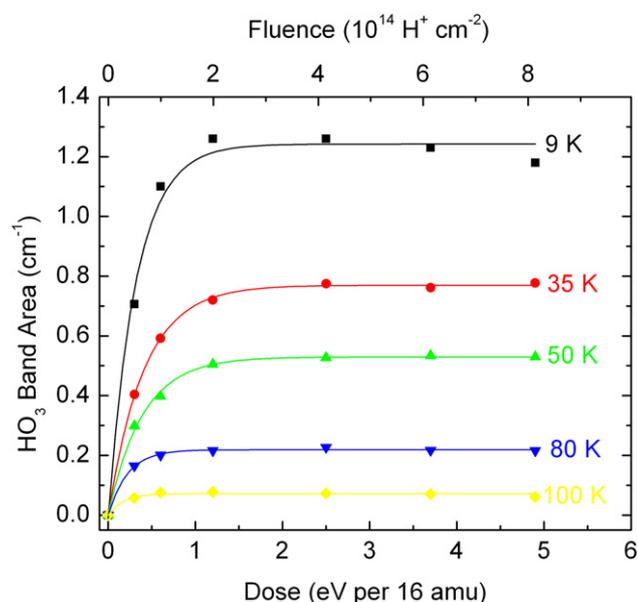


Fig. 6. The dose dependent formation of  $\text{HO}_3$  in a 6:1  $\text{H}_2\text{O} + \text{O}_2$  ice at 9, 35, 50, 80, and 100 K.

35 K. The amount of  $\text{HO}_2$  produced was about the same at 9 and 35 K but was formed slightly slower at 35 K than at 9 K.

### 3.6. $\text{HO}_3$ formation vs temperature

Fig. 6 shows that  $\text{HO}_3$  was formed rapidly with increasing radiation dose. The radical's steady-state IR band area decreases with increasing temperature. We previously showed (Cooper et al., 2006) that  $\text{HO}_3$  forms via H addition to  $\text{O}_3$ , and so a decrease in  $\text{O}_3$  (Fig. 4) reduces the amount of  $\text{HO}_3$  produced and probably accounts for much of the fall in  $\text{HO}_3$  intensity. The  $\text{HO}_3$  decrease between 9 and 35 K was not as great

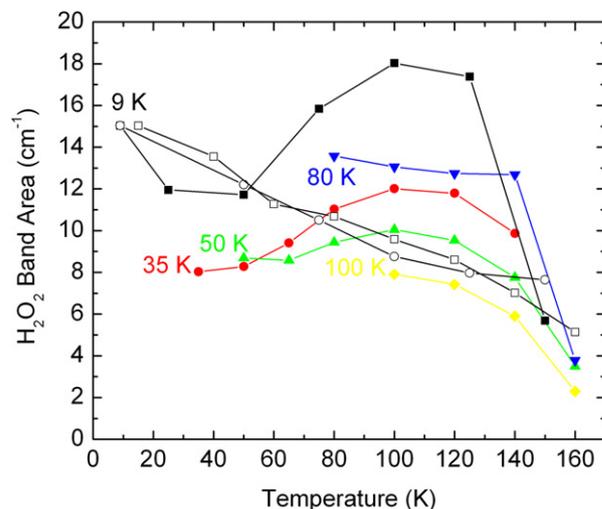


Fig. 7. The change in  $\text{H}_2\text{O}_2$  band area upon warming 6:1  $\text{H}_2\text{O} + \text{O}_2$  ices irradiated to steady-state conditions ( $\sim 5$  eV/16 amu) at 9, 35, 50, 80, and 100 K. The unfilled square and circular data points show the intensity change of the  $\sim 2850$   $\text{cm}^{-1}$  band upon warming unirradiated  $\text{H}_2\text{O}_2$  and a 3%  $\text{H}_2\text{O}_2$  in  $\text{H}_2\text{O}$  ice sample, these data were scaled to the 9 K  $\text{H}_2\text{O} + \text{O}_2$  irradiated data in order to provide a comparison.

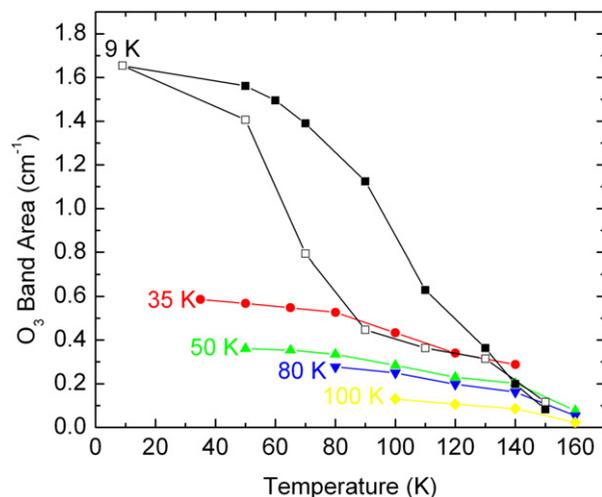


Fig. 8. The change in  $\text{O}_3$  band area upon warming 6:1  $\text{H}_2\text{O} + \text{O}_2$  ices irradiated to steady-state conditions ( $\sim 5$  eV/16 amu) at 9, 35, 50, 80, and 100 K. The unfilled square data points show the  $\text{O}_3$  loss from an unirradiated  $\sim 5:1$   $\text{H}_2\text{O} + \text{O}_3$  mixture, co-deposited at 9 K and warmed to 150 K. The data of the co-deposited mixture was scaled to the 9 K  $\text{H}_2\text{O} + \text{O}_2$  irradiated data in order to provide a comparison.

as it was for  $\text{O}_3$ . The reason for this will be discussed in the next section.

### 3.7. Stability of radiation products vs temperature

Figs. 7 and 8 show the effect of warming an  $\text{H}_2\text{O} + \text{O}_2$  ice irradiated at 9, 35, 50, 80 and 100 K on the area of the  $\text{H}_2\text{O}_2$  and  $\text{O}_3$  bands, respectively.  $\text{H}_2\text{O}_2$  displays a complex behavior during warming while  $\text{O}_3$  decreases with increasing temperature. The band area of  $\text{H}_2\text{O}_2$  in ices irradiated at 9, 35, and 50 K increases to a maximum between 80 and 120 K before decreasing again. In 80 and 100 K irradiated samples, the  $\text{H}_2\text{O}_2$  band

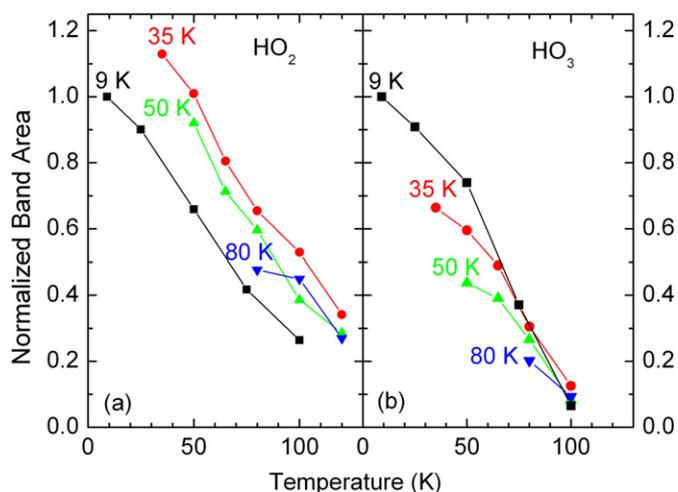


Fig. 9. The change in (a) HO<sub>2</sub> and (b) HO<sub>3</sub> normalized band area upon warming 6:1 H<sub>2</sub>O + O<sub>2</sub> ices irradiated to steady-state conditions ( $\sim 5$  eV/16 amu) at 9, 35, 50, and 80 K.

area stays about the same before decreasing above 120 K. For comparison, Fig. 7 also shows the effect of temperature on non-irradiated pure H<sub>2</sub>O<sub>2</sub> and 3% H<sub>2</sub>O<sub>2</sub> in H<sub>2</sub>O ice samples. The observed band area of H<sub>2</sub>O<sub>2</sub> in each of these samples actually decreases monotonically with increasing temperature. Loeffler et al. (2006b) found that H<sub>2</sub>O<sub>2</sub> sublimation was negligible at the temperatures used in the present experiments. Therefore the monotonic decrease could be due to a decrease in the intrinsic band strength with increasing temperature or is only apparently decreasing due to changes in the optical properties of the sample. Nevertheless, we observe an increase in the band area of H<sub>2</sub>O<sub>2</sub> in irradiated H<sub>2</sub>O + O<sub>2</sub> ices as the temperature is increased up to 120 K.

The band area of O<sub>3</sub> plotted in Fig. 8 decreases monotonically with increasing temperature. For comparison, Fig. 8 also shows the effect of warming an unirradiated 5:1 H<sub>2</sub>O + O<sub>3</sub> mixture co-deposited at 9 K. The O<sub>3</sub> in this unirradiated sample was less thermally stable than O<sub>3</sub> formed radiolytically in irradiated H<sub>2</sub>O + O<sub>2</sub> mixtures.

Fig. 9 shows the effect of warming an irradiated H<sub>2</sub>O + O<sub>2</sub> ice on the amount of HO<sub>2</sub> and HO<sub>3</sub>. For comparison, the integrated bands have been normalized to the area for that species at 9 K. Both radicals are destroyed as the temperature increases, however HO<sub>2</sub> destruction was slightly slower and persists to a slightly higher temperature.

Fig. 10 compares the formation and destruction data at 80 K for all four species of interest. The H<sub>2</sub>O<sub>2</sub> band areas have been reduced by a factor of 25 in order to display the peroxide on the same graph with O<sub>3</sub>, HO<sub>2</sub>, and HO<sub>3</sub>.

## 4. Discussion

### 4.1. O<sub>2</sub> and O<sub>3</sub>

In the present experiments, different H<sub>2</sub>O:O<sub>2</sub> ratios were used, to elucidate the variety of O<sub>2</sub> bonding environments in the ice. Some O<sub>2</sub> molecules will have only H<sub>2</sub>O molecules as

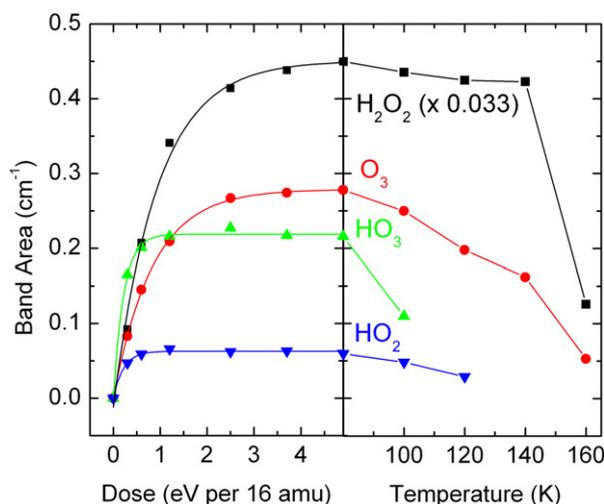


Fig. 10. A comparison of the formation of H<sub>2</sub>O<sub>2</sub>, O<sub>3</sub>, HO<sub>2</sub>, and HO<sub>3</sub> at 80 K and subsequent removal (via either destruction or desorption) from the ice as it is warmed from 80 to 160 K.

nearest neighbors, while some will be surrounded by a mixture of O<sub>2</sub> and H<sub>2</sub>O. At the latter locations, O<sub>2</sub> may actually be present in small aggregates within the ice and will consequently behave like bulk solid oxygen. It is these aggregates that sublime when the sample is warmed to 30 K or higher. Acting as a tracer molecule, the O<sub>3</sub> band position may provide information about the state of O<sub>2</sub> in the ice since the spectral maximum of the  $\nu_3$  mode of O<sub>3</sub> is shifted when bonded to water (Schriver et al., 1990). Under different conditions, the O<sub>3</sub> in our ice samples may form in O<sub>2</sub> aggregates or be dispersed within the ice lattice.

Pure amorphous and crystalline-phase O<sub>3</sub> produce their strongest band,  $\nu_3$ , at 1037.1 and 1026.9 cm<sup>-1</sup>, respectively (Chaabouni et al., 2000), while O<sub>3</sub> in an O<sub>2</sub> matrix produces a band at 1037.8 cm<sup>-1</sup> (Schriver-Mazzuoli et al., 1995) and in an N<sub>2</sub> matrix at 1043.0 cm<sup>-1</sup> (Schriver-Mazzuoli et al., 1996). An O<sub>3</sub> dimer in an O<sub>2</sub> matrix produces a feature at 1042 cm<sup>-1</sup> (Bennett and Kaiser, 2005). Table 1 summarizes band positions for O<sub>3</sub> in various ices. Fig. 2a shows the absorption of O<sub>3</sub> with a spectral maximum at 1039.7 cm<sup>-1</sup> from an irradiated 1:1 H<sub>2</sub>O + O<sub>2</sub> ice mixture at 9 K. At this ratio, O<sub>3</sub> monomers and dimers will be mostly formed in oxygen-rich environments due to the large abundance of O<sub>2</sub>. The absorptions at 1038 and 1042 cm<sup>-1</sup> cannot be resolved at our spectral resolution of 4 cm<sup>-1</sup>, but a contribution from both is likely. Our observed spectral maximum at 1039.7 cm<sup>-1</sup> lies between the maxima for O<sub>3</sub> monomer and dimer, and is consistent with a mixture of both being present.

Spectrum (b) in Fig. 2 shows the ozone absorption in an irradiated 6:1 H<sub>2</sub>O + O<sub>2</sub> ice at 9 K. The lower fraction of O<sub>2</sub> should produce fewer O<sub>3</sub> dimers, and this is supported by the small shift to lower frequency of the spectral maximum to 1038.5 cm<sup>-1</sup>. The O<sub>3</sub> dimer absorption on the high energy side of the band is reduced, resulting in a shift of the peak. When the sample is warmed to 50 K, the maximum shifts to 1037.8 cm<sup>-1</sup> (c), and then again to 1035.8 cm<sup>-1</sup> (d) at 125 K. These changes can be interpreted as the gradual loss of O<sub>3</sub> monomer from an

Table 1  
Summary of spectral maxima ( $\text{cm}^{-1}$ ) for the  $\nu_3(\text{O}_3)$  absorption band

$\text{O}_3$ in various ices	Temperature (K)	$\nu_3$ band position ( $\text{cm}^{-1}$ )	Reference
Pure amorphous $\text{O}_3$	11	1037.1	Chaabouni et al. (2000)
Pure crystalline $\text{O}_3$	55	1026.9	Chaabouni et al. (2000)
$\text{O}_3$ in an $\text{N}_2$ matrix	11	1043.0	Schrivier-Mazzuoli et al. (1996)
$\text{O}_3$ in an $\text{O}_2$ matrix	11	1037.8	Schrivier-Mazzuoli et al. (1995)
$\text{O}_3$ dimer in an $\text{O}_2$ matrix	11	1042	Bennett and Kaiser (2005)
$\text{O}_3$ H-bonded to $\text{H}_2\text{O}$	11	1033.8	Schrivier et al. (1990)
$\text{O}_3$ formed in irradiated $\text{H}_2\text{O} + \text{O}_2$ ices			
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (1:1)	9	1039.7	This work
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (6:1)	9	1038.5	This work
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (10:1)	9	1036.2	This work
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (6:1)	9	1038.5	This work
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (6:1)	50	1037.8	This work
$\text{O}_3$ in $\text{H}_2\text{O} + \text{O}_2$ (6:1)	125	1035.8	This work

Table 2  
Parameters<sup>a</sup> for the pseudo-first order formation of  $\text{H}_2\text{O}_2$ ,  $\text{O}_3$ ,  $\text{HO}_2$ , and  $\text{HO}_3$

		9 K	35 K	50 K	80 K	100 K
$\text{H}_2\text{O}_2$	$b$	$15.3 \pm 0.5$	$8.3 \pm 0.7$	$9.7 \pm 0.6$	$13.9 \pm 0.5$	$8.0 \pm 0.4$
	$k$	$0.85 \pm 0.08$	$1.8 \pm 0.3$	$0.85 \pm 0.16$	$1.1 \pm 0.1$	$1.5 \pm 0.2$
$\text{O}_3$	$b$	$1.6 \pm 0.01$	$0.77 \pm 0.01$	$0.36 \pm 0.01$	$0.28 \pm 0.01$	$0.13 \pm 0.01$
	$k$	$1.2 \pm 0.1$	$1.1 \pm 0.1$	$1.1 \pm 0.1$	$1.2 \pm 0.1$	$2.0 \pm 0.2$
$\text{HO}_2$	$b$	$0.11 \pm 0.01$	$0.12 \pm 0.01$	$0.10 \pm 0.01$	$0.06 \pm 0.01$	$0.05 \pm 0.01$
	$k$	$10.6 \pm 3.6$	$3.7 \pm 0.2$	$4.6 \pm 0.3$	$4.7 \pm 0.5$	$8.6 \pm 0.9$
$\text{HO}_3$	$b$	$1.3 \pm 0.03$	$0.77 \pm 0.01$	$0.53 \pm 0.01$	$0.22 \pm 0.01$	$0.07 \pm 0.01$
	$k$	$3.2 \pm 0.3$	$2.5 \pm 0.05$	$2.6 \pm 0.1$	$4.6 \pm 0.7$	$6.3 \pm 2.2$

<sup>a</sup> Data were fitted to the equation  $A(D) = b(1 - e^{-kD})$ , where  $D$  is the dose in  $\text{eV}/16 \text{ amu}$  and  $A$  is band area with units of  $\text{cm}^{-1}$ .  $b$  and  $k$  are constants that characterize each curve.

$\text{O}_2$ -rich environment, and the gradual increase of  $\text{O}_3$  monomer that is hydrogen bonded to the water molecules in the ice. The latter is known to absorb at  $1033.8 \text{ cm}^{-1}$  (Schrivier et al., 1990). A 10:1  $\text{H}_2\text{O} + \text{O}_2$  ice was also irradiated at 9 K and the  $\text{O}_3$  band position was observed at  $1036.2 \text{ cm}^{-1}$  (not shown), entirely in agreement with our reasoning above.

Table 1 also summarizes the positions of the  $\nu_3$  band of  $\text{O}_3$  for  $\text{H}_2\text{O} + \text{O}_2$  ices irradiated at different temperatures. From these data we can see that the  $\text{O}_2$  sublimation at  $\sim 30 \text{ K}$  has a large effect on the spectral position of  $\text{O}_3$ . The  $\text{O}_2$  sublimating from the sample comes from aggregates behaving like bulk solid  $\text{O}_2$ , while the  $\text{O}_2$  molecules remaining trapped are more dispersed in the water-ice, producing a greater fraction of hydrogen-bonded  $\text{O}_3$  molecules. This is consistent with the shift of the band maximum to a lower frequency (see Table 1). The amount of  $\text{O}_3$  also significantly decreases with increasing irradiation temperature as shown in Fig. 4. It is much easier to form  $\text{O}_3$  in an aggregate of  $\text{O}_2$  molecules than from  $\text{O}_2$  molecules dispersed throughout the ice lattice. The data in Fig. 4 are fitted by a first-order curve and the parameters are listed in Table 2. These fits are consistent with a pseudo-first-order reaction mechanism in which  $\text{O}_2$  is in excess (Eq. (1)).



For a comprehensive investigation of the formation mechanism of  $\text{O}_3$  in the solid-state, see the work by Baragiola et al. (1999). Ozone is made from an  $\text{O}_2$  molecule (with the addition of another O atom), and so the degree of dispersion of  $\text{O}_3$  molecules is an indirect measurement of the degree of dispersion of  $\text{O}_2$  molecules. This is important for studies of Ganymede and Europa where the  $\text{O}_2$  is thought to be present in high densities (Calvin et al., 1996) and possibly aggregates where  $\text{O}_2/\text{O}_3$  ‘microatmospheres’ may exist (Johnson and Jesser, 1997).

#### 4.2. $\text{HO}_2$ and $\text{HO}_3$

Further evidence for the presence of  $\text{O}_2$  aggregates at 9 K comes from the band positions of  $\text{HO}_2$  and  $\text{HO}_3$  observed in  $\text{H}_2\text{O} + \text{O}_2$  ices irradiated at different temperatures (Table 3). A shift in the  $\text{HO}_2$  band position occurs above the sublimation point of  $\text{O}_2$  ( $\sim 30 \text{ K}$ ), indicating that  $\text{HO}_2$  is formed in a different location within the ice above this temperature. The lowering of the frequency with increasing irradiation temperature is consistent with an increase in hydrogen bonding of the  $\text{HO}_2$  radical (Nelander, 1997).  $\text{HO}_3$  also shifts with increasing irradiation temperature and again the direction of the shift indicates the effect of increased hydrogen bonding. The higher temperatures may also allow radicals to gain sufficient thermal energy to reposition themselves into lower-energy configurations within the

Table 3  
Spectral positions for  $\tilde{\nu}_3$  (H<sup>16</sup>O<sub>2</sub>) and  $\tilde{\nu}$  (H<sup>16</sup>O<sub>3</sub>) bands formed at different temperatures

	$\tilde{\nu}_3$ (H <sup>16</sup> O <sub>2</sub> ) (cm <sup>-1</sup> )	$\tilde{\nu}_3$ (H <sup>18</sup> O <sub>2</sub> ) (cm <sup>-1</sup> )	Isotopic shift (cm <sup>-1</sup> )	$\tilde{\nu}_3$ (H <sup>16</sup> O <sub>2</sub> )/ $\tilde{\nu}_3$ (H <sup>18</sup> O <sub>2</sub> )
In H <sub>2</sub> O-ice 9 K	1142	1078	64	1.059
In H <sub>2</sub> O-ice 35 K	1135	–	–	–
In H <sub>2</sub> O-ice 50 K	1135	–	–	–
In H <sub>2</sub> O-ice 80 K	1135	1072	63	1.059
In H <sub>2</sub> O-ice 100 K	1135	–	–	–
In Ar	1101.3	1039.7	61.6	1.0592
	$\tilde{\nu}_3$ (H <sup>16</sup> O <sub>3</sub> ) (cm <sup>-1</sup> )	$\tilde{\nu}_3$ (H <sup>18</sup> O <sub>3</sub> ) (cm <sup>-1</sup> )	Isotopic shift (cm <sup>-1</sup> )	$\tilde{\nu}_3$ (H <sup>16</sup> O <sub>3</sub> )/ $\tilde{\nu}_3$ (H <sup>18</sup> O <sub>3</sub> )
In H <sub>2</sub> O-ice 9 K	1259	1220	39	1.032
In H <sub>2</sub> O-ice 35 K	1259	–	–	–
In H <sub>2</sub> O-ice 50 K	1258	–	–	–
In H <sub>2</sub> O-ice 80 K	1256	1221	35	1.029
In H <sub>2</sub> O-ice 100 K	1254	–	–	–
In Ar	1223	1190	33	1.028

Positions for the 18-oxygen substituted isotopomers are shown at 9 and 80 K.

lattice. We noted in our previous paper (Cooper et al., 2006) that the isotopic shift for HO<sub>3</sub> in H<sub>2</sub>O at 9 K was higher than the shift in Ar (Table 3). Earlier work (Nelander et al., 2000) showed that the HO<sub>3</sub> vibration we followed has both bending and stretching components, and so we proposed (Cooper et al., 2006) that the ice matrix slightly affects the relative contribution of each vibrational component, due to the geometrical constraints of accommodating the large HO<sub>3</sub> species into the H<sub>2</sub>O lattice. This alteration appears to indeed be the case, because the agreement in the isotopic shift at 80 K and the shift in Ar is much better than at 9 K (Table 3).

If we now turn to the dependence of HO<sub>3</sub> formation on irradiation temperature, Fig. 6 shows that the saturation amount of HO<sub>3</sub> decreases with increasing temperature. This trend is expected from the decrease in the amount of O<sub>3</sub> produced, as just discussed. However, the O<sub>3</sub> decrease between 9 and 35 K is far greater (~62%) than that of HO<sub>3</sub> (~38%) over the same temperature range. The likely reason for this is that at 9 K the O<sub>3</sub> is produced in O<sub>2</sub> clusters, as already discussed. In such an environment, an HO<sub>3</sub> will have a short lifetime because of reactions with other HO<sub>3</sub> and HO<sub>2</sub> radicals made concurrently within the aggregate. Since little is known about HO<sub>3</sub> chemistry, we cannot say whether or not this radical will react with O<sub>2</sub> or O<sub>3</sub> molecules. Overall, reactions of HO<sub>3</sub> will limit this radical's concentration in the clusters in which it is formed, and this results in less HO<sub>3</sub> being produced than might be expected from the amount of O<sub>3</sub> present. HO<sub>3</sub> formed from dispersed O<sub>3</sub> molecules at temperatures greater than 30 K is stabilized by hydrogen bonding (Aloisio and Francisco, 1999) and is the likely source of most of the HO<sub>3</sub> in the ice.

Similarly, the HO<sub>2</sub> band area decreases with increasing temperature except for the 9–35 K interval (Fig. 5). The amounts of HO<sub>2</sub> at 9 and 35 K are roughly the same, although at 35 K it takes a higher radiation dose to reach a steady-state level. Once again, the likely reason why the band area at 9 K is lower than expected from the trend is because of formation and subsequent self-reaction within an O<sub>2</sub> aggregate. Supporting this assertion is the higher than expected amount of H<sub>2</sub>O<sub>2</sub> at 9 K in Fig. 3. The trend between 35 and 80 K is for the H<sub>2</sub>O<sub>2</sub> to increase, but the 9 K data is higher than expected from this trend. The

reason for this is likely to be that H<sub>2</sub>O<sub>2</sub> is produced from the disproportionation of HO<sub>2</sub> in O<sub>2</sub> clusters.



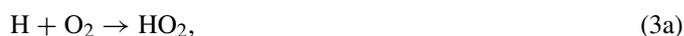
So at 9 K, the HO<sub>2</sub> concentration is limited by the fact that many HO<sub>2</sub> radicals are formed in close proximity in O<sub>2</sub> aggregates through the combination of O<sub>2</sub> and H atoms produced from the dissociation of H<sub>2</sub>O.

The reason for the decrease in HO<sub>2</sub> yield above 35 K may be similar to that of HO<sub>3</sub>. The amount of O<sub>2</sub> remaining in the ice as the temperature is raised will decrease due to thermal desorption. However, this loss cannot be measured via IR spectroscopy because of the extremely weak O<sub>2</sub> IR band. Chemical reactions that occur as the mobility of reactive species increases with higher temperatures may also contribute to the loss of HO<sub>2</sub> above 35 K.

In our earlier paper (Cooper et al., 2006), we discussed possible formation paths for HO<sub>2</sub> and HO<sub>3</sub>. The expanded data here allows us to draw the conclusion that H-addition reactions to O<sub>2</sub> and O<sub>3</sub> are the dominant mechanisms. First-order fits in Figs. 5 and 6 are consistent with pseudo-first-order reactions where O<sub>2</sub> and O<sub>3</sub> are in excess. See Table 2 for a list of the parameters of the curves fitted to the data in Figs. 5 and 6. These first-order fits are made by assuming that the abundances of radiolytic fragments (such as H, OH, and O) are small compared with those of H<sub>2</sub>O, O<sub>2</sub>, and O<sub>3</sub>.

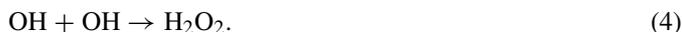
#### 4.3. H<sub>2</sub>O<sub>2</sub>

It would appear that H<sub>2</sub>O<sub>2</sub> can be formed by a number of different mechanisms in the present set of experiments. Three possibilities are H-atom addition (Eqs. (3a) and (3b)), HO<sub>2</sub> disproportionation (Eq. (2)), and OH dimerization (Eq. (4)). It was previously suggested that two H-addition reactions to O<sub>2</sub> (to form HO<sub>2</sub> as an intermediate) enhanced the production of H<sub>2</sub>O<sub>2</sub> in irradiated H<sub>2</sub>O + O<sub>2</sub> ices (Moore and Hudson, 2000; Cooper et al., 2006).



However, at 9 K it is possible that the first step in this sequence (Eq. (3a)) occurs only in O<sub>2</sub> aggregates, so that before a second H can react with HO<sub>2</sub> (Eq. (3b)), two neighboring HO<sub>2</sub> radicals combine with one another, to give H<sub>2</sub>O<sub>2</sub> and O<sub>2</sub> (Eq. (2)). This explains the greater yield of H<sub>2</sub>O<sub>2</sub> at 9 K compared with 35 K. Above 35 K there are fewer O<sub>2</sub> aggregates in the ice, and so HO<sub>2</sub> is formed and remains trapped before reacting with a second H to produce H<sub>2</sub>O<sub>2</sub>.

At 80 K an increase in the H<sub>2</sub>O<sub>2</sub> band area was measured. The most commonly proposed H<sub>2</sub>O<sub>2</sub> formation mechanism, is that of the reaction of two OH radicals.



These OH radicals are strongly hydrogen bonded to water molecules (Cooper et al., 2003), and it is not until ~80 K that they can diffuse within a water-ice lattice (Johnson and Quickenden, 1997, and references therein). Energetic OH radicals can diffuse short distances along ion tracks and react at 10 K, but bulk diffusion probably does not occur. The increase of H<sub>2</sub>O<sub>2</sub> at 80 K we observe is probably an effect of this OH mobility. At 100 K the H<sub>2</sub>O<sub>2</sub> abundance decreases dramatically and reflects increased destruction of H<sub>2</sub>O<sub>2</sub> by reactive species such as OH radicals or electrons.

Without the presence of O<sub>2</sub>, a detectable quantity of H<sub>2</sub>O<sub>2</sub> is not seen at 80 K in our proton-irradiated H<sub>2</sub>O. This result is consistent with Moore and Hudson (2000). They suggested that one possible reason why H<sub>2</sub>O<sub>2</sub> was not observed in proton MeV irradiation experiments on pure H<sub>2</sub>O at 80 K was because of the rapid reaction of electrons with H<sub>2</sub>O<sub>2</sub>. When electron-scavenging molecules were co-deposited with water, H<sub>2</sub>O<sub>2</sub> was produced in observable quantities. Molecular oxygen may also be acting as an electron scavenger in the present experiments via the formation of the superoxide ion (O<sub>2</sub><sup>-</sup>). The role of electron scavenging needs to be investigated further but is beyond the scope of the present paper. Other groups (Loeffler et al., 2006c; Gomis et al., 2004a, 2004b) recently have shown that observable amounts of H<sub>2</sub>O<sub>2</sub> can be produced in pure H<sub>2</sub>O at 80 K using keV-energy ions with greater stopping power than 0.8 MeV protons. Researchers Zheng et al. (2006b) and Loeffler et al. (2006c) reported a decrease in H<sub>2</sub>O<sub>2</sub> production in pure water-ice with increasing temperature, and so electron scavenging may play a critical role in the radiation stability of H<sub>2</sub>O<sub>2</sub> in the ice in our experiments. However, Gomis et al. (2004a) found that the yields were ion-dependent. Ices irradiated with C<sup>+</sup>, H<sup>+</sup> and O<sup>+</sup> ions produced more H<sub>2</sub>O<sub>2</sub> at 77 K than 16 K, while N<sup>+</sup> and Ar<sup>+</sup> had no temperature dependence on the H<sub>2</sub>O<sub>2</sub> yield.

Yet a fourth reaction sequence, discussed by Zheng et al. (2006a), is O addition to H<sub>2</sub>O to form H<sub>2</sub>O<sub>2</sub>. Although there is no clear evidence from our data for it, O atoms produced from the dissociation of O<sub>2</sub> may react with H<sub>2</sub>O to form H<sub>2</sub>O<sub>2</sub>.

#### 4.4. Chemical destruction processes

Once the ice at each temperature was irradiated and the formation of products approached steady-state values, the sample was warmed. Figs. 7–9 show the changes in the band areas of H<sub>2</sub>O<sub>2</sub>, O<sub>3</sub>, HO<sub>2</sub>, and HO<sub>3</sub>. The loss of O<sub>3</sub> (Fig. 8) may be due

to desorption from the ice, but the increased rate of loss between 80 and 120 K may also indicate a chemical destruction mechanism. Fig. 9 shows that HO<sub>2</sub> and HO<sub>3</sub> are destroyed at ~100–120 K. The diffusion of these species in ice is not well understood. The OH radical, which is comparable to HO<sub>2</sub> in terms of the strength of hydrogen bonding (Cooper et al., 2003; Nelander, 1997), is mobile at these temperatures, but the larger physical sizes of HO<sub>2</sub> and HO<sub>3</sub> may significantly affect their mobility.

The increase in H<sub>2</sub>O<sub>2</sub> (Fig. 7) after warming irradiated 9, 35 and 50 K H<sub>2</sub>O + O<sub>2</sub> ices is likely the result of trapped OH radicals reacting with one another, as in Eq. (4). OH radicals are produced at all temperatures. An OH radical that can reactively scatter to within close proximity of another OH can combine with it to form H<sub>2</sub>O<sub>2</sub>, but at the lower temperatures if an OH does not react then it will be trapped in the ice lattice. As the ice's temperature increases, the OH mobility rises and more H<sub>2</sub>O<sub>2</sub> is produced. However, OH also destroys H<sub>2</sub>O<sub>2</sub> as follows:



In other words, the H<sub>2</sub>O<sub>2</sub> concentration is enhanced, but also limited by the availability of OH radicals. In experiments on 80 and 100 K irradiated samples, no increase in H<sub>2</sub>O<sub>2</sub> is seen during warming, indicating that the rate of H<sub>2</sub>O<sub>2</sub> formation is about equal to the rate of destruction of H<sub>2</sub>O<sub>2</sub> by Eq. (5). Alternatively there may not be enough OH radicals stored during irradiation at 80 K or 100 K to bring about an increase in the H<sub>2</sub>O<sub>2</sub> concentration with further warming. For all cases investigated here, the H<sub>2</sub>O<sub>2</sub> concentration starts to decrease once the irradiated ice temperature exceeds 120 K. The decrease between 120 and 160 K is not attributed to thermal losses since Zheng et al. (2006a) showed that warmed H<sub>2</sub>O-ice containing H<sub>2</sub>O<sub>2</sub> did not release H<sub>2</sub>O<sub>2</sub> until the 160–180 K region. The decrease we observe therefore suggests that OH radicals are rapidly destroying H<sub>2</sub>O<sub>2</sub>. The product of this destruction, HO<sub>2</sub> is not seen to increase as would be expected, because it too is rapidly destroyed by OH (Eq. (6)).



However, destruction of HO<sub>2</sub> is slower than that of HO<sub>3</sub>, so that HO<sub>2</sub> remains in the ice to a temperature ~20 K higher than HO<sub>3</sub>. A reservoir of H<sub>2</sub>O<sub>2</sub> would buffer the HO<sub>2</sub> concentration and explains this observation. The chemistry of H<sub>2</sub>O<sub>2</sub>, O<sub>3</sub>, HO<sub>2</sub> and HO<sub>3</sub> with OH radicals is supported by recent work by Loeffler et al. (2006b) who observed a decrease in O<sub>3</sub> concentration in irradiated H<sub>2</sub>O<sub>2</sub> associated with a chemical destruction mechanism. The reaction of OH with O<sub>3</sub> may also produce HO<sub>2</sub> via Eq. (7) and could also contribute to HO<sub>2</sub> persisting in the ice at a higher temperature than HO<sub>3</sub>.



#### 4.5. Relevance to icy satellites

We have shown that irradiated H<sub>2</sub>O + O<sub>2</sub> ices produce H<sub>2</sub>O<sub>2</sub>, O<sub>3</sub>, HO<sub>2</sub>, and HO<sub>3</sub>. All of these new species are ox-

idizing agents and may provide chemical energy to the surfaces and potentially to any liquid sub-surfaces that may exist on the icy Galilean satellites. We have formed H<sub>2</sub>O<sub>2</sub> at much higher amounts in the presence of O<sub>2</sub> in the ice than without, which raises the question of how H<sub>2</sub>O<sub>2</sub> is formed on such icy satellites and what is its relation to the chemistry of O<sub>2</sub>?

The products of liquid water's radiolysis usually are discussed in terms of chemical reactions that first produce OH, then H<sub>2</sub>O<sub>2</sub>, then HO<sub>2</sub>, and finally O<sub>2</sub>. In contrast, the radiolysis of water-ice, which is more relevant to planetary scientists, has proven more difficult to interpret. If the combination of two OH radicals (Eq. (4)) produced from the radiolytic destruction of water-ice forms H<sub>2</sub>O<sub>2</sub>, then this molecule could be destroyed by radiation to form O<sub>2</sub> (Loeffler et al., 2006a). Alternatively, if O<sub>2</sub> is produced first (Johnson et al., 2005; Petrik et al., 2006), then subsequent H-addition reactions (Eqs. (3a) and (3b)) could produce H<sub>2</sub>O<sub>2</sub>. Recent work (Petrik et al., 2006) has suggested that in 87 eV electron irradiation of water-ice, H<sub>2</sub>O<sub>2</sub> forms first via Eq. (4), before Eq. (5) produces HO<sub>2</sub>. The latter then undergoes an electronic excitation to give O<sub>2</sub>. However, Petrik et al. (2006) also proposed that HO<sub>2</sub> is a stable precursor for O<sub>2</sub> at temperatures as high as 130 K. In fact, the O<sub>2</sub> yield in their experiments, as in earlier studies (Johnson et al., 2005; Orlando and Sieger, 2003), increased with temperature. We do not see HO<sub>2</sub> being stable above 100 K, as the amount of HO<sub>2</sub> present after saturation doses (>1 eV/16 amu) decreases with increasing temperature. The greater penetrating range of MeV protons and much higher dose used in the present work compared with that of Petrik et al. (2006) may contribute to this difference. Further comprehensive and systematic investigations of O<sub>2</sub> formation are needed to sort out the influence of different temperature and radiation environments.

As an alternative to the above, H<sub>2</sub>O<sub>2</sub> and O<sub>2</sub> may form independently of one another directly from H<sub>2</sub>O and then subsequently interconvert. Further work needs to be carried out to determine this. What is clear though is that H<sub>2</sub>O<sub>2</sub> and O<sub>2</sub> are likely to be present together on icy satellite surfaces because of their intimately related radiation chemistry in water-ice.

In addition, O<sub>3</sub> may also be spatially correlated with H<sub>2</sub>O<sub>2</sub> and O<sub>2</sub> on icy satellites since O<sub>3</sub> is formed easily in O<sub>2</sub> clusters. If O<sub>3</sub> forms in O<sub>2</sub>-rich regions, then the  $\nu_3$  band maximum will be at  $\sim 1038\text{--}1040\text{ cm}^{-1}$ , but if the peak is at  $1035\text{--}1036\text{ cm}^{-1}$ , then this would indicate that the O<sub>3</sub> is dispersed in the water-ice. The O<sub>3</sub> molecule would in effect be acting as a tracer for O<sub>2</sub> and would provide an indication of its state.

HO<sub>2</sub> and HO<sub>3</sub> radicals may also be present on the Galilean satellites, and might be expected to be found in greater column densities at the colder polar regions. The band areas of these species decrease with increasing temperature and by  $\sim 100$  K, we were unable to detect their infrared bands in our laboratory samples. Interestingly, the HO<sub>2</sub>·H<sub>2</sub>O complex is calculated (Aloisio et al., 1999) to have a UV absorption at 255 nm, compared with 192 nm for the free HO<sub>2</sub> radical. HO<sub>2</sub> in water-ice may contribute to the 'ozone-like absorber' band observed near 260 nm by the Galileo ultraviolet spectrometer (UVS) instrument (Hendrix et al., 1999).

## 5. Conclusions

The major conclusions from this work are summarized below.

- 1) The position of the  $\nu_3$  band of O<sub>3</sub> is sensitive to the environment in which the molecule is situated. When an O<sub>3</sub> is surrounded by O<sub>2</sub> molecules the spectral maximum occurs at  $1038.5\text{ cm}^{-1}$ . The peak shifts to  $1035.6\text{ cm}^{-1}$  when O<sub>3</sub> is surrounded by water molecules.
- 2) HO<sub>2</sub> and HO<sub>3</sub> are produced in O<sub>2</sub>-containing ices and are present at a detectable level up to  $\sim 100$  K.
- 3) The H<sub>2</sub>O<sub>2</sub> growth observed on warming irradiated H<sub>2</sub>O + O<sub>2</sub> ices suggests that some OH radicals produced below 80 K are trapped and do not react until the 80–120 K temperature range.
- 4) For O<sub>3</sub>, HO<sub>2</sub>, and HO<sub>3</sub>, the trend upon warming of the irradiated H<sub>2</sub>O + O<sub>2</sub> ices is for these species to decrease in abundance. Some O<sub>3</sub> is lost via desorption from the ice, however a loss via chemical destruction appears to occur. Mobile OH radicals are the most likely source for the destruction of these species.
- 5) A variety of oxidizing species, including H<sub>2</sub>O<sub>2</sub>, O<sub>3</sub>, OH, HO<sub>2</sub>, and HO<sub>3</sub> are all produced in H<sub>2</sub>O + O<sub>2</sub> ices and likely part of the oxidant inventory on icy satellites such as Europa and Enceladus. These species may provide chemical energy to sub-surface water that may potentially harbor non-Earth originating life.

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