

Correlation of Pc5 wave power inside and outside the magnetosphere during high speed streams

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Abstract. We show a clear correlation between the ULF wave power (Pc5 range) inside and outside the Earth's magnetosphere during high speed streams in 1995. We trace fluctuations beginning 200 R_E upstream using Wind data, to fluctuations just upstream from Earth's bow shock and in the magnetosheath using Geotail data and compare to pulsations on the ground at the Kilpisjarvi ground station. With our 5-month data set we draw the following conclusions. ULF fluctuations in the Pc5 range are found in high speed streams; they are non-Alfvénic at the leading edge and Alfvénic in the central region. Compressional and Alfvénic fluctuations are modulated at the bow shock, some features of the waveforms are preserved in the magnetosheath, but overall turbulence and wave power is enhanced by about a factor of 10. Parallel (compressional) and perpendicular (transverse) power are at comparable levels in the solar wind and magnetosheath, both in the compression region and in the central region of high speed streams. Both the total parallel and perpendicular Pc5 power in the solar wind (and to a lesser extent in the magnetosheath) correlate well with the total Pc5 power of the ground-based H -component magnetic field. ULF fluctuations in the magnetosheath during high speed streams are common at frequencies from 1–4 mHz and can coincide with the cavity eigenfrequencies of 1.3, 1.9, 2.6, and 3.4 mHz, though other discrete frequencies are also often seen.

Key words. Interplanetary physics (MHD waves and turbulence) – Magnetospheric physics (solar wind-magnetosphere interactions; MHD waves and instabilities)

1 Introduction

The possible role of Pc5 pulsations in energizing electrons to relativistic levels in the magnetosphere (e.g. Rostoker et al., 1998; Baker et al., 1998; Mathie and Mann, 2000) has increased the importance of understanding their excitation mechanisms. Pc5 pulsations, both in the equatorial outer

magnetosphere (e.g. Kokubun et al., 1989; Anderson et al., 1991) and on the ground (e.g. Greenstadt et al., 1979; Engebretson et al., 1991; Glassmier, 1995), have been linked to high speed solar wind streams in 1995. It is fairly well established that bow shock associated ULF fluctuations and pressure pulses can convect downstream, cause magnetosheath fluctuations and may be responsible for magnetospheric pulsations, particularly in the Pc3–4 band (e.g. Greenstadt et al., 1983; Sibeck et al., 1989; Fairfield et al., 1990; Engebretson et al., 1991; Lin et al., 1991). However, gaps remain in understanding the energy transfer process from high speed streams in the solar wind into the magnetosphere, particularly in the Pc5 frequency range. Below, we review the relevant knowledge of ULF waves and fluctuations (particularly Pc5 range) in high speed streams, in the magnetosheath and in the magnetosphere. We also review theory and observations of transmission across the bow shock and magnetopause, and the current understanding of cavity and waveguide modes.

1.1 Waves and fluctuations in the solar wind

High speed solar wind streams occur during the declining phase of the solar cycle, originating from the Sun's polar coronal holes. The polar coronal holes producing the fast solar wind reach their maximum latitudinal extent near sunspot minimum, confining the slow solar wind to a narrow equatorial belt. Co-rotating Interaction Regions (CIRs) often occur at the leading edges of high speed streams as the high speed streams overtake the slower solar wind (Belcher and Davis, 1971; Tsurutani et al., 1995). Inward and outward propagating waves are found in these CIR regions and may be the result of turbulence driven by velocity shear (Coleman, 1968). The large-amplitude fluctuations found in the compression regions or colliding stream regions near the leading edge of streams cannot propagate away from these regions (Burlaga, 1970) and thus must persist within these regions for a considerable distance. In the central region of high speed streams, only outwardly propagating Alfvén waves are found. Belcher and Davis (1971) suggested that these waves

may be generated in the solar atmosphere by turbulent processes. Theoretically, slow and fast MHD waves usually are quickly damped in collisionless plasmas with moderate to high plasma β (Barnes, 1966), leaving only outward propagating Alfvén waves. These persist and can be quite pure in the inner solar system. Belcher and Davis showed the Alfvén waves to be broad-band Pc5 range and below, with peaks at 1.67, 1.39, 1.11, 0.83, 0.37 mHz; though using Mariner 5 data they were limited to a Nyquist frequency of 1.67 mHz and could see nothing higher. Waves and fluctuations in high speed streams will impact the magnetosphere first at the bow shock.

1.1.1 Waves and fluctuations across boundaries

McKenzie and Westphal (1969, 1970) calculated amplitudes and directions of waves diverging from a fast hydromagnetic shock perturbed by a small amplitude hydromagnetic wave, using the Rankine-Hugoniot conservation equations across the shock, and Snell's law. They found that fast magnetoacoustic longitudinal waves are greatly amplified on passage through the shock and that Alfvénic waves are only moderately amplified. They suggested that this amplification might contribute to the turbulent nature of the Earth's magnetosheath and that fluctuations in the magnetosheath would tend to be predominantly longitudinal rather than Alfvénic in nature. More recent MHD simulation studies indicate that Alfvén/slow-mode waves and other discontinuities can be generated in the magnetosheath by the interaction between the bow shock and various MHD discontinuities or Alfvén waves in the upstream solar wind (Lin et al., 1996). Observationally, Sibeck et al. (1997), using simultaneous Wind and Geotail data, found clear evidence for Alfvénic fluctuations propagating into the magnetosheath. The magnetosheath now is well known as a turbulent region, with most fluctuations originating in the solar wind or associated with the quasi-parallel bow shock (e.g. Crooker et al., 1981; Sibeck et al., 2000). Kwok and Lee (1984) calculated transmitted and reflected MHD waves at the magnetopause when it is a rotational discontinuity. They found that the reflected fast magnetosonic wave is significantly amplified and thus may contribute to magnetosheath turbulence.

Satellite observations have shown that the magnetopause can stably exist as either a tangential or rotational discontinuity (e.g. Russell and Elphic, 1978; Sonnerup et al., 1981). Under conditions when the magnetopause is a tangential discontinuity, $B_n = 0$, the magnetopause is closed and MHD wave transmission is not efficient (e.g. McKenzie, 1970; Wolfe and Kaufmann, 1975). When the magnetopause is a rotational discontinuity, $B_n \neq 0$, Kwok and Lee (1984) determined that MHD wave transmissions occurred over a wide range of incident angle and that the transmitted waves were usually amplified. They suggested that MHD wave transmission at an open magnetopause can be a significant mechanism for energy transport from the magnetosheath to the magnetosphere. Furthermore, Kwok and Lee (1984) suggested that the valid range of wavelength,

λ , and wave frequency, f , for transmitted waves would be $\lambda \gg 1000$ km (larger than the thickness of the magnetopause) and $f \ll 150$ mHz based on $v_A = 150$ km/s at the magnetopause. This designation favors ULF waves, particularly Pc5 and lower frequency waves.

1.1.2 Waves in the magnetosphere

Dungey (1954) originally proposed that the regular periods of geomagnetic pulsations might be due to standing Alfvén waves excited on geomagnetic field lines, later termed field line resonances (FLRs). When waves are observed in the magnetosphere, the peak amplitude occurs at the latitude that coincides with the field line that resonates at the wave frequency. Pc5 pulsations are common at latitudes between about 60° and 70° (Hughes, 1994 and sources therein). In space these map from about the plasmopause to near the magnetopause on the dayside. Field line resonances are observed in the Pc5 range, with discrete frequencies at approximately 1.3, 1.9, 2.6, and 3.4 mHz (Ruohoneimi et al., 1991; Samson et al., 1992; Samson and Rankin, 1994). Samson et al. (1992) showed that these discrete frequencies were compatible with MHD waveguide and cavity modes in the magnetosphere. The Earth's magnetospheric cavity extends from an outer boundary, possibly the magnetopause (or bow shock), to an internal turning point, possibly near the plasmopause. In the cavity model energy is input into the magnetosphere and the cavity as a whole rings at its own eigenfrequencies, efficiently transporting energy at those frequencies to field lines in the magnetosphere and producing the classic field line resonance signature. The waveguide model is similar except that the cavity remains open downtail; waveguide modes propagate antisunward at the natural frequencies of the magnetosphere (e.g. Harrold and Samson, 1992; Mann et al., 1999).

It is commonly agreed that toroidal Pc5 pulsations are caused by an external source in the solar wind. The most frequently cited source of Pc5 pulsations in the magnetosphere is the Kelvin-Helmholtz instability at the magnetopause (e.g. Dungey, 1955; Miura, 1992; Anderson, 1994; Engebretson et al., 1998). Mann et al. (1999) have recently shown that for very large flow speeds at the magnetopause flanks, the Kelvin-Helmholtz instability can energize body type waveguide modes. Other possible sources of Pc5 pulsations have been proposed, such as upstream shock-related pressure oscillations that drive magnetopause surface waves with periods in the Pc5 range (Sibeck et al., 1989). Fairfield et al. (1990) suggested that upstream pressure variations may be linked to magnetospheric compressions. Engebretson et al. (1998 and sources therein) suggested that if the compression regions at the leading edges of high speed streams contain waves in the Pc5 range, they could provide a source of wave energy to the magnetosphere or that the waves could act as seed perturbations to drive boundary displacements that are amplified by the Kelvin-Helmholtz instability. By contrast, Kepko et al. (2002) showed observations of pressure fluctuations at the same discrete frequencies in

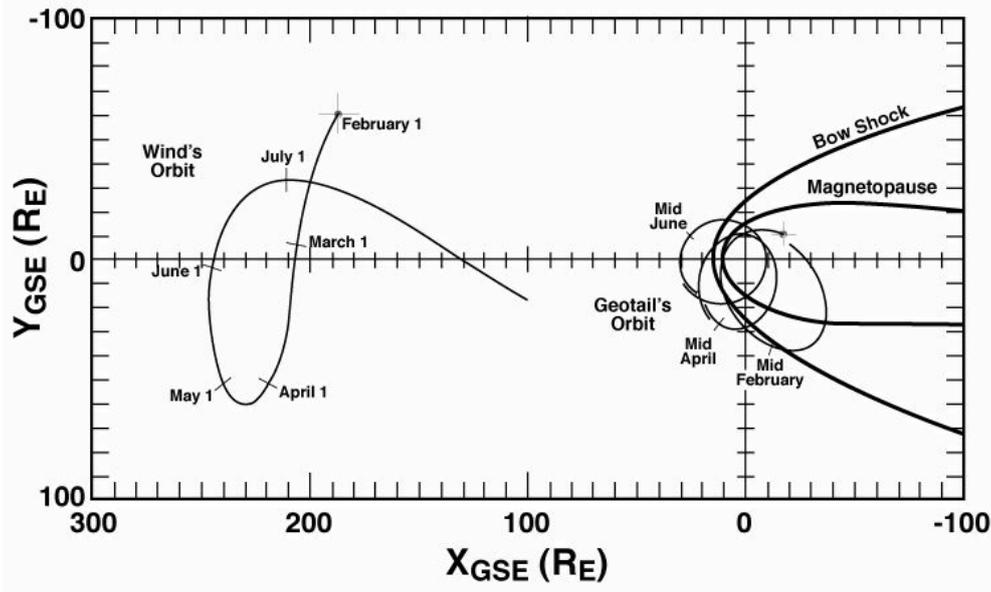


Fig. 1. Wind orbit from February through July 1995 and example Geotail orbits in mid February 1995, mid April 1995 and mid June 1995.

the solar wind as in the magnetosphere and suggested that the solar wind may be a direct source for discrete Pc5 pulsations. Wright and Rickard (1995) showed that broad-band fluctuation power in the solar wind can lead to enhanced excitation of the magnetospheric cavity or waveguide modes, even if the spectral content of the upstream and magnetosphere waves are different.

In this paper we investigate the linkage of waves and fluctuations in the solar wind, magnetosheath, and on the ground. We show a clear correlation between Pc5 wave power inside and outside Earth's magnetosphere during high speed streams in 1995. We trace the fluctuations beginning 200 R_E upstream using Wind data, to fluctuations just upstream from the Earth's bow shock and in the magnetosheath using Geotail data and compare to waves on the ground at the Kilpisjarvi ground station. We show observations of waveforms, power spectral densities and total power in all regions and discuss probable transfer mechanisms.

2 Instruments and methods

For this study we use data from the Wind and Geotail satellites and from the Kilpisjarvi ground-based station. For the time period of this study, February through June 1995, Wind is located far upstream ($\sim 200 R_E$) with y_{GSE} varying between ± 60 , as shown in Fig. 1. Geotail is orbiting the Earth, just transitioning to an orbit with an apogee of $\sim 30 R_E$. We show three orbits during the 5-month time frame in Fig. 1. The time interval of this study is dictated by the time when Geotail began its near-Earth orbit to when high speed streams became less regular.

We use data from the Magnetic Field Investigation (MFI) (Lepping et al., 1995) and the Solar Wind Experiment (SWE)

(Ogilvie et al., 1995) on Wind. We use magnetometer data from the Magnetic Field Measurement (MGF) (Kokubun et al., 1994) on Geotail. We use Kilpisjarvi (KIL) ground-based magnetometer data from the IMAGE magnetometer array (Luhr et al., 1998; www.geo.fmi.fi/image). KIL is located at 69.02° geographic latitude, and is thus in a good position to measure Pc5 pulsations. For the space-based magnetometers we use 1-min resolution data to provide detail on the Pc5 portion of the spectrum (0–8.3 mHz) and 3-s resolution data to generate a broader ULF wave spectrum from 0–166 mHz. At Kilpisjarvi we have used 10-s resolution data (0–50 mHz) from the morning sector, summing only those 2-h bins which lie entirely within the range 05:00–11:00 LT.

For the satellite magnetometers we used the following analysis procedure. We first put the data into field-aligned coordinates and then calculated the power spectra separately for the parallel component and for the vector average of the perpendicular component. For the one-min data we used a sliding 128 point FFT (2 h) window and summed power in all the bands (1–8.3 mHz). For the 3-s data we used a sliding 512 point FFT (1/2 h) window and summed power in the 1–10 mHz band. Each data set was detrended over the sliding point window (2 or 1/2 h) by subtracting an average value and then tapered using a Parzen window. The data were processed using a fast Fourier transform algorithm to create power spectral densities (PSD) and then the PSDs were summed to obtain total power. The window was advanced for each data point to create a running spectrum and in addition, an average value was determined for every 6- and 24-h worth of data.

For the ground-based magnetometer we used a sliding 720 point FFT window to the H component, advanced by successive half hour intervals through the day. A daily Pc5 power value was calculated by summing power in the 1–10 mHz

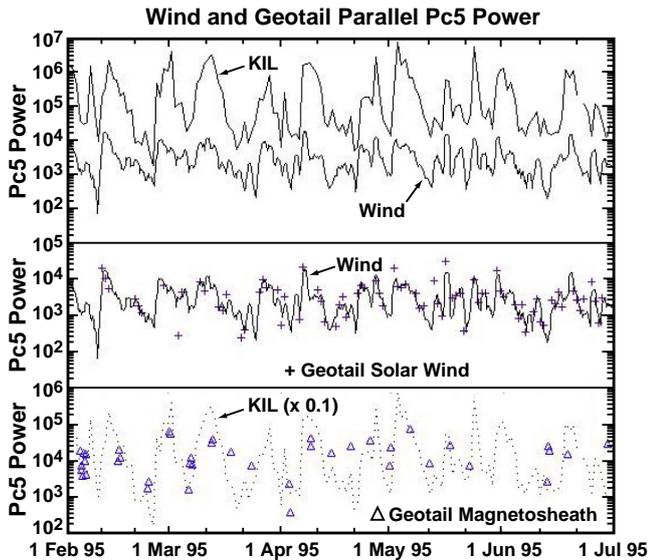


Fig. 2. Top panel: Pc5 pulsation power at Kilpisjarvi ground-based station (upper curve) compared to parallel Pc5 power at Wind (lower curve). Middle panel: parallel Pc5 power at Wind (curve) and Geotail (crosses) with both in the solar wind. Bottom panel: Pc5 pulsation power at Kilpisjarvi (dotted curve) and parallel Pc5 power at Geotail in the magnetosheath (triangles).

band from an average of the windows centered in the local time morning sector. These data are nearly identical to those published in Mathie and Mann (2000), except that in the latter all bins in the morning sector were summed, not averaged.

3 Observations

Figure 2 compares Pc5 wave power inside and outside of the magnetosphere during 5 months in 1995. The peaks in power coincide with the peaks in solar wind speed in the high speed streams, as shown previously by Mathie and Mann (2000). The top panel of Fig. 2 shows ULF wave power in the Pc5 range from the H component at Kilpisjarvi (upper curve) and the parallel component of Pc5 power from Wind (lower curve). These two curves represent the extreme locations examined in this study, i.e. far upstream and on the ground. The correlation between these two curves is good, as can be seen by eye; the correlation coefficient is 0.61. The middle panel shows the same parallel Pc5 power from Wind to guide the eye along with the sparse points of parallel Pc5 power from Geotail in the solar wind (crosses). The Geotail data in the solar wind are not continuous because Geotail's orbit also takes it into the magnetosheath and magnetosphere. These two sets of data in the middle panel are nearly coincident, but in some cases, the power at Geotail is slightly higher. In these cases Geotail is connected to a quasi-parallel bow shock for an extended time and the Pc5 power is enhanced. The bottom panel shows the Pc5 wave power from Geotail in the dusk magnetosheath (triangles), along with the H -component Pc5 power from KIL shifted down by an order

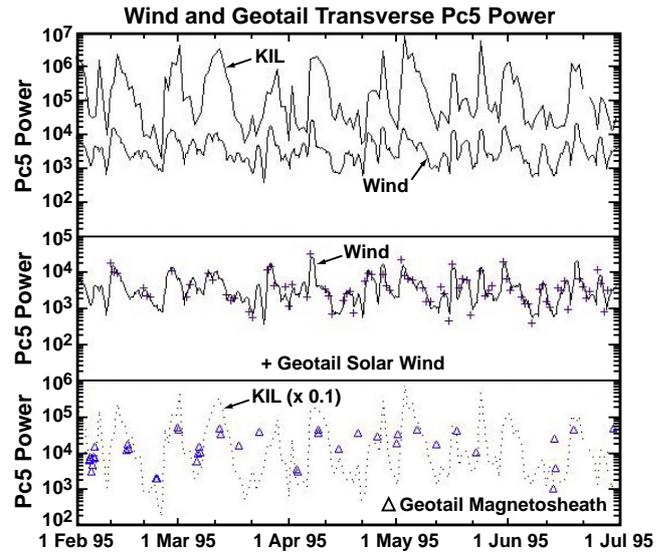


Fig. 3. Top panel: Pc5 pulsation power at Kilpisjarvi ground-based station (upper curve) compared to perpendicular Pc5 power at Wind (lower curve). Middle panel: perpendicular Pc5 power at Wind (curve) and Geotail (crosses) with both in the solar wind. Bottom panel: Pc5 pulsation power at Kilpisjarvi (dotted curve) and perpendicular Pc5 power at Geotail in the magnetosheath (triangles).

of magnitude. The magnitude of the dusk magnetosheath Pc5 power lies between the solar wind and ground-based levels but follows a similar pattern, though the sparse points make the correlation difficult to see without the downshift in the KIL power curve to guide the eye. There were only 3 dawn-side magnetosheath points and each occurred inside a compression region, so it was not possible to establish any trend or see any correlation between the dawnside magnetosheath and KIL.

Figure 3 is similar to Fig. 2, but now we compare the H component at KIL with perpendicular power at Geotail and Wind. In comparing Fig. 2 and Fig. 3 we see that the parallel (compressional) power is at approximately the same level as the perpendicular (transverse) power in both the solar wind and magnetosheath. The correlation between the H component Pc5 power at Kilpisjarvi and perpendicular Pc5 power at Wind is also good, with a correlation coefficient of 0.58. The middle panel of Fig. 3 shows that the perpendicular power at Wind and Geotail in the solar wind correlates as well as it did in the parallel case. There are again cases when Geotail power is higher due to being connected to a quasi-parallel bow shock. This is more often true starting in April, when Geotail's orbit swings around to spend more time on the dawn side. The bottom panel of Fig. 3 is similar to the bottom panel of Fig. 2, i.e. transverse Pc5 wave power from Geotail in the dusk magnetosheath (triangles) along with H -component Pc5 power from KIL shifted down by an order of magnitude. The correlation here is as good as it was in the parallel case.

We look in detail at one classic high speed stream during

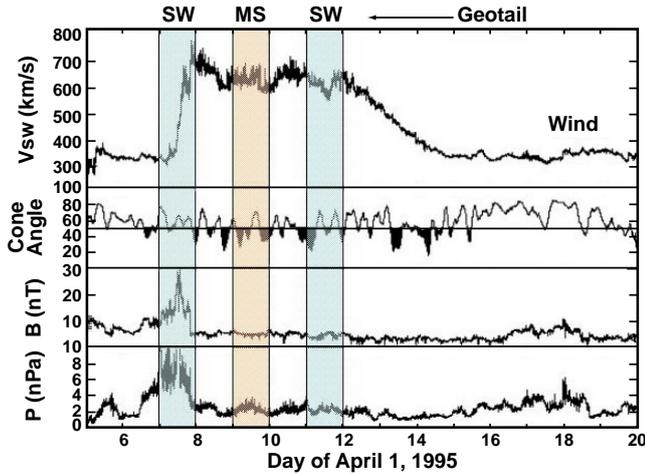


Fig. 4. Classic high speed stream, 5–20 April 1995. From top to bottom: solar wind high speed stream, cone angle, magnetic field magnitude, ram pressure.

the five month interval in 1995, 5–20 April. The top panel of Fig. 4 shows that the solar wind speed has an initial steep increase, with the speed remaining high for 4 days before gradually trailing off. At the leading edge of the stream there is a compression region, as seen by the increases in dynamic pressure and magnetic field magnitude. This compression regions occurs nearly a day before the increase in speed. This is typical of other streams in this study and agrees with the results of Engebretson et al. (1998), who showed that the peak pulsation power at two ground stations typically occurred about 1 day before the peak solar wind velocity. The high speed streams in 1995 follow this general pattern, but not all are as steep or as long lasting, and some have more than one compression region.

The second panel shows the IMF cone angle, the angle between the x -GSE axis and the magnetic field. During the high speed stream, the cone angle is less than 50° more than half the time, whereas outside the stream the cone angle is greater than 50° . This is also typical of high speed streams during this interval. During this high speed stream in April, Geotail’s orbit passed through the solar wind at the leading edge and then through the magnetosheath and solar wind in the central region (sub-intervals marked in Fig. 4). The type of waves seen are characteristic of each region of the high speed stream, and differ in the solar wind, magnetosheath, and on the ground, as we show in the following examples. Ideally, we would look at examples when Wind, Geotail, and KIL are all on the dawn side. However, with this data set we are limited to two of the three at any one time.

3.1 Leading edge of high speed stream

At the leading edge we show three examples. The first is a comparison of Wind and Geotail data (7 April 1995), when both are in the solar wind on the dusk side, with

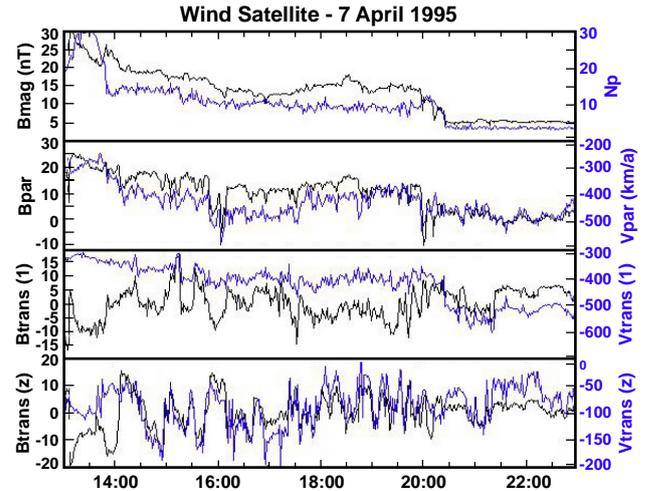


Fig. 5. Wind data in the solar wind inside the compression region, 7 April 1995. Top panel: magnetic field and density; second panel: $b_{||}$ and $v_{||}$, bottom two panels: transverse components of \mathbf{b} and \mathbf{v} .

Wind $\sim 200 R_E$ upstream and Geotail just in front of the bow shock. Then we compare Wind and Geotail data, when both are on the dawn side (13:00–20:00 UT, 19 June 1995), Wind is still $\sim 200 R_E$ upstream but Geotail is in the magnetosheath. We have no intervals when Geotail is in the dawn magnetosheath with KIL also on the dawn side. We show instead a comparison of Wind and KIL data when both are on the dawn side (02:00–10:00 UT 19 June 1995).

Our first interval corresponds to the first sub-interval in Fig. 4, when both Wind and Geotail are in the solar wind on the dusk side. We show perturbations in magnetic field and velocity in field-aligned coordinates during the compression region on 7 April 1995: one parallel and two transverse components in the bottom three panels of Fig. 5. In the top panel we show magnetic field magnitude and ion density. Both of these are high at the beginning of the plot, decrease just before 14:00 UT and again at about 20:30 UT. Because the magnetic field magnitude and the density are changing significantly up until 20:30 UT, we expect the fluctuations not to be Alfvénic in this region. Aperiodic Alfvén waves would show a correlated variation between the components, as defined by the relation, $\mathbf{b} = \pm(4\pi\rho)^{1/2}\mathbf{v}$ (Belcher and Davis, 1971). This is not the case, as can be seen in the bottom 3 panels of Fig. 5. After the drop in magnetic field magnitude and ion density at about 20:30 UT the values remain fairly steady for the remainder of the plot.

Figure 6 follows the same format as Fig. 5 but has data from Geotail in the solar wind. There are many similarities between Figs. 5 and 6. The top panel shows magnetic field magnitude and density drops about 45 min earlier in Fig. 5 compared to Fig. 6 which is expected because of the difference in location. If the fluctuations are trapped in the compression region and the compression region encompasses a large area we might expect to see similar fluctuations at Wind and at Geotail even with almost $200 R_E$ separation between

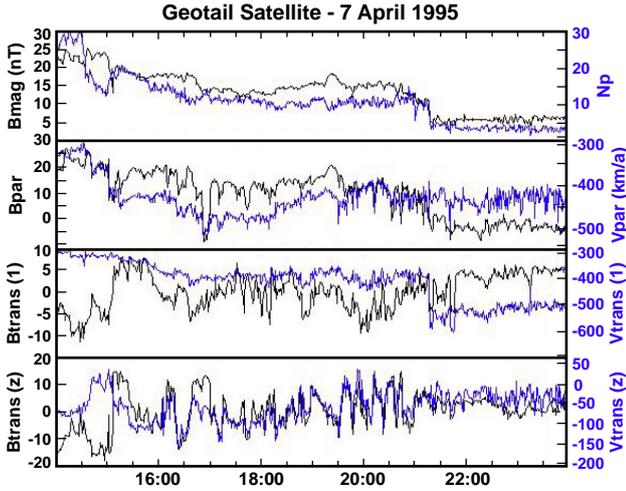


Fig. 6. Geotail data in the solar wind inside the compression region, 7 April 1995. Top panel: magnetic field and density; second panel: $b_{||}$ and $v_{||}$, bottom two panels: transverse components of \mathbf{b} and \mathbf{v} .

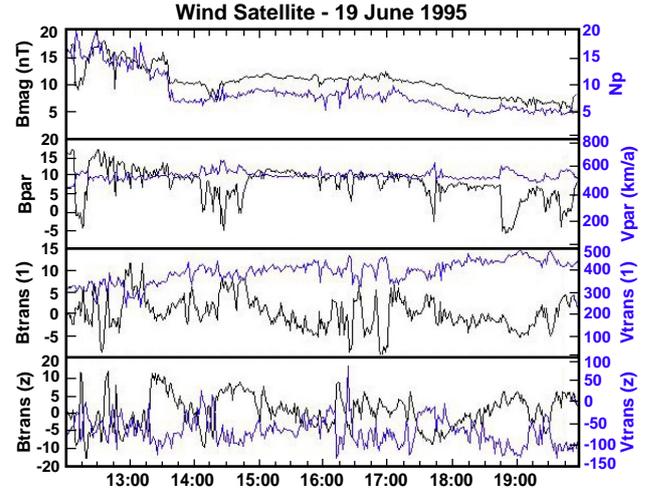


Fig. 8. Wind data in the solar wind inside the compression region, 19 June 1995. Top panel: magnetic field and density; bottom three panels: field-aligned and perpendicular components of \mathbf{b} and \mathbf{v} .

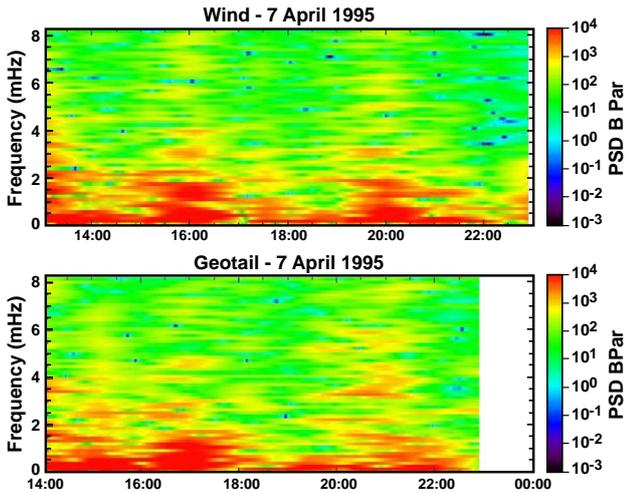


Fig. 7. PSD for Wind (top panel) and Geotail (bottom panel) on 7 April 1995.

them in x and $20\text{--}30 R_E$ separation in y . There are similarities between the fluctuations but they are not identical. The level of fluctuations of the components are similar and it is possible to pick out some similar peaks and dips in the components. For example in the second panels there are similar dips in B_{par} at $\sim 16:00$ UT in Fig. 5 and $\sim 17:00$ UT in Fig. 6. Additionally, the v_{trans} component in the third panels begins low and the trend is to increase slightly in magnitude and then become steady with a sharp increase in magnitude coinciding with a drop in B magnitude and density $\sim 20:30$ UT in Fig. 5 and $\sim 21:20$ UT in Fig. 6. Although they do not match up in all details there is enough similarity to believe that they are portions of one extended region.

Power Spectral Densities (PSDs) in the Pc5 range associated with the parallel (compressional) component are shown

for Wind (top panel) and Geotail (bottom panel) on 7 April 1995 in Fig. 7. Both satellites are in the solar wind and the color scales are identical for both panels in this figure. The two panels are not identical, but like the waveforms there are many similarities. In both panels there are three large PSD enhancements extending over much of the Pc5 range. At Wind these enhancements are centered at about 13:00 UT, 16:00 UT, and 20:00 UT, while at Geotail they occur approximately an hour later in each case. The peak power is primarily in the 0–4 mHz range, though at Geotail there are a few peaks up to 6 mHz. The discrete power peaks are not identical at the two satellites, but they are similar. For example, the enhancement centered at 16:00 UT at Wind and 17:00 UT at Geotail has broad band power under 1.7 mHz in each case with discrete peaks at 2.3, 3.0, and 3.8 mHz for Wind and 2.5, 2.9, 4.0 and 4.6 mHz for Geotail.

We chose another interval in which both Wind and Geotail were on the dawn side with Wind still $\sim 200 R_E$ upstream but with Geotail in the magnetosheath. Figures 8 and 9 again show \mathbf{b} and \mathbf{v} fluctuations in field-aligned coordinates, as in Figs. 5 and 6. Figure 8 shows the magnetic field and plasma fluctuations at Wind. There is a drop in magnetic field magnitude and density at $\sim 13:40$ UT, shown in the top panel, followed by some smaller fluctuations in each. After 13:40 UT the pressure also decreases (not shown) but is still higher than that typically found in the central region. Before 13:40 UT Wind is clearly in the compression region and the fluctuations are non-Alfvénic, as can be seen in the bottom three panels of Fig. 8. After 13:40 UT the nature of the waves changes; they exhibit some anticorrelation between the magnetic field and velocity components that could be an indication of shear Alfvén waves or fluctuations (Kivelson and Russell, 1995) and a transition to the central region of the stream. However, the fluctuations in the magnetic field and velocity components are somewhat irregular during the intervals after

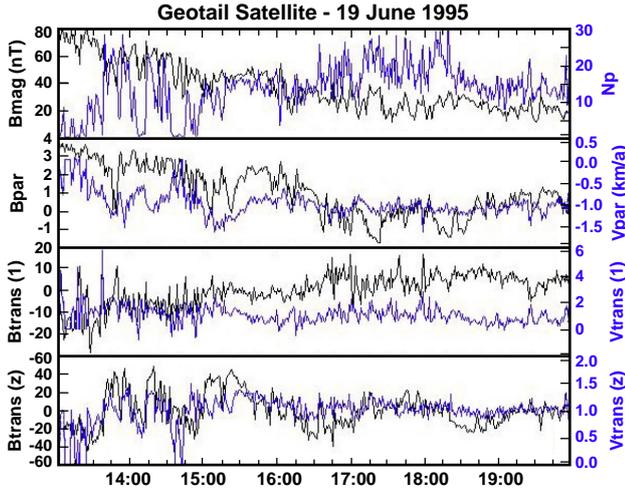


Fig. 9. Geotail data in the magnetosheath inside the compression region, 19 June 1995. Top panel: magnetic field and density; bottom three panels: field-aligned and perpendicular components of b and v .

13:40 UT, when fluctuations in density and magnetic field magnitude increase, e.g. between 15:45 UT and 17:15 UT. The magnetosheath fluctuations for magnetic field and velocity at Geotail (Fig. 9) are enhanced over those seen at Wind, especially in the compression region up to about 15:00 UT. The fluctuations in density remain high throughout the interval and the amplitude is especially high during the interval 16:30 UT to 18:20 UT. The anticorrelations are not evident at Geotail. We previously showed larger ULF wave fluctuations in the magnetosheath in Figs. 2 and 3, so this result is not surprising. We calculated a time lag between the two satellites based on x separation and a solar wind speed of about 40 min. Starting at about 14:05 UT in Fig. 8 and at about 14:40 UT in Fig. 9, we can pick out a similar feature in the B -parallel component, a multiple dip, at both Wind and Geotail. There are other similarities, but there are more differences, most likely due to modulation at the bow shock or reflection at the magnetopause.

Figure 10 shows the power spectral density (PSD) associated with the parallel (compressional) component at Wind (top panel) and Geotail (bottom panel). The color scales are different in each panel, in order to pull out the peaks in the spectrum. The total power increases as we move from the top panel to the bottom panel. We first compare general trends in the solar wind (top panel) with general trends in the magnetosheath (bottom panel). There are five intervals of PSD increase at Wind (top panel): at the beginning of the interval until about 13:15 UT; from about 13:45 UT until about 15:15 UT; centered at about 16:15 UT, from about 17:15 UT until about 18:15 UT; and centered at about 19:00 UT. These intervals are not exactly mirrored at Geotail (bottom panel) but there are some similarities. At Geotail we could define two intervals of PSD increase, starting at 13:00 UT until about 16:00 UT, that could be related to the first two

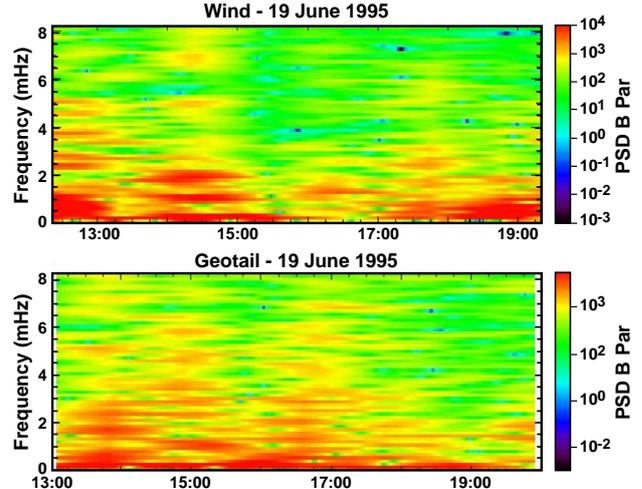


Fig. 10. PSD for Wind (top panel) and Geotail (bottom panel) on 19 June 1995.

intervals at Wind, though in the magnetosheath the intervals have nearly merged and have similar frequency peaks below 4 mHz. The third interval at Geotail has peaks throughout the Pc5 frequency range in contrast to the third interval at Wind which has power only below 2 mHz. The fourth interval at Wind has no clear corresponding power at Geotail, and the fifth interval is small, both at Wind and Geotail.

Throughout the region there are peaks in Pc5 wave power that fall primarily between 1 and 5 mHz in the solar wind (Fig. 10 top panel) and between 1 and 7 mHz in the magnetosheath (bottom panel). We can compare the power associated with the waveform features starting at Wind at 14:05 UT and at Geotail at 14:40 UT (Figs. 8 and 9). We see some similar peaks in the PSD at these times in Fig. 10. Wind (top panel) has 2 dominant broad-band peaks centered 0.9–1.2 mHz and 1.7–2.2 mHz, and other minor peaks at 2.6 mHz, 3.4 mHz, 5 mHz, and 7 mHz. Geotail (bottom panel) has a broad-band peak near 1 mHz, and other peaks at 2.4, 3.0 and 3.6, 5, and 6 mHz. The frequencies associated with the other intervals do not compare as well between the solar wind and magnetosheath. In summary, there are some similar general trends in PSD in the solar wind and magnetosheath at the leading edge of high speed streams, and, as noted above, there is an interval in this example where there are peaks in PSD at the same frequency for Wind and Geotail. But there are also cases where there is power in the solar wind with little power in the magnetosheath and vice versa.

The third interval is earlier in the day on 19 June 1995 and is a comparison between Wind $\sim 200 R_E$ upstream and KIL, ground-based station. There are no similarities between the waveforms at Wind and KIL which is not surprising since Wind measures the solar wind magnetic field while KIL measures the Earth's intrinsic field, as well as influences from magnetospheric and ionospheric currents. We show them for reference in Figs. 11 and 12. Figure 11 follows the same format as previous waveform plots shown here, that is, the mag-

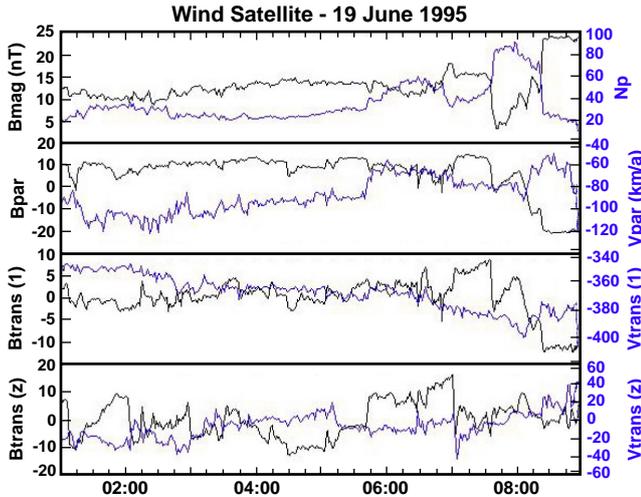


Fig. 11. Wind data in the solar wind inside the compression region, 19 June 1995. Top panel: magnetic field and density; bottom three panels: field-aligned and perpendicular components of \mathbf{b} and \mathbf{v} .

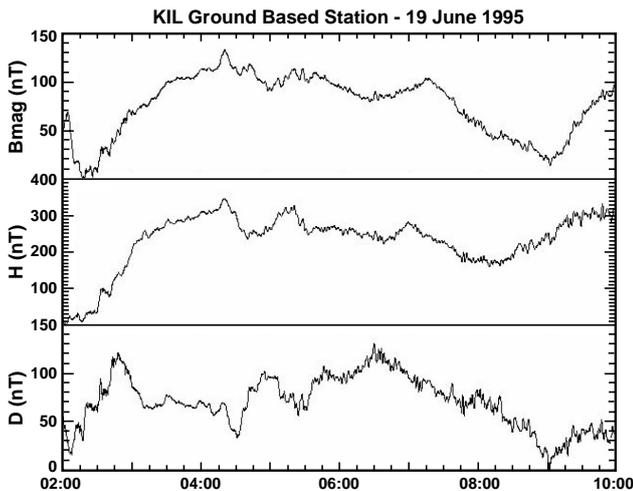


Fig. 12. KIL ground-based data during the time when the compression region of a high speed stream impacts the Earth's magnetosphere on 19 June 1995. Top panel: magnetic field magnitude; middle panel: H component; bottom panel: D component.

nitude of the magnetic field and density in the top panel, and the field-aligned and perpendicular components of magnetic field and velocity fluctuations in the bottom three panels. The fluctuations at Wind are again non-Alfvénic, as we have seen consistently for the leading edge of high speed streams. We show the ground-based magnetometer data at KIL (Fig. 12) as variometer data, that is, as a range rather than absolute values. We show the magnitude and H and D components in top to bottom panels, and do not have velocity or density data. There is no correlation between the Wind and KIL large-scale waveforms, so we next look at the associated ultra low frequency wave power (Pc5 range).

Figure 13 shows the power spectral density (PSD) associ-

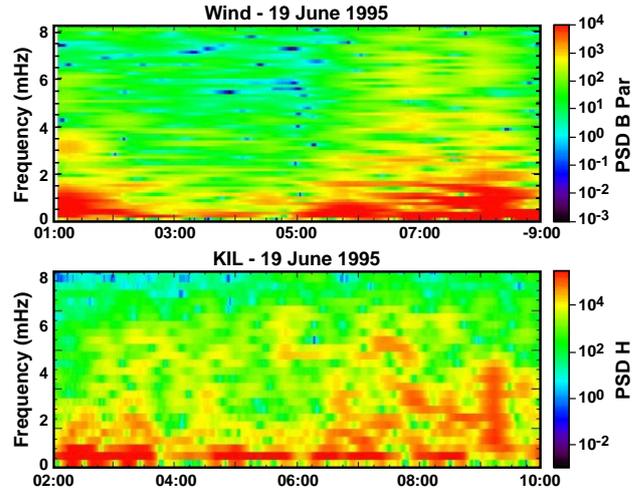


Fig. 13. PSD for Wind (top panel) and KIL (bottom panel) on 19 June 1995.

ated with the parallel (compressional) component at Wind (top panel) and the H component at KIL (bottom panel). These are the same components as were used in Fig. 2 (top panel). The color scales are different in each panel in Fig. 13, in order to pull out the peaks in the spectrum. The KIL power bands are somewhat wider than those at Wind due to the resolution of the data. There are general trends in the two panels that are similar: an enhancement in power between 0 and 4 mHz at the beginning of the plot; a depletion in PSD between 03:00 UT and 05:00 UT (top) and between 04:00 UT and 06:00 UT (bottom); and increased PSD between 06:00 UT and 09:00 UT (top) and between 07:00 UT and 10:00 UT (bottom). The peaks in PSD occur in the frequency range of 0–4 mHz for Wind, but extend to a higher frequency for KIL \sim 06:00 UT, \sim 07:30 UT, and \sim 09:15 UT. Although there is similarity in the general trends, there is no one-to-one match up between individual frequency peaks at Wind with those at KIL. For example, at \sim 09:15 UT at KIL there is broad-band power from 1–5 mHz, while at Wind at \sim 08:15 UT the highest power (also broad-band) is from 0–1.5 mHz with a lower broad-band power from 1.6 to 2.1 mHz, and more discrete peaks at 2.6–2.7, 3.4, and 3.6 mHz. It is interesting that the cavity eigenfrequencies of 1.3, 1.9, 2.6, and 3.4 are all evident in the solar wind data during the largest broad-band increase at KIL \sim 09:15 UT.

3.2 Central region of high speed stream

In the central region we show two examples. First, we show a comparison of Wind and Geotail data (9 April 1995) when both are on the dusk side, with Wind $\sim 200 R_E$ upstream and Geotail in the magnetosheath. There are no intervals in the central regions of high speed streams with Geotail on the dawn side, but we show a comparison of Wind and KIL data when both are on the dawn side (2 March 1995).

For the first example both Wind and Geotail are on the dusk side in Figs. 14 and 15, respectively. This is the second

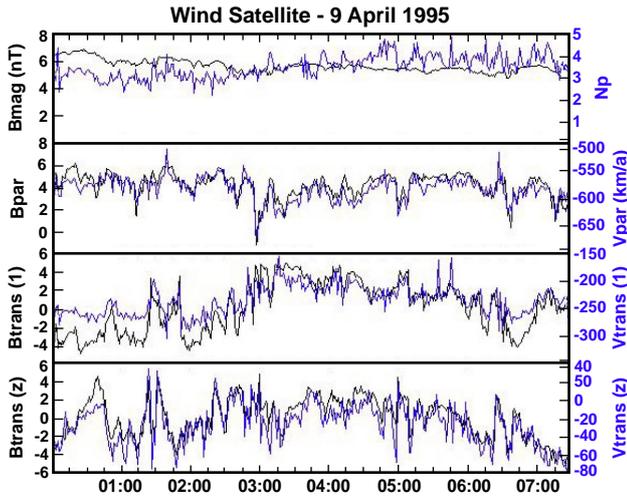


Fig. 14. Wind data in the solar wind in the central region of the high speed stream on 9 April 1995. Top panel: magnetic field and density; bottom three panels: field-aligned and perpendicular components of \mathbf{b} and \mathbf{v} .

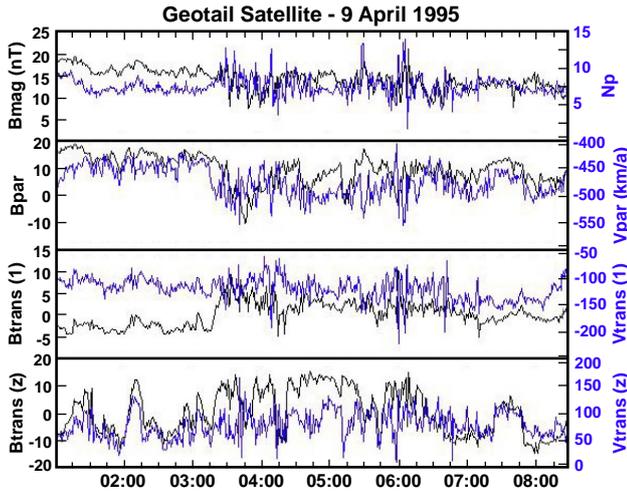


Fig. 15. Geotail magnetosheath data in the central region of the high speed stream on 9 April 1995. Top panel: magnetic field and density; bottom three panels: field-aligned and perpendicular components of \mathbf{b} and \mathbf{v} .

sub-interval shown in Fig. 4 and it occurs on 9 April 1995. Figures 14 and 15 follow the same format as previous waveform plots shown here, that is, the magnitude of the magnetic field and density in the top panel and the field-aligned and perpendicular components of magnetic field and velocity fluctuations in the bottom three panels. In Fig. 14, the bottom three panels show the Alfvénic nature of the fluctuations at the Wind satellite in the solar wind. The magnetic field magnitude and density in the top panel are not as highly variable as they were in the compression region, although there is some evidence of a compressional wave in the magnetic field magnitude. Each of the components of magnetic field and velocity are correlated (bottom three panels of Fig. 14)

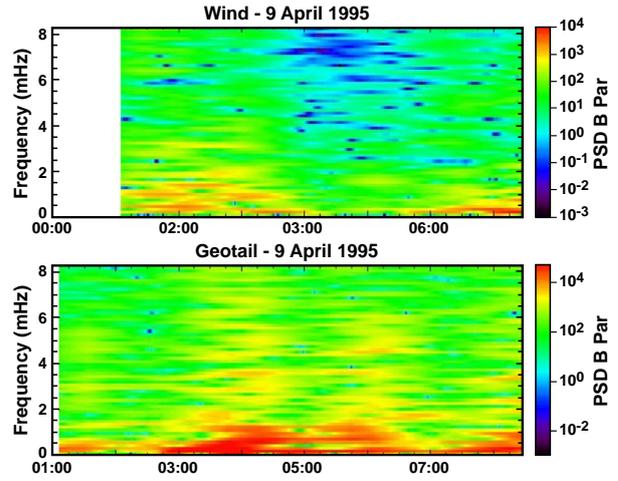


Fig. 16. PSD for Wind (top panel) and Geotail (bottom panel) on 9 April 1995.

but the correlation is not perfect.

In the magnetosheath in Fig. 15 the amplitude of the fluctuations is increased and the fluctuations are less Alfvénic. Based on timing considerations these fluctuations should be from the same region of the high speed stream. It is also possible to spot some similarities such as the higher level of B_{par} at the beginning of each plot (second panels), followed by a drop at about 03:00 UT in Fig. 14 and at about 03:45 UT in Fig. 15. The absolute value of each component is higher in the magnetosheath than in the solar wind but this is expected after crossing the bow shock. The Alfvén fluctuations in the central region of high speed streams appear to retain more of their original character after crossing the bow shock than do the compressional fluctuations at the leading edge.

Figure 16 shows the power spectral density associated with the parallel (compressional) component at Wind (top panel) and Geotail (bottom panel) on 9 April 1995. The color scales are different in each panel, in order to pull out the peaks in the spectrum. The total power increases as we move from the top panel to the bottom panel. The power at both Wind and Geotail is lower than in Fig. 7, but this is due to being in the central region of the high speed stream. There is a power enhancement centered at about 04:00 UT at Geotail that may be related to the enhancement at Wind centered at about 03:00 UT, taking into account timing differences between the satellite locations. The individual power peaks within this enhancement are similar but not identical for the two satellites. Peaks at Wind (top panel) are evident at 1.0, 1.5, and 1.9 mHz. Geotail power (bottom panel) is broadband below 1.5 mHz with another broad-band increase centered near 1.8 mHz and narrower peaks at 2.8 and 3.6 mHz. After this enhancement, there is a power reduction at Wind but at Geotail there is another power enhancement similar to the first one. Then we can see the edge of enhancements at Wind and Geotail that again, due to timing considerations, may be related.

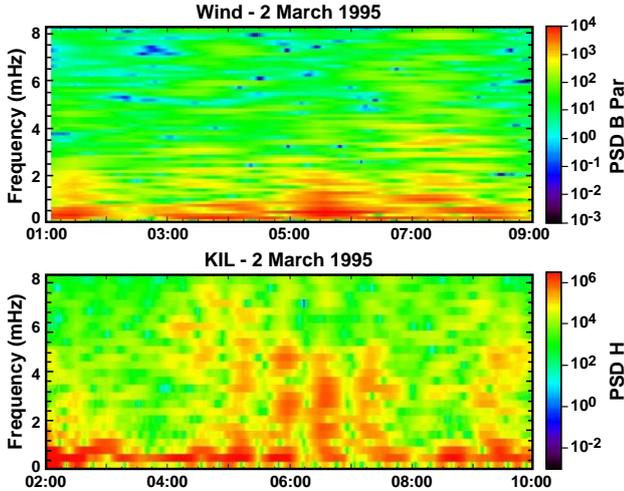


Fig. 17. PSD for Wind (top panel) and KIL (bottom panel) on 2 March 1995.

For the second example both Wind and KIL are making observations on the dawn side, with Wind $200 R_E$ upstream ($y = -5 R_E$). We do not show the waveforms for this example because there is no similarity between them. The fluctuations at Wind are Alfvénic, which is expected in the central region of a high speed stream. We are interested in comparing the Pc5 frequency range at Wind and KIL. Figure 17 shows the power spectral density associated with the parallel (compressional) component at Wind (top panel) and the H component at KIL (bottom panel) on 2 March 1995. There are four enhanced intervals at Wind: between 01:00 UT and 02:00 UT; centered at about 03:30 UT; between 05:00 UT and 06:00 UT; and centered at about 07:30 UT. There is an enhanced interval between 02:00 UT and 03:00 UT at KIL that may be related to the first interval at Wind, in each case the power is enhanced at frequencies less than about 2 mHz and is fairly broad-band. However, between 04:00 UT and 08:00 UT the KIL PSD (bottom panel) bears little resemblance to the Wind PSD (top panel). Both intervals are enhanced over ambient levels, but the nature of the enhancements are quite different. At Wind, the enhancements occur for frequencies less than about 2 mHz, except for the enhancement centered at about 07:30 UT, in which case the power is enhanced up to about 4 mHz. However, at the corresponding time at KIL, there is a power reduction. The KIL PSD plot (bottom panel) shows four distinct intervals of about 30 min each with broad-band enhancements stretching up to about 6 mHz. We again note that the cavity eigenfrequencies of 1.3 and 1.9 are both evident in the solar wind data during the largest broad-band increase at KIL at about 06:30 UT. With the possible exception of the interval at 06:30, it would seem that between 04:00 UT and 08:00 UT the power enhancements at Wind and KIL are not directly related.

4 Discussion and conclusions

A fundamental question that remains unanswered is “what drives Pc5 power in the magnetosphere?” There are two favorite hypotheses: (1) that fluctuations in the solar wind directly drive ULF power in the magnetosphere (e.g. Kepko et al., 2002) or that broad-band fluctuation power in the solar wind can lead to enhanced excitation of the magnetospheric cavity or waveguide modes, even if the spectral content of the upstream and magnetosphere are different (e.g. Wright and Rickard 1995); or (2) that the fast solar wind streams lead to an enhanced Kelvin-Helmholtz instability on the flanks, particularly the dawn flank, and that this leads to enhanced Pc5 power inside the magnetosphere through the surface mode hypothesis (e.g. Dungey, 1955; Miura, 1992; Anderson, 1994), or by the energisation of body waveguide modes (e.g. Mann et al., 1999). We focus our discussion on fluctuations and wave transmission to investigate (1) above. We can inquire if ULF fluctuations themselves, particularly in the Pc5 range, are found in high speed solar wind streams. If so, are they modulated crossing the bow shock? How do the magnetosheath fluctuations compare with the solar wind fluctuations? How do the solar wind or magnetosheath fluctuations interact with the magnetopause and magnetosphere? In answering these questions we can reach a deeper understanding of how energy is transmitted from the solar wind to the magnetosphere and what drives Pc5 power.

ULF fluctuations in the Pc5 range are found in high speed streams. All of the solar wind examples shown here, both at the leading edge and in the central region of high speed streams, exhibit fluctuations in the Pc5 range. These examples are representative of the 5-month data set and the discussion relates to all. The Pc5 frequency range dominates over higher frequency ULF fluctuations, with the power being several orders of magnitude greater than the Pc3 or Pc4 power inside high speed streams (not shown). Lower frequency fluctuations (under 1 mHz) have higher or comparable power to the Pc5 range, but these fluctuations may be less geo-effective than the Pc5 range. At the leading edge of high speed streams, fluctuations are generated in the compression region that forms when the fast wind catches up to the slow solar wind (e.g. Belcher and Davis, 1971). Belcher and Davis also identified the Alfvén waves in the central regions of high speed streams as solar generated. We also note the existence of recurrent sector boundary crossings during the first 6 months of 1995 (T. Hoeksema, Wilcox Magnetic Observatory). The northward progressions of the current sheet are nearly coincident with the high speed streams.

We have shown similar fluctuations from leading edge compression regions $200 R_E$ upstream and just in front of the bow shock (Figs. 5 and 6, respectively). Central region Alfvén fluctuations $200 R_E$ upstream also propagate inside the high speed streams and are similar to fluctuations observed just in front of the bow shock. This suggests that the fluctuations retain their characteristics over large spatial scales in the direction of propagation. We note a curiosity in this data set. Belcher and Davis (1971) showed that the com-

ponent of fluctuations parallel to the magnetic field direction was the smallest component; the largest was perpendicular to the magnetic field direction. By contrast, we find the two components to be essentially equivalent.

Pc5 range fluctuations are modulated in crossing the bow shock. We have shown a comparison between Wind and Geotail at the leading edge of a high speed stream in which Wind is in the solar wind and Geotail is in the magnetosheath (Figs. 8 and 9, respectively). As the fluctuations cross the bow shock the amplitude increases substantially and most similarity in the waveform is lost. The power is increased by more than a factor of 10 in the magnetosheath compared to the solar wind. By contrast, Alfvén fluctuations $200 R_E$ upstream retain similar features to the fluctuations in the magnetosheath (Figs. 14 and 15), but again turbulence increases. The Alfvén wave power is increased by slightly less than a factor of 10. This is in keeping with McKenzie and Westphal (1969; 1970), who found that fast magnetoacoustic longitudinal waves are greatly amplified on passage through the shock and that Alfvénic waves are moderately amplified. The overall power in the waves is enhanced by about a factor of 10, as seen in Figs. 2 and 3. Other fluctuations and discontinuities may be generated in the magnetosheath by the interaction between the bow shock and MHD discontinuities or Alfvén waves in the solar wind, as simulated by Lin et al. (1996). MHD waves may be reflected at the magnetopause back into the magnetosheath, as calculated by Kwok and Lee (1984). The end result would be that the fluctuations seen by Geotail in the magnetosheath would be a mix of waves and discontinuities. There could be some resemblance to the solar wind fluctuations, but it would be unusual to see identical fluctuations in the solar wind and magnetosheath, as evidenced by our observations.

How do the solar wind and magnetosheath fluctuations interact with the magnetopause and magnetosphere? Ideally to answer this question we would look at fluctuations in the solar wind, magnetosheath and magnetosphere simultaneously. With this data set, however, there are no intervals in which we have simultaneous dawn-side magnetosheath and ground-based waves. To obtain some insight into this question we have compared fluctuations in the solar wind to fluctuations in the magnetosheath and also compared fluctuations in the solar wind to pulsations on the ground, both at the leading edge and in the central region of high speed streams. We discussed above the comparison of solar wind and magnetosheath fluctuations. To recap, we have seen differences at the leading edge and in the central region. The solar wind Pc5 peaks are primarily under 4 mHz at the leading edge but only 2 mHz or less in the central region. The fluctuations in the magnetosheath not only have more power, but power that extends to higher Pc5 frequencies at the leading edge compared to the central region of high speed streams. The leading edge compression regions appear to be more active and powerful than the central region. How does this affect the ground-based measurements? For our example at the leading edge compression region of a high speed stream on 19 June 1995, we found similar general trends in the PSD at Wind

and KIL (Fig. 13), but no one-to-one match up between individual features at Wind with those at KIL. The similarity in general trends suggests that this may be a driven system. For our example in the central region of a high speed stream on 2 March 1995, we found, with the exception of the first enhanced interval (between 02:00 UT and 03:00 UT at KIL), that there was little resemblance between the PSD at Wind and KIL. The later intervals were enhanced over ambient levels, but the nature of the enhancements were quite different.

We have been discussing a particular solar wind structure, i.e. high speed streams, in which ULF fluctuations in the Pc5 range play a dominant role. We have shown a clear correlation between total power (Pc5 range) in the solar wind, in the magnetosheath, and on the ground over a 5-month period of high speed streams in 1995. The correlation extends from $200 R_E$ upstream, to just upstream from the bow shock, to the magnetosheath and on the ground (Figs. 2 and 3). At some times, particular frequencies of the spectral power in the solar wind and magnetosheath are nearly coincident, though at other times the frequencies don't match up. Between the solar wind and ground-based measurements, the frequencies are not the same, though they generally do fall within the same range from 1–4 mHz (e.g. Fig. 13), and can, but don't always, coincide with the cavity eigenfrequencies of 1.3, 1.9, 2.6, and 3.4 mHz given by Samson et al. (1992) and Samson and Rankin (1994). As we noted previously, in some cases large increases in broad-band power at KIL were coincident with power enhancements in discrete frequency bands in the solar wind spanning the range of the broad-band power seen on the ground.

Without the dawn-side magnetosheath measurements we can only speculate on the driver of Pc5 power in the magnetosphere, based on our 5-month database illustrated in these examples. Longitudinal or compressional fluctuations appear to have a different effect on the magnetosheath and ground-based measurements than do Alfvén fluctuations. For compressional fluctuations at the leading edge, it could be that the MHD cavity is driven by random boundary motion with a broad-band frequency spectrum in the correct range, as suggested by Wright and Rickard (1995). The compressional fluctuations in the Pc5 range also could drive magnetopause surface waves with periods in the Pc5 range. Engebretson et al. (1998 and sources therein) suggested that if the compression regions at the leading edges of high speed streams contain waves in the Pc5 range, they could provide a source of wave energy to the magnetosphere, or that the waves could act as seed perturbations to drive boundary displacements that are amplified by the Kelvin-Helmholtz instability. It is also possible that the magnetopause is open at these times and that MHD waves are transmitted from the magnetosheath to the magnetosphere as suggested by Kwok and Lee (1984). For the Alfvén fluctuations in the central region, features seen on the ground at KIL but not in the solar wind at Wind must have another generating mechanism. Alfvén fluctuations may not drive the magnetosphere in the same way that compressional fluctuations do. The ground-based morning sector H -component Pc5 power is highly correlated with the

solar wind speed (correlation coefficient ~ 0.8), being better correlated with solar wind speed than with the solar wind Pc5 power that we have considered explicitly in this paper (correlation coefficient ~ 0.61). Features seen on the ground but not in the solar wind might be related to hypothesis (2) fast solar wind streams leading to an enhanced Kelvin-Helmholtz instability on the flanks, and to enhanced Pc5 power inside the magnetosphere.

We have examined the 5-month period in the descending phase of the last solar cycle. More data should be available for high speed streams in the current cycle's descending phase, allowing for comparisons of solar wind, magnetosheath and ground-based data simultaneously on the dawn side. The answer to the driver of Pc5 pulsations in the magnetosphere may then be identified. The dawn side of the bow shock is recognized as the local time region that contains the maximum in Pc5 power on the ground (Engebretson et al., 1998), yet some ground-based statistics show a local time distribution peaked approximately equally at dawn and dusk (Anderson, 1994 and sources therein). There is also some disparity between ground-based and satellite observations of Pc5 pulsations in the magnetosphere. Observations in space show a dawn maximum from magnetic field measurements, but a broad dayside distribution from electric field measurements (Anderson, 1994 and sources therein). These controversies remain to be solved.

With our limited data set, we draw the following conclusions:

- ULF fluctuations in the Pc5 range are found in high speed streams; they are non-Alfvénic at the leading edge and Alfvénic in the central region.
- Compressional and Alfvénic fluctuations are modulated at the bow shock; Alfvénic features of the waveforms are better preserved in the magnetosheath. Overall turbulence and wave power is enhanced by about a factor of 10.
- Parallel (compressional) and perpendicular (transverse) power are at comparable levels, both in the compression region and in the central region of high speed streams. This is true in the solar wind and in the magnetosheath.
- Both the total parallel and perpendicular Pc5 power in the solar wind (and, to a lesser extent in the magnetosheath) correlate well with the total Pc5 power of the ground-based H -component magnetic field.
- ULF fluctuations in the solar wind and magnetosheath during high speed streams are common at frequencies from 1–4 mHz and can coincide with the cavity eigenfrequencies of 1.3, 1.9, 2.6, and 3.4 mHz, however, other discrete frequencies are also often seen.

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References

- Anderson, B. J., Potemra, T. A., Zanetti, L. J., and Engebretson, M. J.: Statistical correlations between Pc3-5 pulsations and solar wind/IMF parameters and geomagnetic indices, in *Physics of Space Plasmas (1990)*, SPI Conf. Reprint Ser., vol. 10, edited by Chang T., Crew, G.B., and Jasperse, J. R., page 419 Scientific Publishing Inc., Cambridge, Mass., 419–429, 1991.
- Anderson, B. J.: An overview of spacecraft observations of 10 s to 600 s period magnetic pulsations in the Earth's magnetosphere, in: *Solar Wind Sources of Magnetospheric Ultra-Low-Frequency Waves*, edited by Engebretson, M. J., Takahashi, K., and Scholer, M., AGU, Wash., D.C., 1994.
- Baker, D. N., Pulkkinen, T. I., Li, K. X., and Kanekal, S. G. et al.: Coronal mass ejections, magnetic clouds and relativistic magnetospheric electron events: *ISTP J. Geophys. Res.*, 103, 17 279–17 291, 1998.
- Barnes, A.: Collisionless damping of hydromagnetic waves, *Phys. Fluids*, 9, 1483, 1966.
- Belcher, J. W., and Davis, L.: Large-amplitude Alfvén waves in interplanetary medium, 2, *J. Geophys. Res.*, 74, 2303, 1971.
- Burlaga, L. F., and Ogilvie, K. W.: Magnetic and thermal pressures in the solar wind, *Sol. Phys.*, 15, 61, 1970.
- Coleman, P. J.: Turbulence, viscosity, and dissipation in the solar wind plasma, *Astrophys. J.*, 153, 371, 1968.
- Crooker, N. U., Eastman, T. E., Frank, L. A., Smith, E. J., and Russell, C. T.: Energetic magnetosheath ions and the interplanetary magnetic field orientation, *J. Geophys. Res.*, 86, 4455, 1981.
- Dungey, J. W.: *Electrodynamics of the outer atmosphere*, Rep. 69, Ions. Res. Lab. Pa. State Univ., University Park, 1954.
- Dungey, J. W.: *Electrodynamics of the outer atmosphere*, in *Proceedings of the Ionosphere Conference*, p. 225, The Physical Society of London, 1955.
- Engebretson, M. J., Lin, N., Baumjohann, W., and Luehr, H., et al.: A comparison of ULF fluctuations in the solar wind, magnetosheath, and dayside magnetosphere 1. Magnetosheath Morphology, *J. Geophys. Res.*, 96, 3441, 1991.
- Engebretson, M. J., Glassmeier, K.-H., and Stellmacher, M.: The dependence of high-latitude Pc5 power on solar wind velocity and phase of high-speed solar wind streams, *J. Geophys. Res.*, 103, 26 271, 1998.
- Fairfield, D. H., Baumjohann, W., Paschmann, G., Luehr, H., and Sibeck, D.G.: Upstream pressure variations associated with the bow shock and their effects on the magnetosphere, *J. Geophys. Res.*, 95, 3773, 1990.
- Glassmier, K. H.: Ultralow-frequency pulsations: Earth and Jupiter compared, *Adv. Space Res.*, 16, 209, 1995.
- Greenstadt, E. W., and Olson, J. V.: A contribution to ULF activity in the Pc 3-4 range correlated with IMF radial orientation, *J. Geophys. Res.*, 82, 4991, 1977.
- Greenstadt, E. W., Olson, J. V., Loewen, P. D., Singer, H. J., and Russell, C. T.: Correlation of Pc3/4/5 activity with solar wind speed, *J. Geophys. Res.*, 84, 6694, 1979.
- Greenstadt, E. W., Mellott, M. M., McPherron, R. L., Russell, C. T., Singer, H. J., and Knecht, D. J.: Transfer of pulsation-related wave activity across the magnetopause: Observations of corre-

- sponding spectra by ISEE-1 and ISEE-2, *Geophys. Res. Lett.*, 10, 659, 1983.
- Greenstadt, E. W., and Russell, C. T.: Stimulation of exogenic, daytime geomagnetic pulsations: a global perspective, in: *Solar Wind Sources of Magnetospheric Ultra-Low-Frequency Waves*, edited by Engebretson, M. J., Takahashi, K., and Scholer, M., AGU, Wash., D.C., 1994.
- Harrold, B. G. and Samson, J. C.: Standing ULF modes of the magnetosphere: a theory, *Geophys. Res. Lett.*, 19, 1811, 1992.
- Hughes, W. J.: Magnetospheric ULF waves: a tutorial with a historical perspective, in: *Solar Wind Sources of Magnetospheric Ultra-Low-Frequency Waves*, edited by Engebretson, M. J., Takahashi, K., and Scholer, M., AGU, Wash., D.C., 1994.
- Kepko, L., Spence, H. E., and Singer, H. J.: ULF waves in the solar wind as direct drivers of magnetospheric pulsations, *J. Geophys. Res.*, 29, 39–1, 2002.
- Kivelson, M. G. and Russell, C. T.: *Introduction to Space Physics*, Cambridge University Press, Cambridge, U.K. 1995.
- Kokubun, S., Erikson, K. N., Fritz, T. A., and McPherron, R. L.: Local time asymmetry of Pc4-5 pulsations and associated particle modulations at synchronous orbit, *J. Geophys. Res.*, 94, 6607, 1989.
- Kokubun, S., Yamamoto, T., Acuna, M. H., Hayashi, Shiokawa, K. K., and Kawano, H.: The GEOTAIL magnetic field experiment, *J. Geomag. Geoelectr.*, 46, 7–21, 1994.
- Kwok, Y. C., and Lee, L. C.: Transmission of magnetohydrodynamic waves through the rotational discontinuity at the Earth's magnetopause, *J. Geophys. Res.*, 89, 10 697–10 708, 1984.
- Lepping, R. P., Acuna, M. H., Burlaga, L. F., Farrell, W. M., et al.: The Wind Magnetic Field Investigation, in *The Global Geospace Mission*, ed. C.T. Russell, Kluwer Academic Publishers, 207–229, 1995.
- Lin, N., Engebretson, M. J., McPherron, R. L., Kivelson, M. G., Baumjohann, W., Luehr, H., et al.: A comparison of ULF fluctuations in the solar wind, magnetosheath, and dayside magnetosphere 2: Field and Plasma conditions in the magnetosheath, *J. Geophys. Res.*, 96, 3455, 1991.
- Lin, Y., Lee, L. C., and Yan, M.: Generation of dynamic pressure pulses downstream of the bow shock by variations in the interplanetary magnetic field orientation, *J. Geophys. Res.*, 101, 479, 1996.
- Luhr, H., Aylward, A., Buchert, S. C., Pajunpdd, A., Pajunpdd, K., Holmboe, T., and Zaleski, S. M.: Westward moving dynamic substorm features observed with the IMAGE magnetometer network and other ground-based instruments. *Ann. Geophysicae*, 16, 425–440, 1998.
- Mann, I. R., Wright, A. N., Mills, K. J., and Nakariakov, V. M.: Excitation of magnetospheric waveguide modes by magnetosheath flows, *J. Geophys. Res.*, 104, 333–353, 1999.
- Mathie, R. A. and Mann, I. R.: A correlation between extended intervals of ULF wave power and storm-time geosynchronous relativistic electron flux enhancements, *Geophys. Res. Lett.*, 27, 3261, 2000.
- McKenzie, J. F.: Hydromagnetic wave interaction with the magnetopause and bow shock, Transmission of Alfvén waves through the Earth's bow shock, *Planet. Space Sci.*, 18, 1–23, 1970.
- McKenzie, J. F., and Westphal, K. O.: Transmission of Alfvén waves through the Earth's bow shock, *Planet. Space Sci.*, 17, 1029–1037, 1969.
- McKenzie, J. F. and Westphal, K. O.: Interaction of hydromagnetic waves with hydromagnetic shocks, *Phys. Fluids*, 13, 630–640, 1970.
- Miura, A.: Kelvin-Helmholtz instability at the magnetospheric boundary: Dependence on the magnetosheath sonic Mach number, *J. Geophys. Res.*, 97, 10 665, 1992.
- Ogilvie, K. W., Chornay, D. J., Fritzenreiter, R. J., Hunsaker, F., et al.: SWE, Comprehensive Plasma Instrument for Wind Spacecraft, in *Global Geospace Mission*, ed. C.T. Russell, Kluwer Academic Pub., 55, 1995.
- Rostoker, G., Skone, S., and Baker, D. N.: On the origin of relativistic electrons in the magnetosphere associated with geomagnetic storms, *Geophys. Res. Lett.*, 25, 19, 3701, 1998.
- Ruohoniemi, J. M., Greenwald, R. A., Baker, K. B., and Samson, J. C.: HF radar observations of Pc5 field line resonances in the midnight/early morning MLT sector, *J. Geophys. Res.*, 96, 15 697, 1991.
- Russell, C. T., and Elphic, R. C.: Initial ISEE magnetometer results: magnetopause observations, *Space Sci. Rev.*, 22, 681, 1978.
- Samson, J. C., Harrold, B. G., Ruohoniemi, J. M., Greenwald, R. A., and Walker, A. D. M.: Field line resonances associated with MHD waveguides in the magnetosphere, *Geophys. Res. Lett.*, 19 441–19 444, 1992.
- Samson, J. C. and Rankin, R.: The coupling of solar wind energy to MHD cavity modes, waveguide modes, and field line resonances in the Earth's magnetosphere, in: *Solar Wind Sources of Magnetospheric Ultra-Low-Frequency Waves*, edited by Engebretson, M. J., Takahashi, K., and Scholer, M.: AGU, Wash., D.C., 1994.
- Sibeck, D. G., Baumjohann, W., Elphic, R. C., Fairfield, D. H., et al.: The magnetospheric response to 8-minute period strong-amplitude upstream pressure variations, *J. Geophys. Res.*, 94, 2505, 1989.
- Sibeck, D. G., Takahashi, K., Kokubun, S., Mukai, T., Ogilvie, K. W., and Szabo, A.: A case study of oppositely propagating Alfvénic fluctuations in the solar wind and magnetosheath, *Geophys. Res. Lett.*, 24, 3133, 1997.
- Sibeck, D. G., Phan, T.-D., Lin, R. P., Lepping, R. P., Mukai, T., and Kokubun, S.: A survey of MHD waves in the magnetosheath: International Solar Terrestrial Program observations, *J. Geophys. Res.*, 105, 129, 2000.
- Sonnerup, B. U. O., Paschmann, G., Papamastorakis, I., Sckopke, N., Haerendel, G., Bame, S. J., Asbridge, J. R., Gosling, J. T., and Russell, C. T.: Evidence of magnetic field reconnection at the Earth's magnetopause, *J. Geophys. Res.*, 86, 10 049, 1981.
- Southwood, D. J.: Some features of field line resonances in the magnetosphere, *Planet. Space Sci.*, 22, 483, 1974.
- Tsurutani, B. T., Gonzalez, W. D., Gonzalez, A. L. C., Tang, F., Arballo, J. K., and Okada, M.: Interplanetary origin of geomagnetic activity in the declining phase of the solar cycle, *J. Geophys. Res.*, 100, 21 717, 1995.
- Wolfe, A. and Kaufmann, R. L.: MHD wave transmission and production near the magnetopause, *J. Geophys. Res.*, 80, 1764, 1975.
- Wright, A. N. and Rickard, G. J.: A numerical study of resonant absorption in a magnetohydrodynamic cavity driven by a broadband spectrum, *Astrophys. J.*, 444, 458–470, 1995.