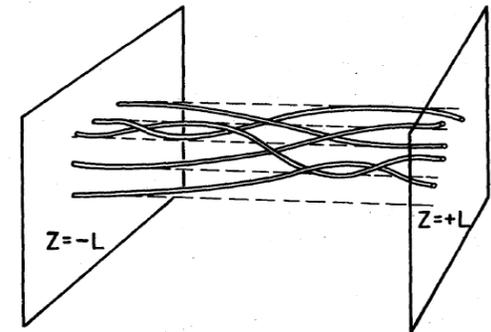
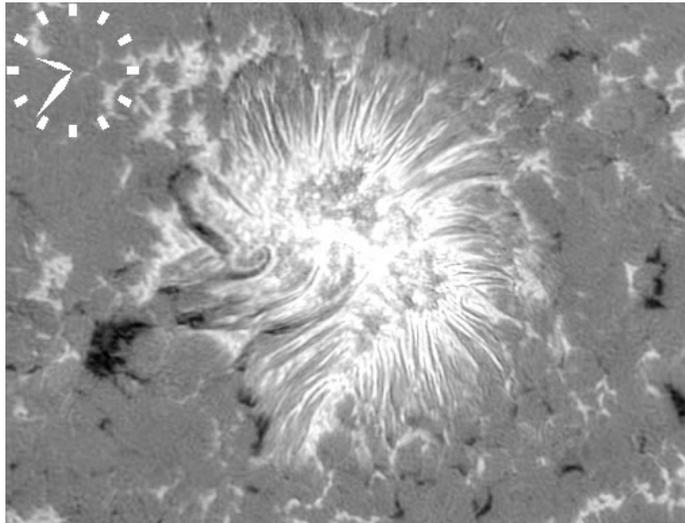


# Reconnection and heating in braided coronal loops

D. Pontin, GSFC 5<sup>th</sup> March 2025

# Braiding mechanism for coronal heating

- Field lines tangled by convective photospheric motions
- In corona  $R_m \gg 1$  so evolution 'ideal' almost everywhere: topology preserved
- Parker argues that perturbed field can't relax to smooth equilibrium except in certain non-generic cases.
- Consequence: tangential discontinuities, i.e. current sheets, form  $\rightarrow$  reconnection and heating.



Parker, E.N., ApJ, **174**, 499 (1972).

REVIEW ARTICLE

The Parker problem: existence of smooth force-free fields and coronal heating

David I. Pontin<sup>1,2</sup> · Gunnar Hornig<sup>2</sup>

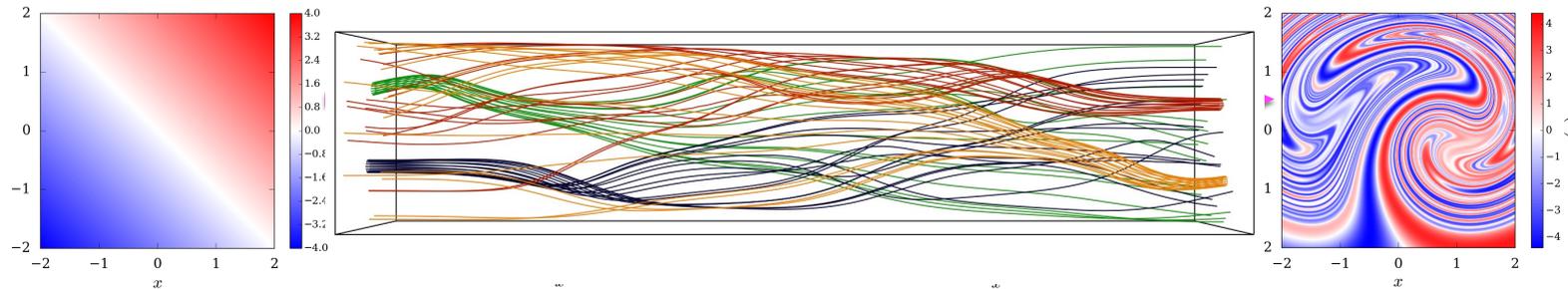
Living Reviews in Solar Physics (2020)17:5  
<https://doi.org/10.1007/s41116-020-00026-5>

# Parker's braiding mechanism

- 1) Braiding by surface flows
- 2) formation of current sheets
- 3) reconnection and energy deposition
- 4) plasma response

- Parker's claim is that no smooth equilibrium exists for braided fields → substantial debate without a consensus.
- No hard proof of current sheet (T.D.) formation in a geometry matching Parker's.
- But this does not matter....!
- Current sheets get progressively thinner as braid complexity increases (van Ballegooijen 1988, Mikić+ 1989, Pontin & Hornig 2015)
- **In solar corona continual braiding will inevitably lead to reconnection onset**

## (Finite) current layers in equilibrium



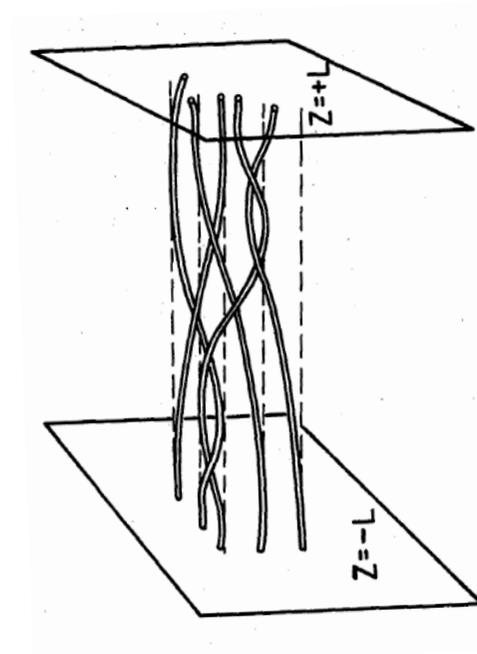
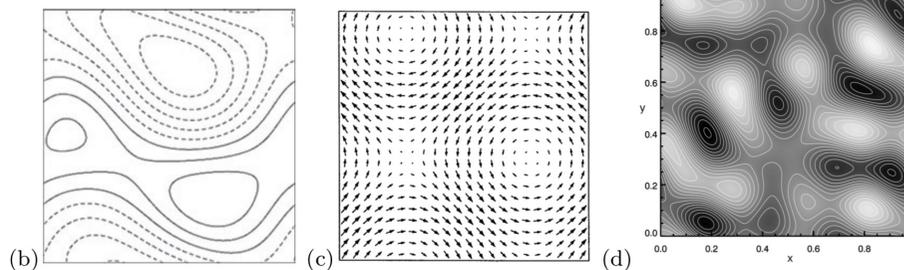
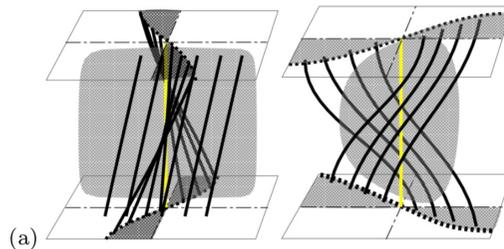
- Recall equilibrium satisfies  $\mathbf{J} \times \mathbf{B} = (\nabla \times \mathbf{B}) \times \mathbf{B} \approx 0$
- Suppose that  $\nabla \times \mathbf{B} = \alpha \mathbf{B}$  so that  $\mathbf{B} \cdot \nabla \alpha = 0$  and  $\alpha$  const along B-lines
- Suppose  $\alpha$  has distribution with length scale  $\ell$  on (e.g.) lower boundary
- Mapping along field lines,  $\alpha$  naturally exhibits length scales on order of field line mapping on upper boundary:  $\ell \times \lambda_{\min}$  ( $\lambda_{\min}$  smallest eigenvalue of mapping Jacobian DF)
- $\alpha = \mathbf{J} \cdot \mathbf{B} / B^2$ , and so assuming  $|\mathbf{B}| \sim O(1)$  then  $J_{\parallel}$  also has length scales on order of  $\ell \times \lambda_{\min}$ .
- Thus: if smooth force-free equilibrium exists, it must contain current layers at least as thin as layers in field line mapping

[Pontin & Hornig 2015, c.f. van Ballegooijen 1988]

# Flux braiding simulations

# Continuous driving in a resistive plasma ("flux braiding")

- Many different approaches:
  - Driven vs relaxation
  - Full MHD vs Reduced MHD
  - Domain aspect ratios from 1:1 to 1:10
  - Different driving velocities (sequence of shears, superposition of vortices,...)



## REVIEW ARTICLE

The Parker problem: existence of smooth force-free fields and coronal heating

David I. Pontin<sup>1,2</sup> · Gunnar Hornig<sup>2</sup>

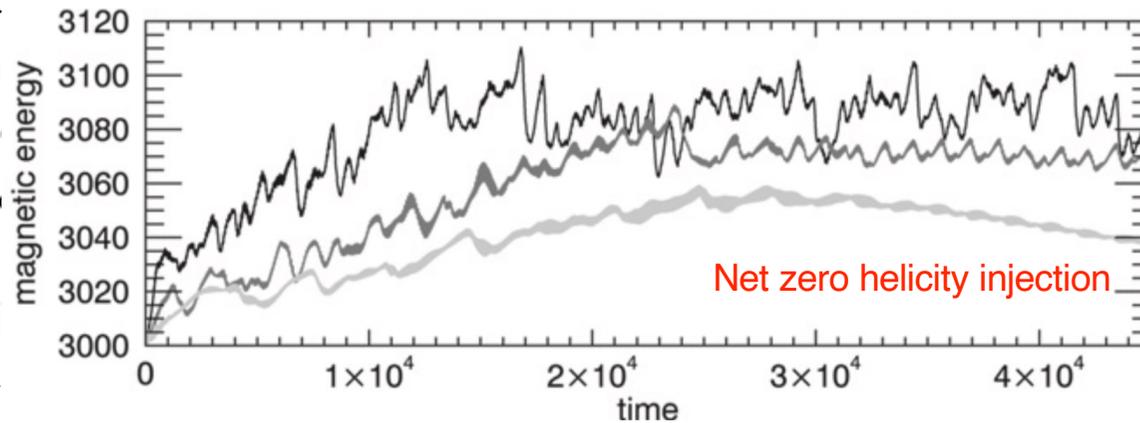
Living Reviews in Solar Physics (2020)17:5  
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## Main results of flux braiding experiments

- Thin ribbons of current form, thickness decreases exponentially in time.
- Evolution resembles anisotropic turbulent cascade (either driven or decaying).
- Scaling exponents in stat. steady state depend on loop length, field strength...
- $E_m > E_k$  by a factor of 10 at least.
- $E_m$  (equivalently  $\langle B_{\perp} \rangle$ ) scales with driving speed  $v_0$ . Also scales with  $R_m$ , at least for  $R_m \leq 10^3$ .
- Geometry of driving has minimal importance, but helicity injection rate or correlation time may affect characteristics of energy release.
  
- Outstanding issues:
  - Most studies of scaling with  $R_m$  use RMHD: for large  $R_m$ ,  $B_{\perp}$  may grow sufficiently large to invalidate assumptions.
  - ...and Poynting flux into volume depends on coronal state.
  - Coronal state is a function of  $R_m$
  - Question of timescales: driving timescale vs energy release timescale
  - Effect of plasma/magnetic field structuring of chromosphere/carpet

## Main results of flux braiding experiments

- Thin ribbons of current form, thickness decreases exponentially in time.
- Evolution resembles anisotropic turbulent cascade (either driven or decaying).
- Scaling exponent  $\beta \approx 1$  (at least for correlation time  $\tau_c \ll \tau_{\text{inertial}}$ ).
- $E_m > E_k$  by a factor of  $\sim 10$ .
- $E_m$  (equivalently  $\tau_c$ ) increases with  $R_m$  at least for  $R_m \leq 10^3$ .
- Geometry of driver may affect character of cascade.

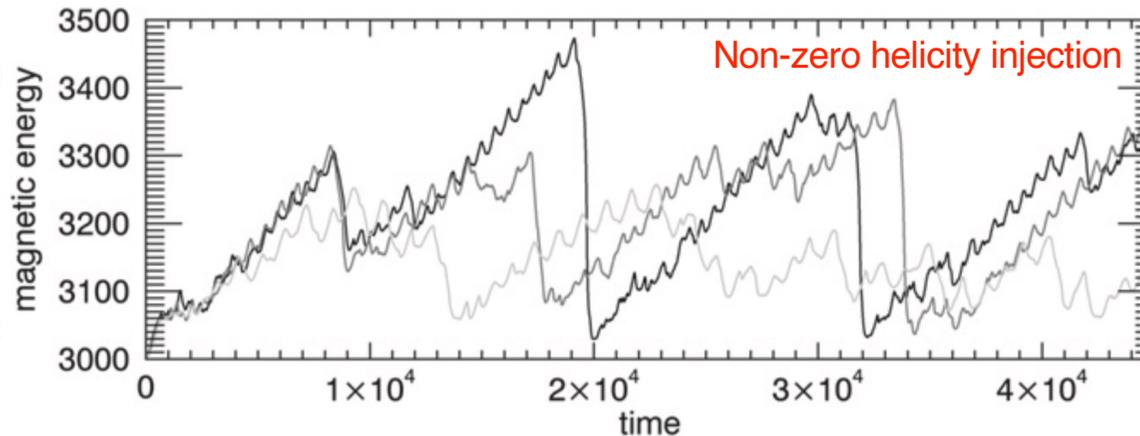


h...

least for

correlation time

- Outstanding issues
  - Most studies of flux braiding tend to invalidate as  $R_m$  increases.
  - ...and Poynting flux is not conserved.
  - Coronal state is not stationary.
  - Question of time scale for energy transfer.
  - Effect of plasma/magnetic field structuring of chromosphere/corona.



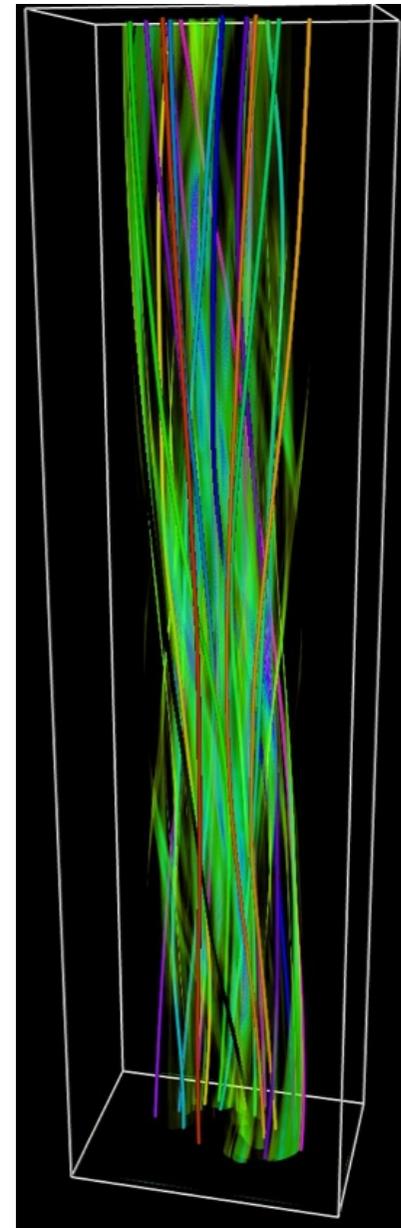
sufficiently large

## Main results of flux braiding experiments

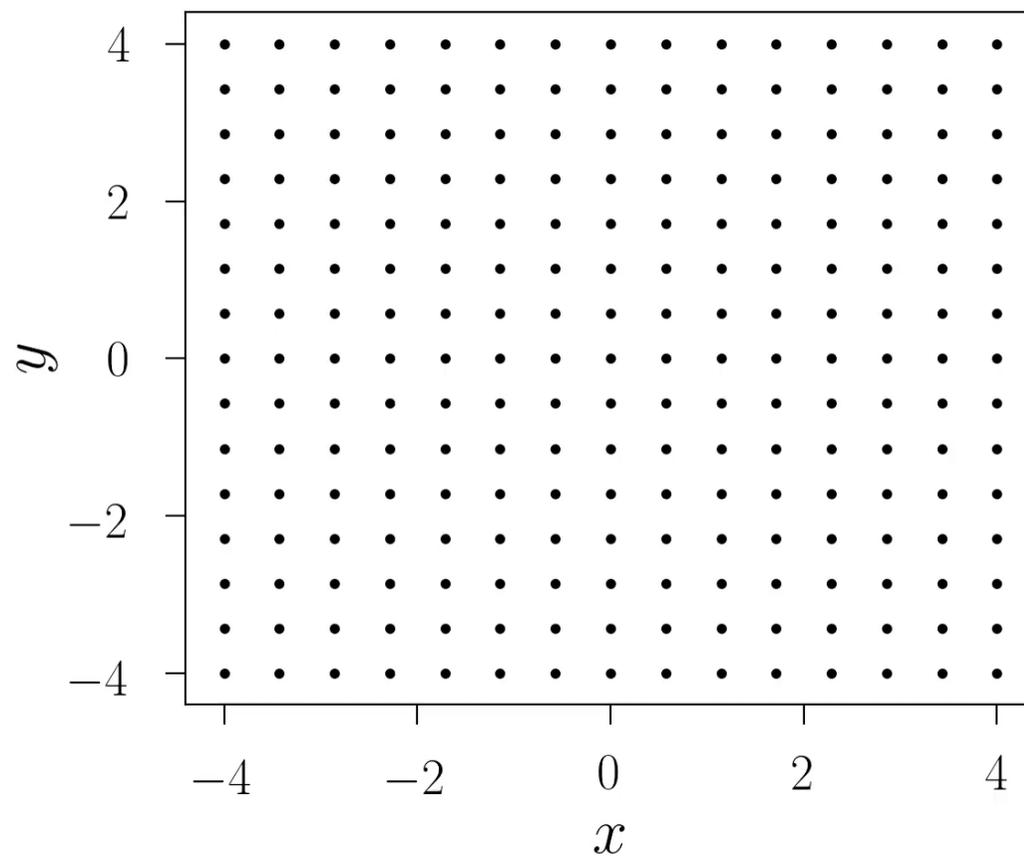
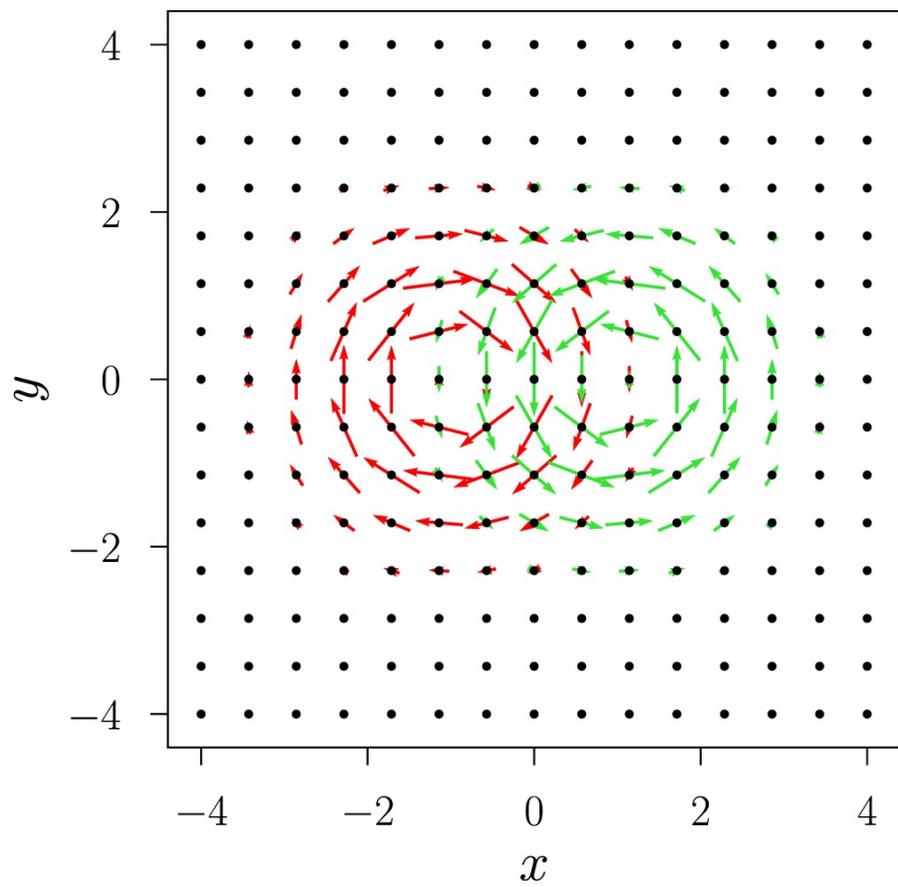
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  - Effect of plasma/magnetic field structuring of chromosphere/carpet

## MHD simulation – flux braiding

- Start from uniform magnetic field
- Apply driving continuously on line-tied boundaries
- Other boundaries periodic
- Solve resistive MHD equations (no thermal conduction, radiative losses). Hyper-resistivity/viscosity
- 6x6x48 Mm
- Plasma-beta: 0.05
- 640x264x264; 1280x528x528; 1280x1056x1056
- Low and Med runs go to  $t \sim 500$  which is  $\sim 12$  periods of driver. High run goes only until  $t \sim 250$

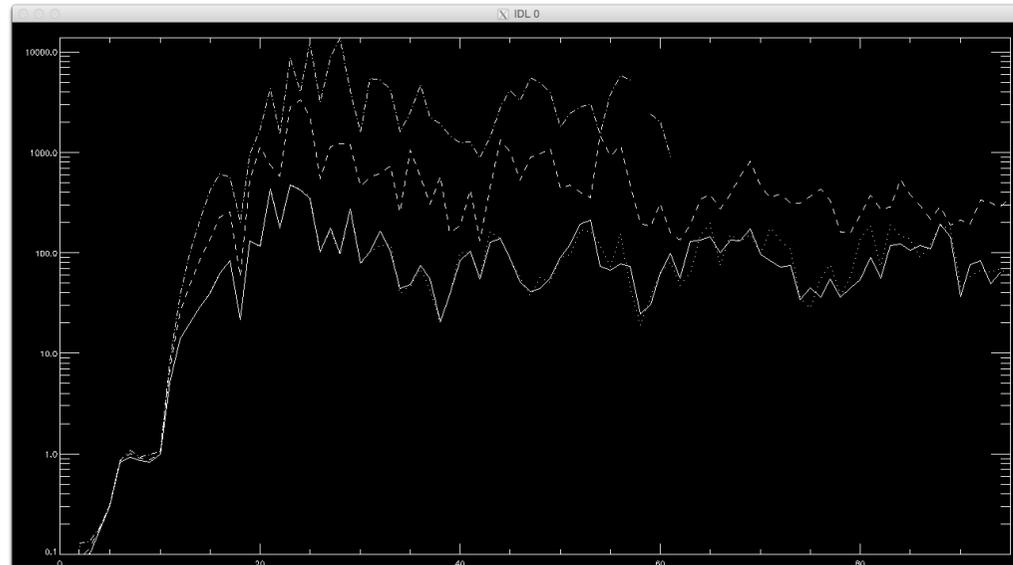


# Boundary driving

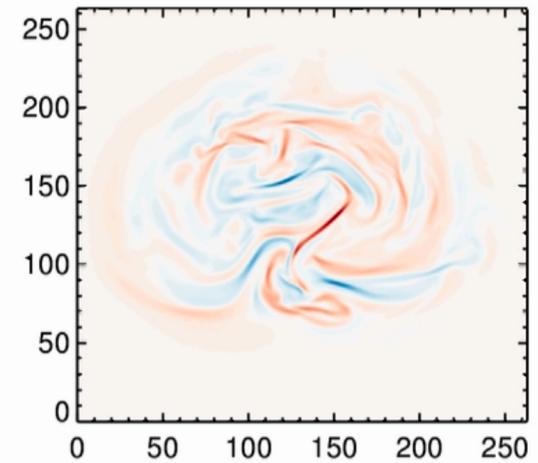
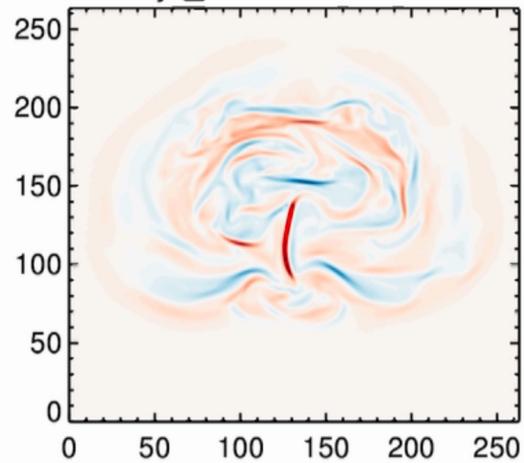
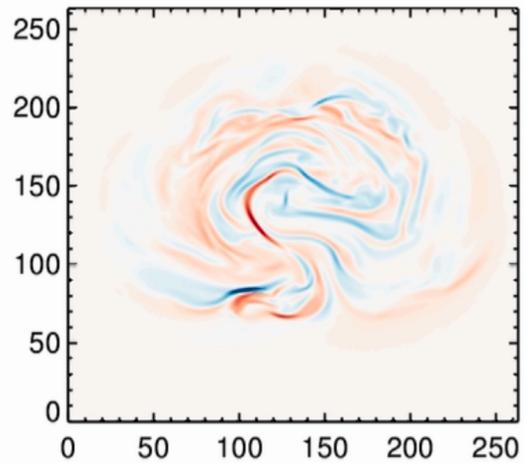


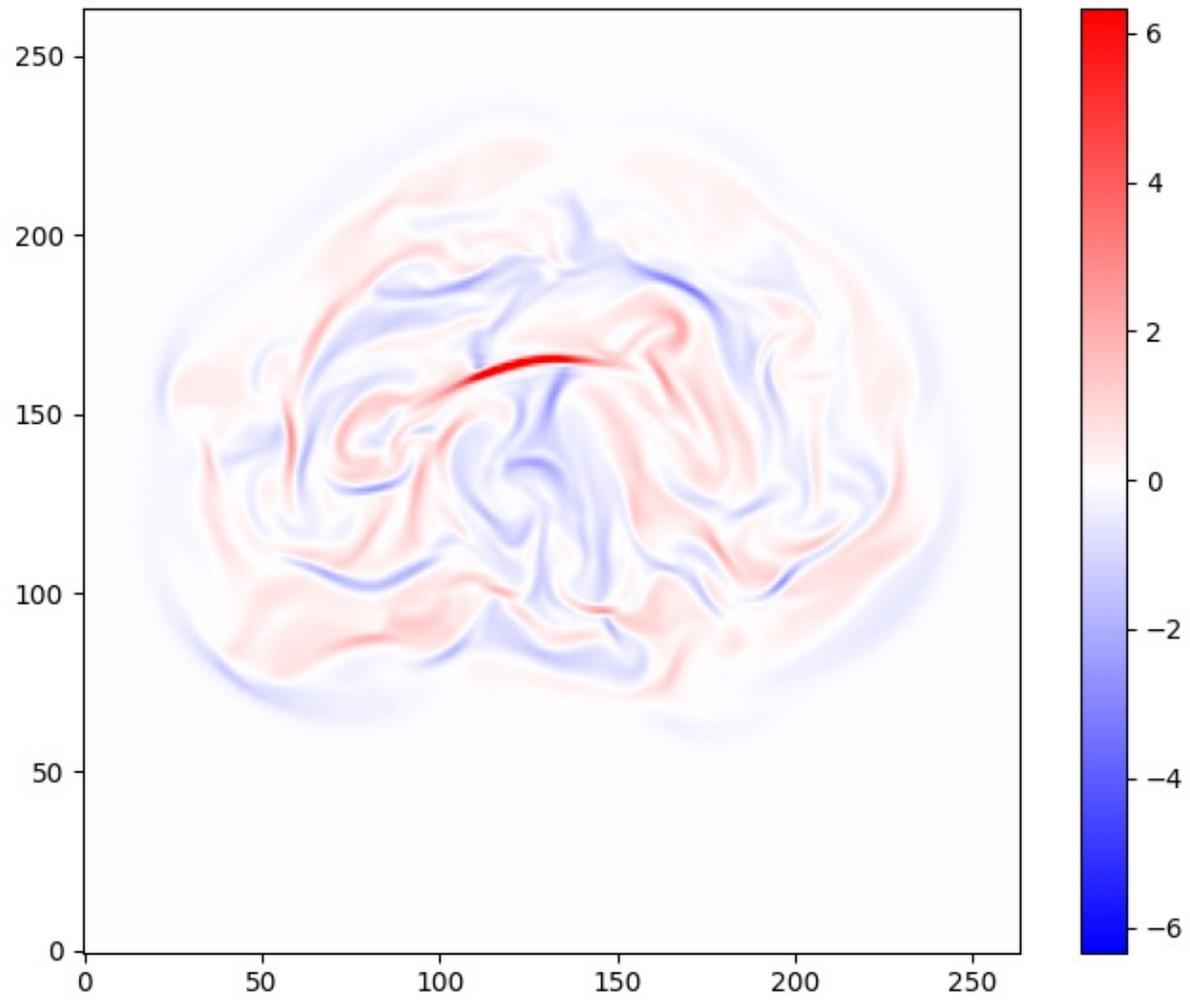
# Evolution of axial current with time

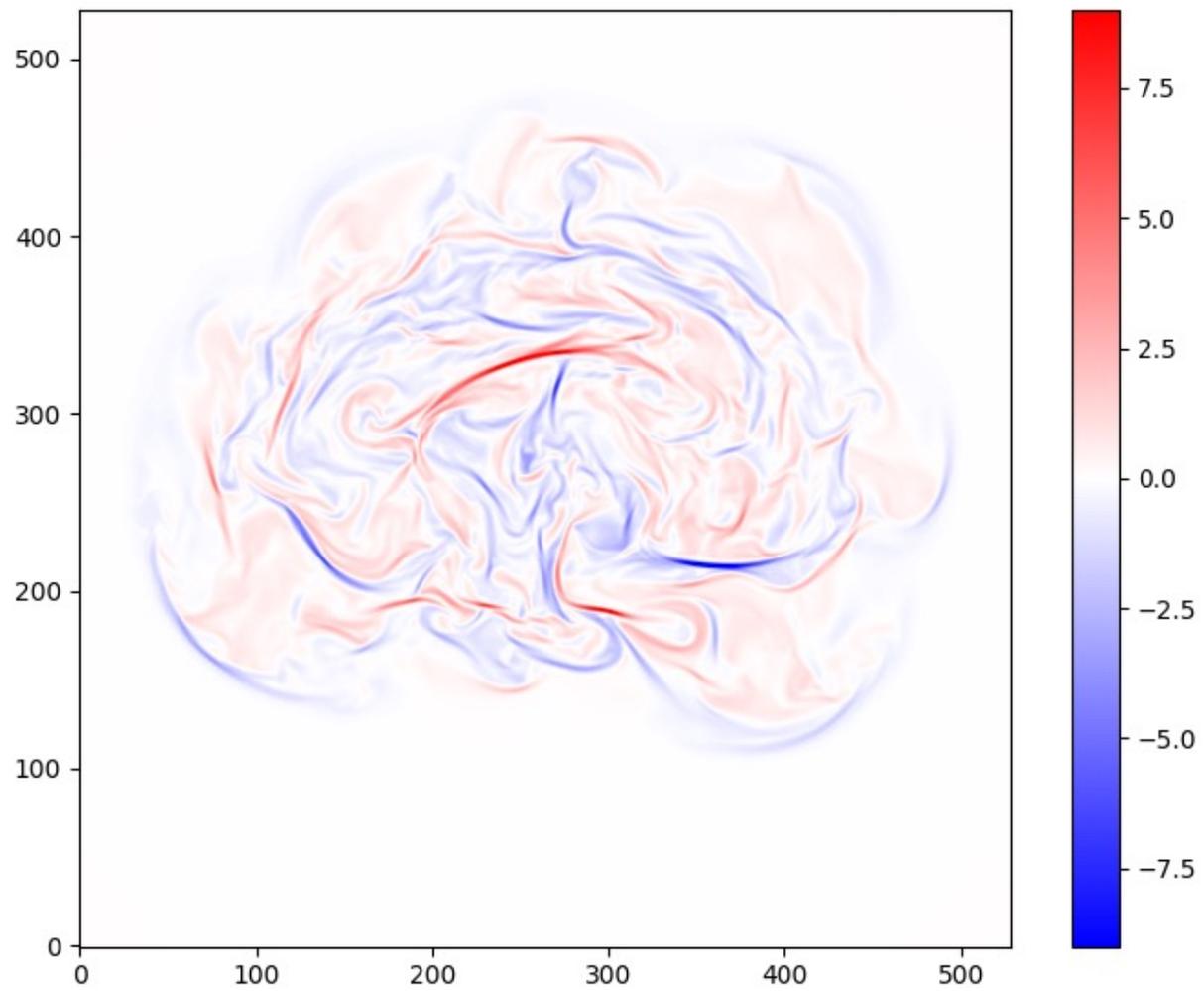
$J_{\max}$

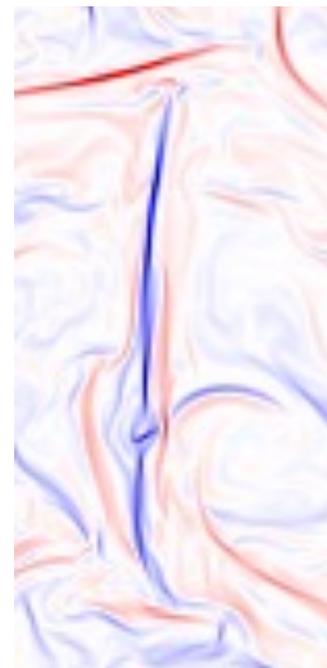
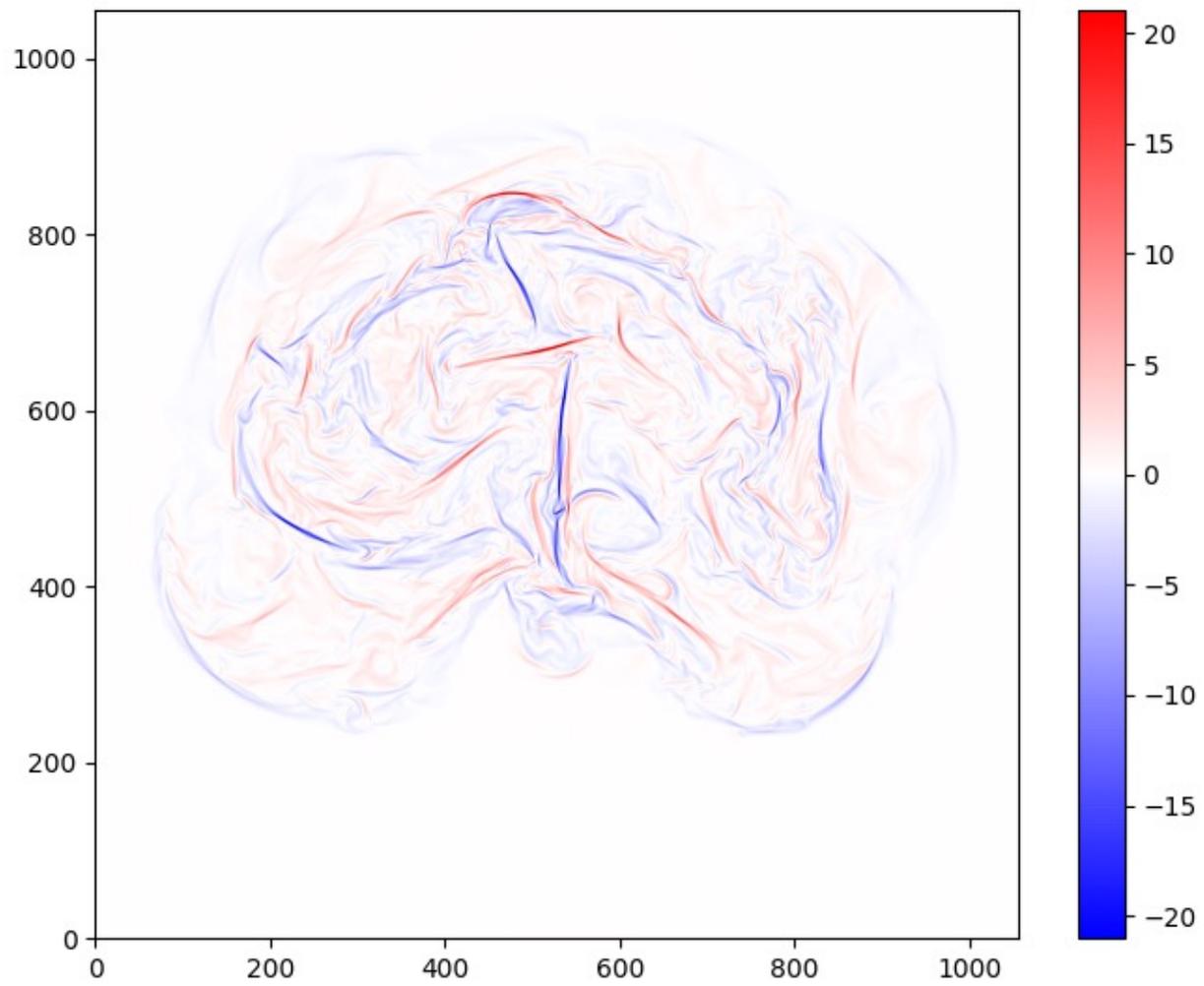


$jx_{\max} = 7.53631$



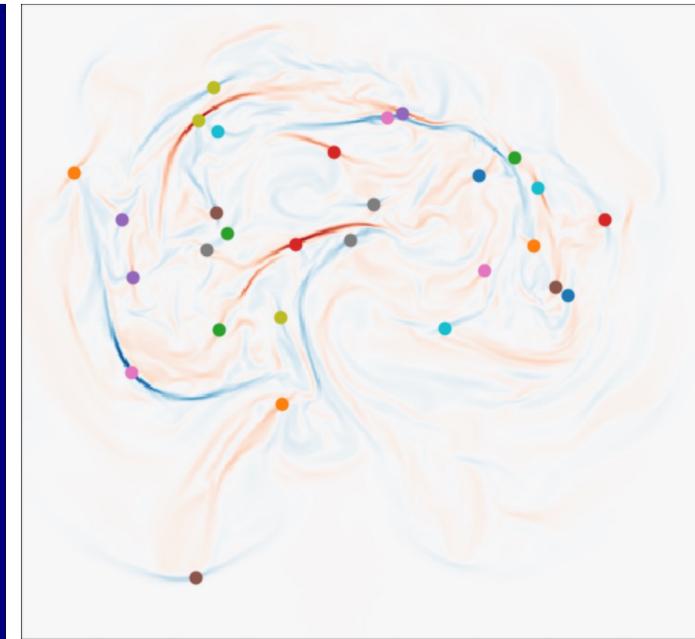
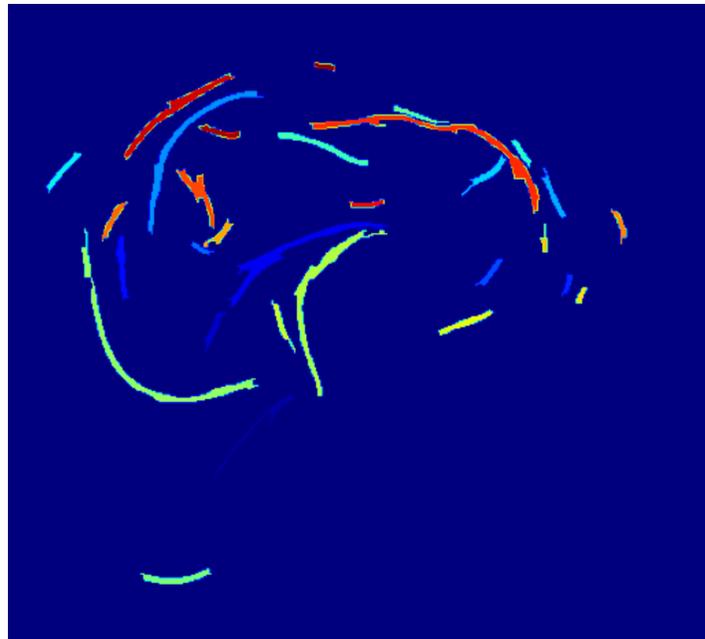
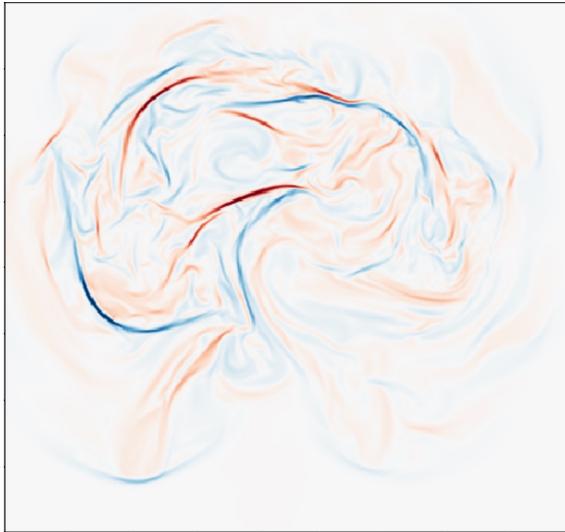




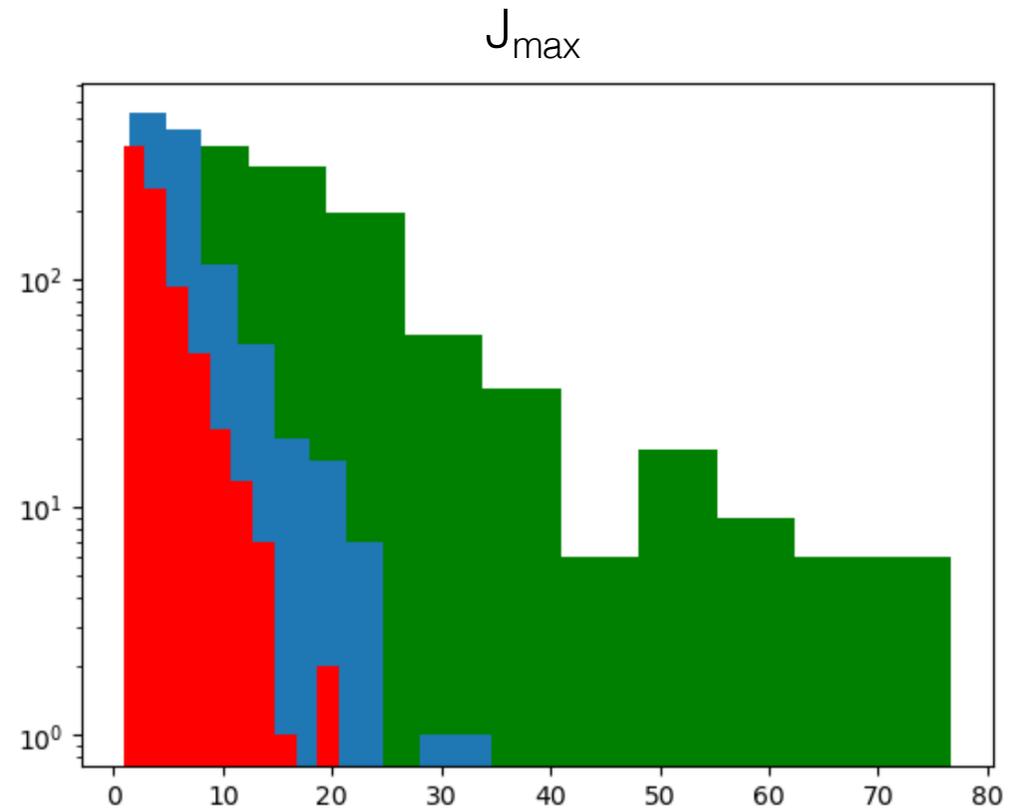
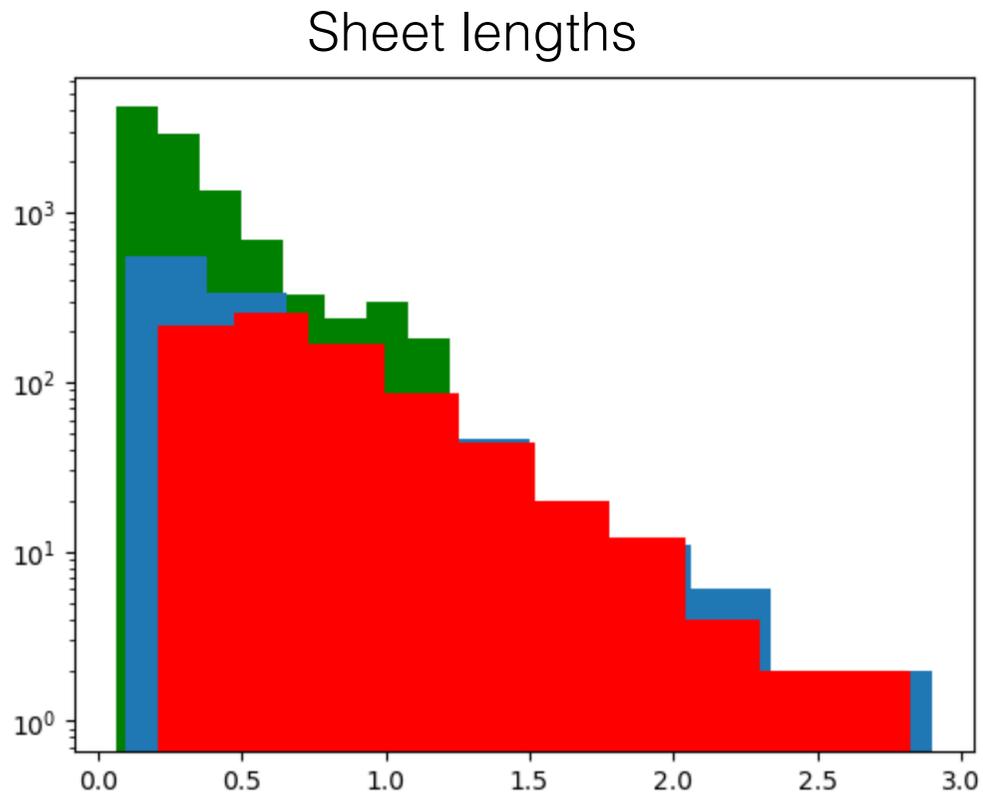


## Identify rec regions

- “Canny” edge detection
- Separate  $J_z > 0$  and  $J_z < 0$
- Morphological closing
- Region labelling
- Identify end and centres of regions
- So far just in midplane

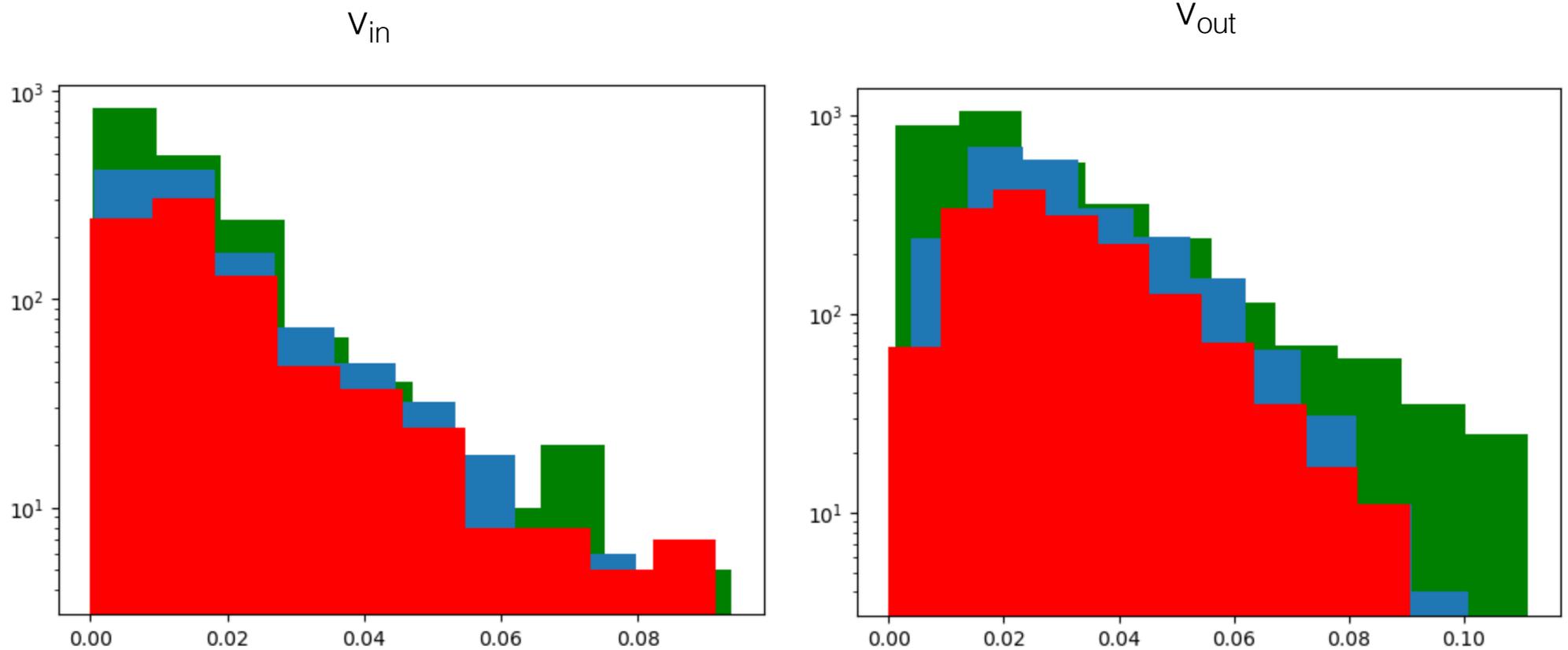


## Distribution of current sheet lengths



Note: normalisation of distributions is nonsense!

## Distribution of current sheet in/outflows



Note: normalisation of distributions is nonsense!

# An aside: non-thermal line-broadening in a flux braiding simulation

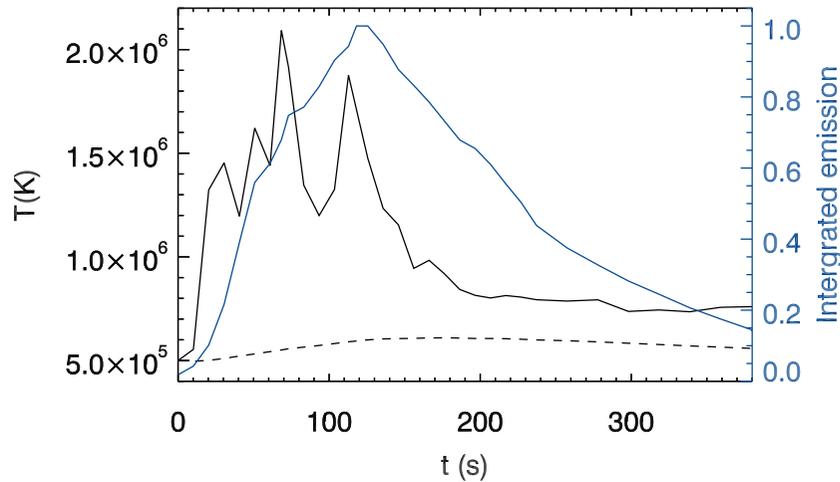
A&A 639, A21 (2020)  
<https://doi.org/10.1051/0004-6361/202037582>  
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**Astronomy  
&  
Astrophysics**

**Non-thermal line broadening due to braiding-induced turbulence in  
solar coronal loops**

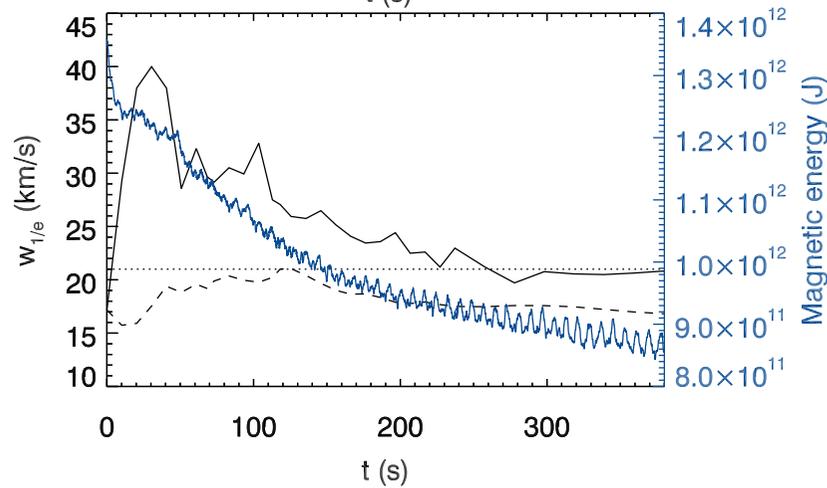
D. I. Pontin<sup>1,2</sup>, H. Peter<sup>3</sup>, and L. P. Chitta<sup>3</sup>

# Results: line broadening



Maximum temperature

Emission integrated over loop



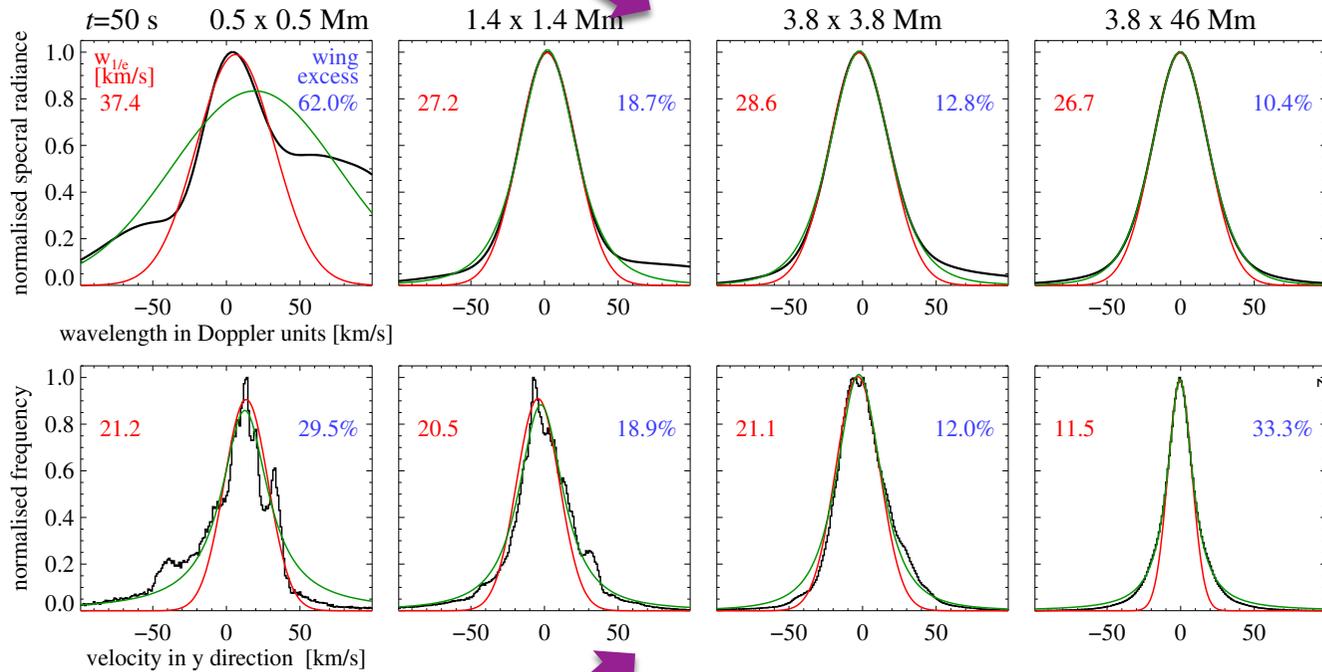
1/e-width of emission line

Magnetic energy

- When loop is brightest,  $w_{1/e} \sim 25\text{-}30$  km/s  $\Rightarrow w_{nth} \sim 14\text{-}21$  km/s

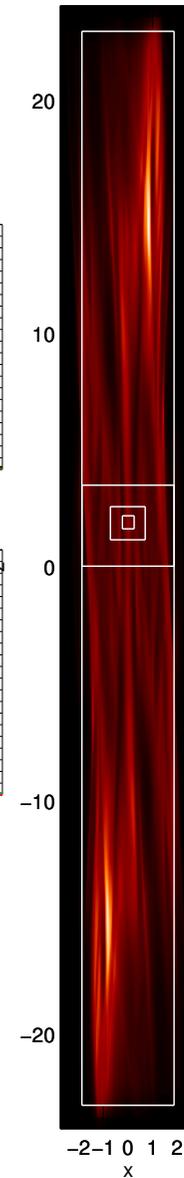
# Results: non-Gaussian spectra

Integrated emission (normalised)

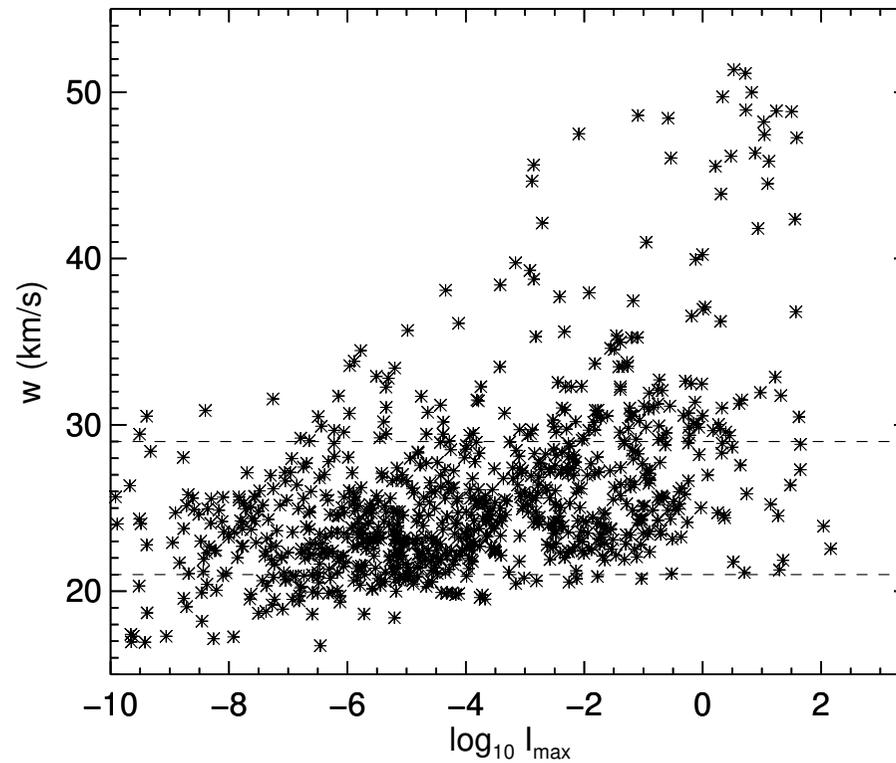


Line-of-sight velocity histograms

- Property (2): broadening  $\sim$  indep of resolution
- Property (4): enhanced power in wings compared to Gaussian ( $\sim 20\%$ ) – inherited from velocity distribution



## Results: intensity vs line width



- Property 3: correlation between line width and intensity
- Speculate: strongest flows assoc with recently reconnected flux bundles

## Conclusions (line broadening study)

- **Recall:** Coronal loop emission spectra (forming at around 1 to a few MK) exhibit the following properties:
  - 1) the non-thermal broadening is of the order of 15–30 km s<sup>-1</sup>
  - 2) the broadening is approximately independent of instrument resolution (down to ~0.35''),
  - 3) the correlation of line intensity and the non-thermal broadening
  - 4) the line profiles are non-Gaussian with enhanced power in the wings.
- These 4 properties reproduced self-consistently by loop heating through magnetic braiding.
- Wing enhancements are directly related to underlying velocity distribution, a signature of intermittency.