



The Thickness of Current Sheets and Implications for Coronal Heating

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Topological Dissipation (Parker 1972)

Continuous 3D magnetic fields (without separatrices or X-points) **MUST** develop infinitely thin current sheets when subjected to continuous motions at a line-tied boundary.

Infinitely thin current sheet = tangential discontinuity = singular current sheet

Straightforward explanation of coronal heating.

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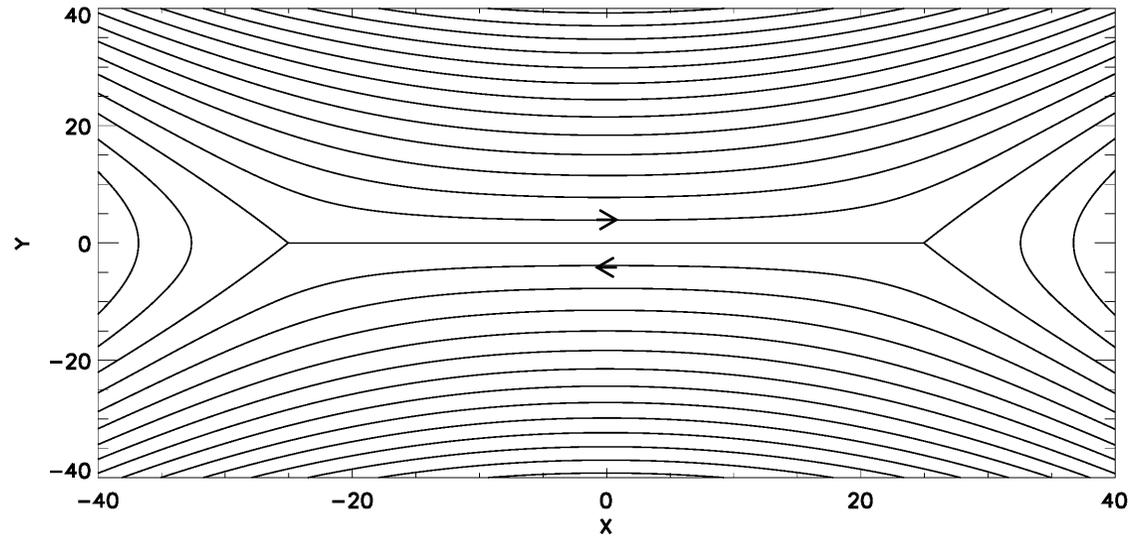
Straightforward explanation of coronal heating.

No – infinitely thin sheets are problematic! Magnetic stress and energy cannot build up.

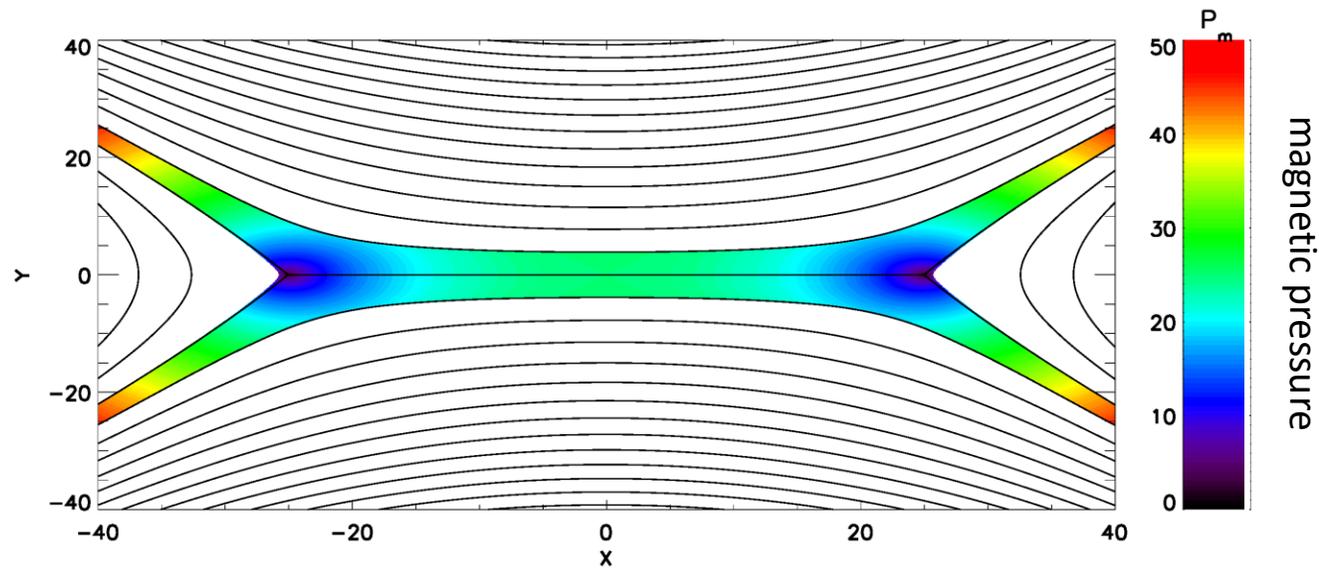
$$F = \frac{1}{4\pi} v_h B_h B_v$$

$$\frac{B_h}{B_v} \approx \tan(10^\circ)$$

Infinitely thin current sheet

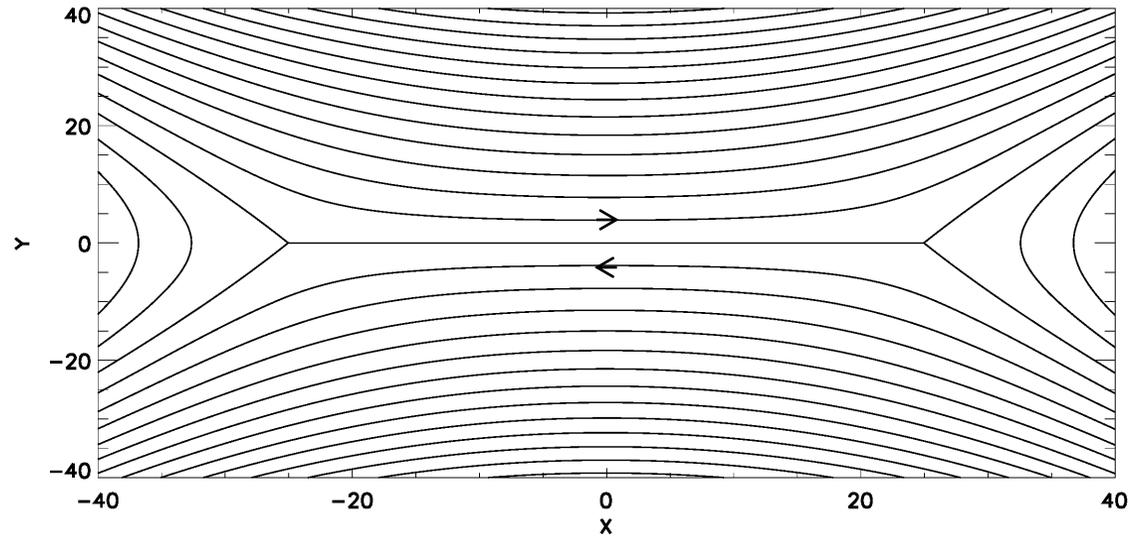


$$\frac{1}{c} (\mathbf{J} \times \mathbf{B}) = 0$$

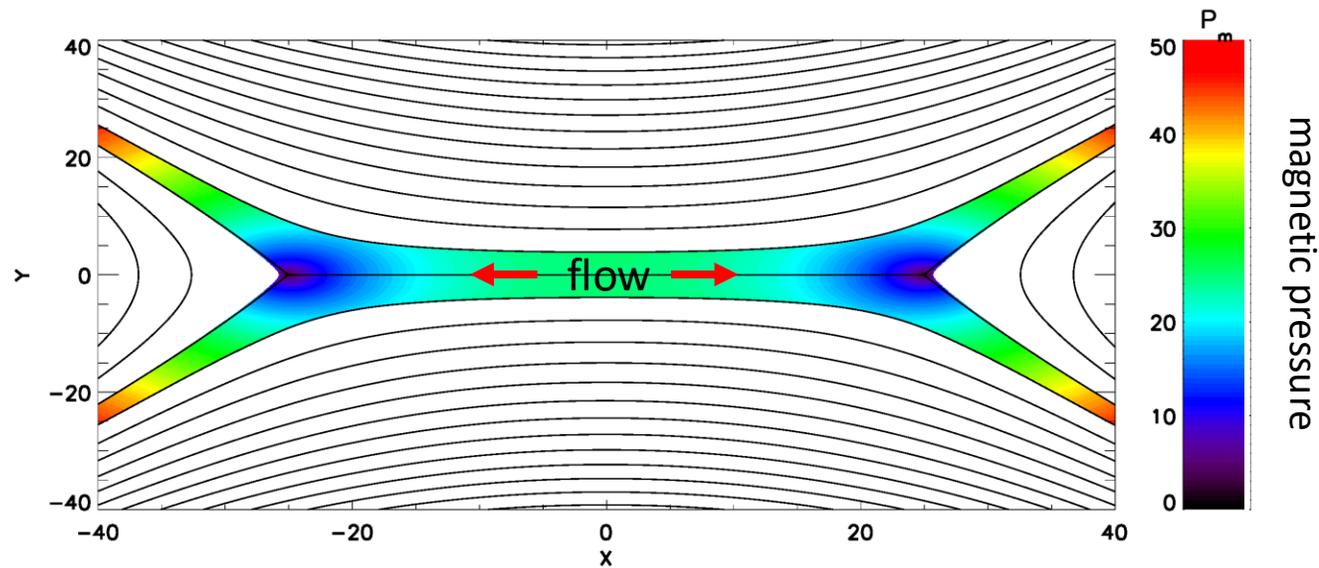


$$\nabla \left(\frac{B^2}{8\pi} \right) = \frac{1}{4\pi} (\mathbf{B} \cdot \nabla) \mathbf{B}$$

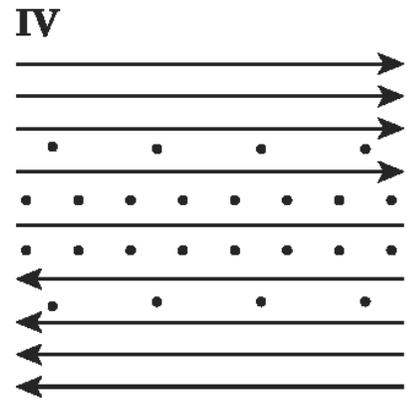
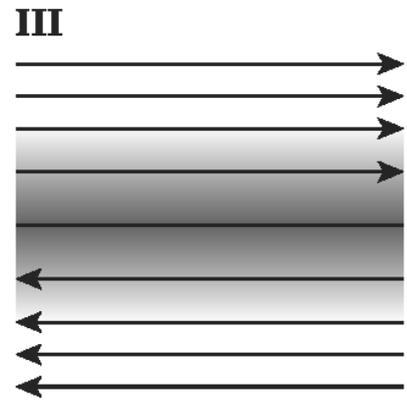
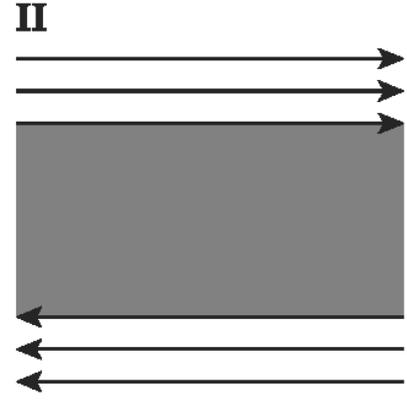
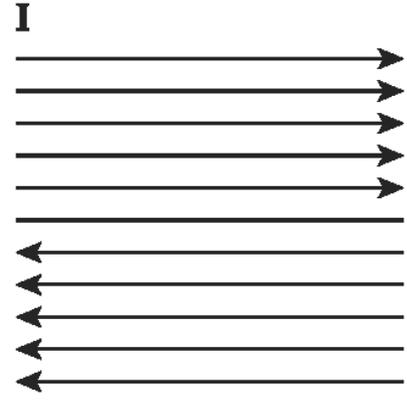
Infinitely thin current sheet

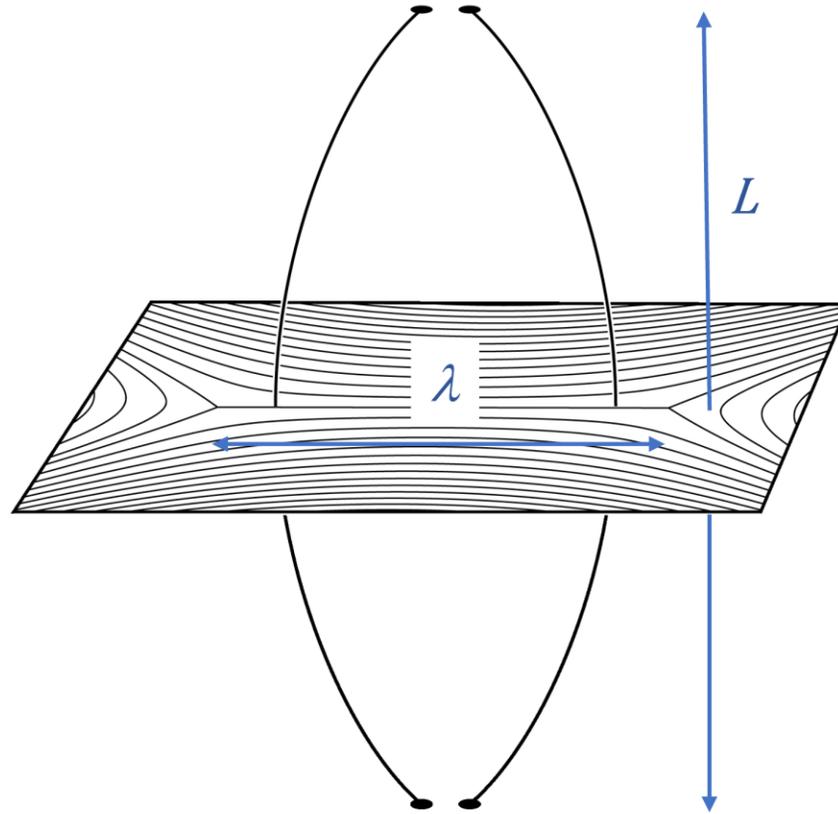


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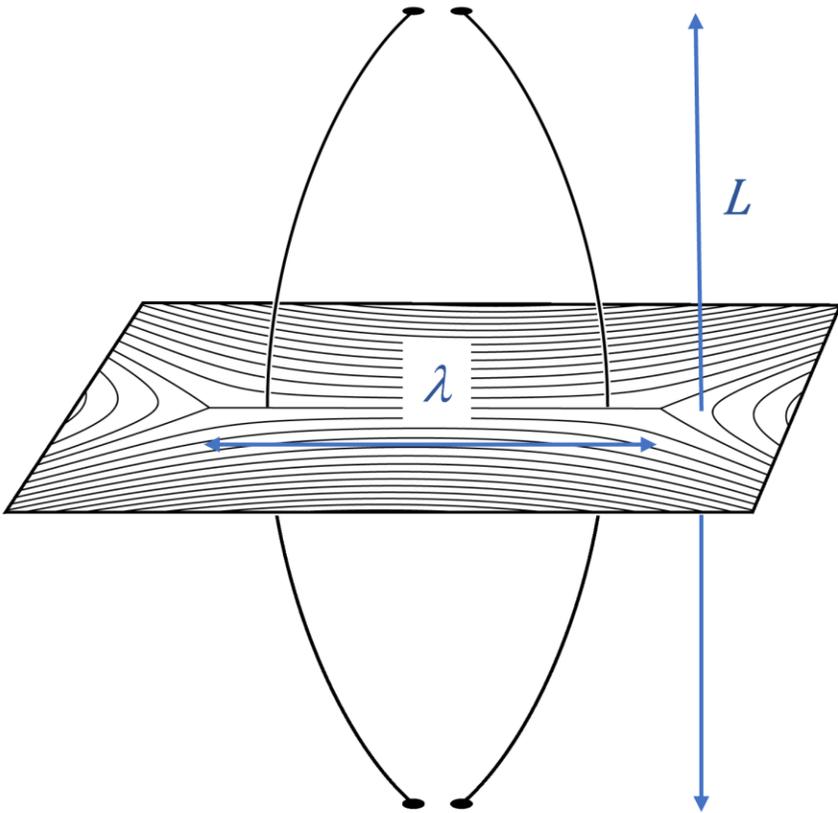
$$\nabla \left(\frac{B^2}{8\pi} \right) = \frac{1}{4\pi} (\mathbf{B} \cdot \nabla) \mathbf{B}$$





Tension in curved line-tied guide field.

Opposes pressure gradient force along sheet.



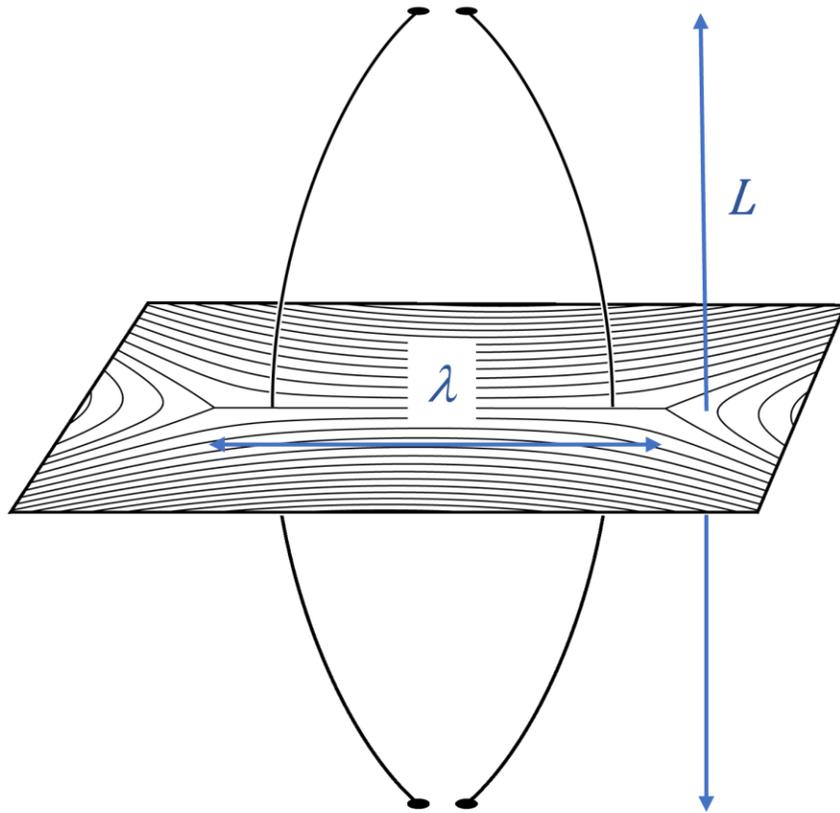
$$\frac{F_t}{F_p} \approx \frac{\lambda}{R_c} \left(\frac{B_{g0}}{B_{x0}} \right)^2$$

$$R_{c,min} \approx \frac{1}{\lambda} \left(\frac{L}{2} \right)^2$$

$$\frac{F_t}{F_p} \leq 4 \left(\frac{\lambda B_{g0}}{L B_{x0}} \right)^2$$

$$\frac{B_{x0}}{B_{g0}} > \left(\frac{B_{x0}}{B_{g0}} \right)_{crit} = \frac{2\lambda}{L}$$

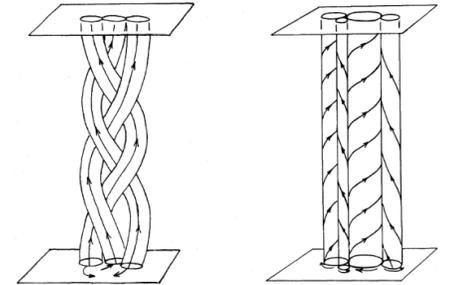
B_{g0}, B_{x0} = guide field and reconnecting field components outside of sheet



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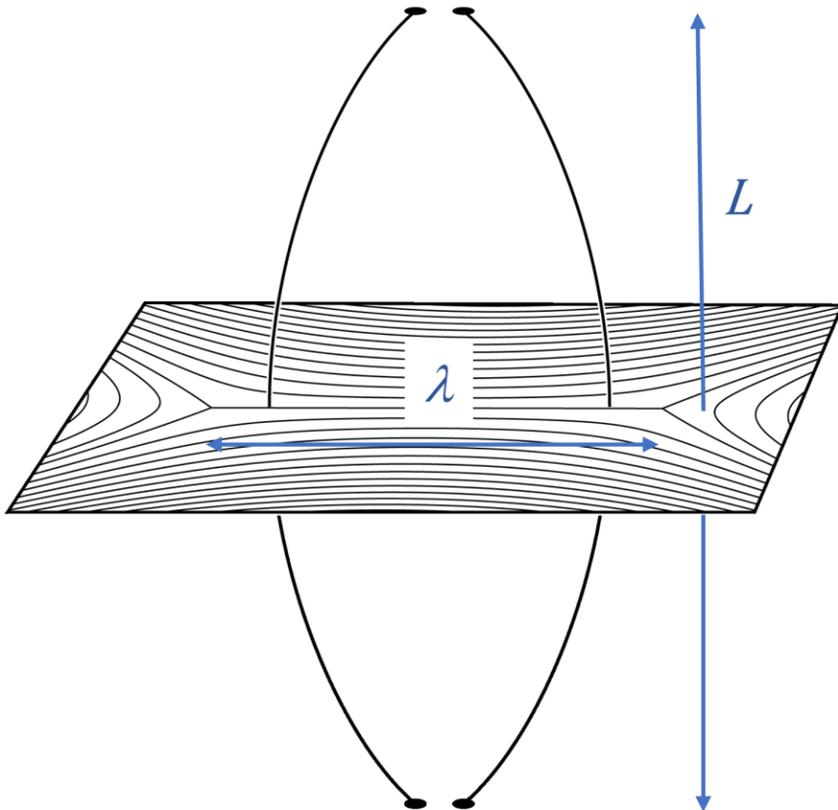
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$$\frac{B_{x0}}{B_{g0}} > \left(\frac{B_{x0}}{B_{g0}} \right)_{crit} = \frac{2\lambda}{L}$$

$$L = 50,000 \text{ km}$$

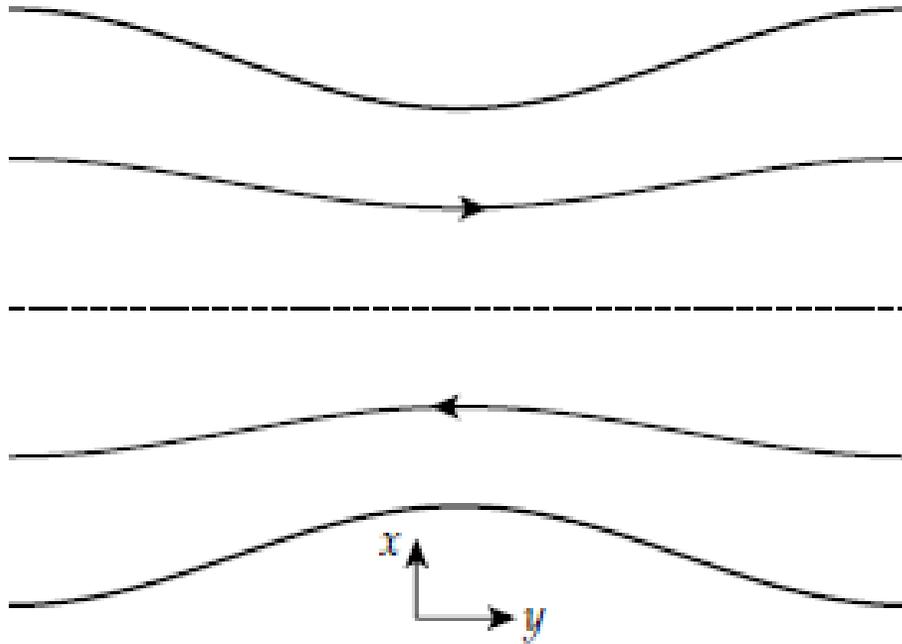
$$\lambda = 2,000 \text{ km}$$

$$\rightarrow (B_{x0}/B_{g0})_{crit} = 0.08 \quad (10^\circ \text{ shear angle})$$

Consistent with observed heating requirements in active regions.

Hahn-Kulsrud-Taylor (HKT) Problem

Zhou, Huang, Qin, & Bhattacharjee (2018)

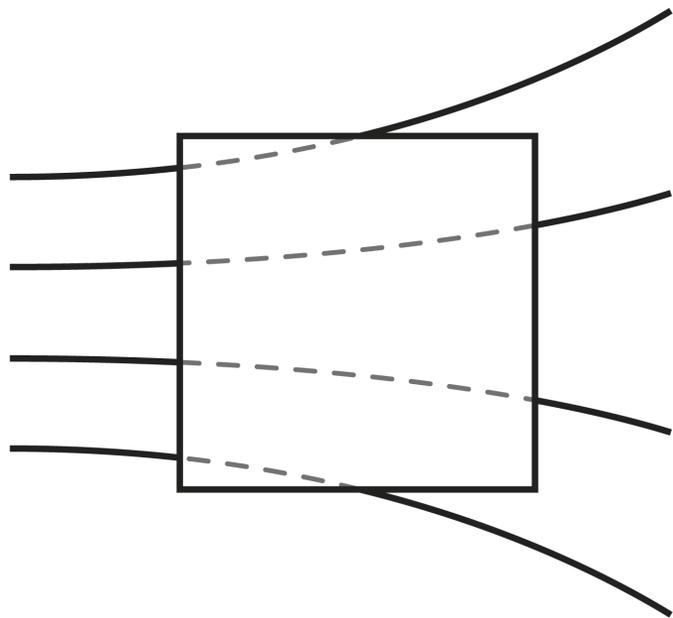
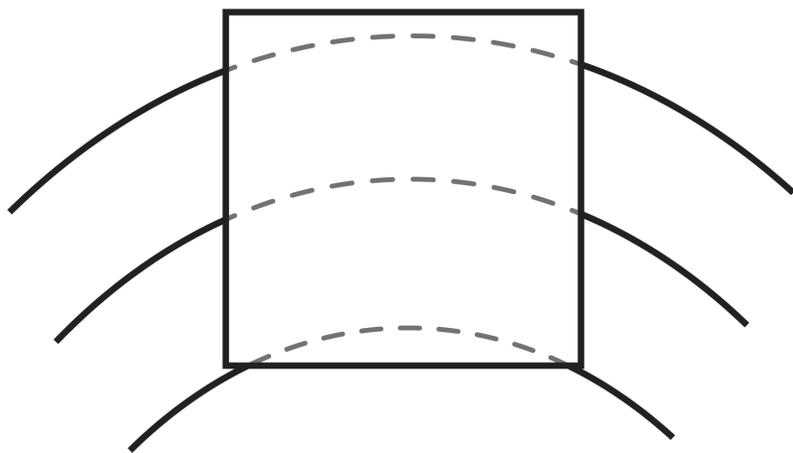


Huang et al. (2022)

$$L_{crit} = 29 \lambda$$

Our formula:

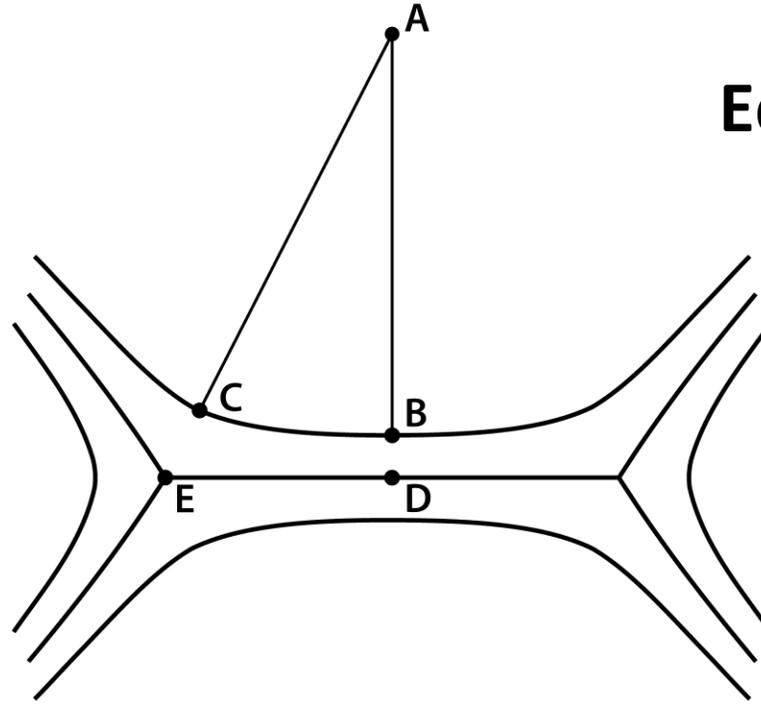
$$L > L_{crit} = 2\lambda \left(\frac{B_{g0}}{B_{x0}} \right) \\ = 11 \lambda$$



$$T_{ij} = \frac{1}{4\pi} \left(B_i B_j - \frac{1}{2} B^2 \delta_{ij} \right)$$



$B = 0$ in sheet

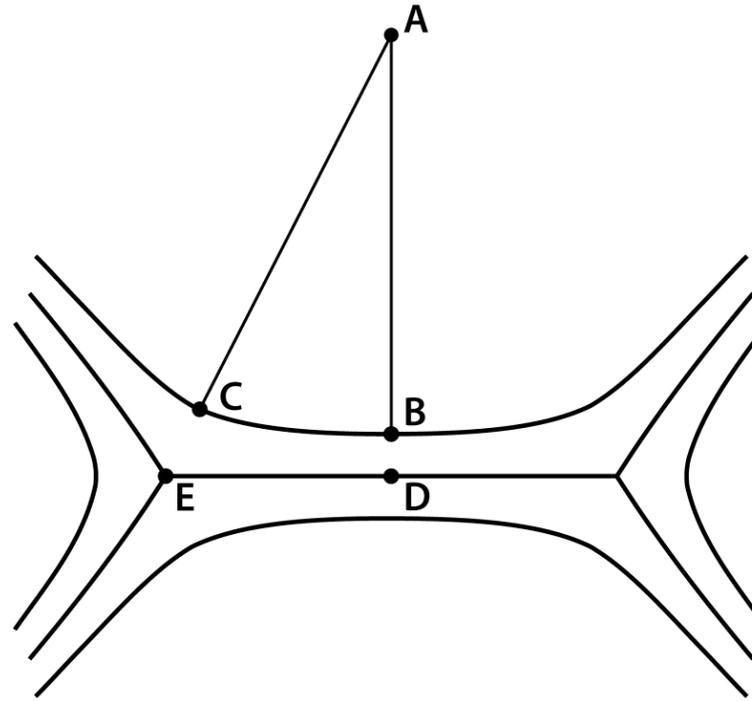


Equilibrium possible?

- Plasma pressure must be constant along field lines (and within entire sheet).
- Total pressure is constant across a magnetic boundary.
- Therefore, $B_C = B_B$.
- But $B_C < B_B$ to satisfy force balance outside the sheet. [tension force = $B^2/(4\pi R_c)$]
- No equilibrium exists at finite thickness (2D and 2.5D)



$B = 0$ at center of sheet
 P_{max} at center of sheet



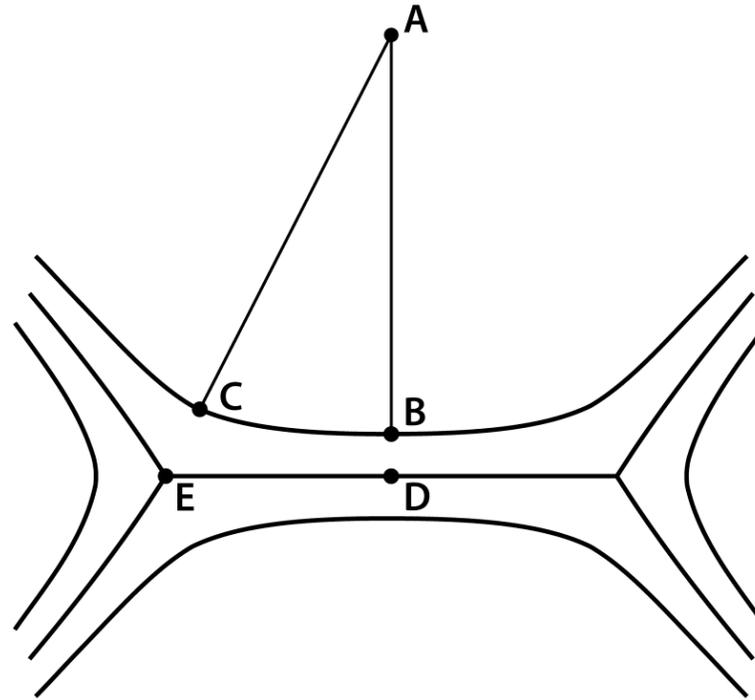
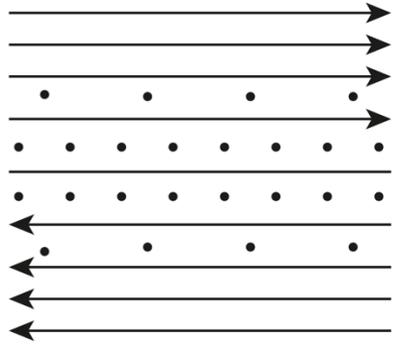
$$P_E < B_C^2/8\pi$$

$$B_C^2/8\pi < B_B^2/8\pi$$

$$P_D \approx B_B^2/8\pi$$

$$P_E < P_D$$

But $P_E = P_D$ in equilibrium. Horizontal outflow from center of sheet.



Replace plasma with a straight guide field.

$P =$ guide field pressure, $B_{guide}^2/8\pi$

Field rotates by 180° across sheet.

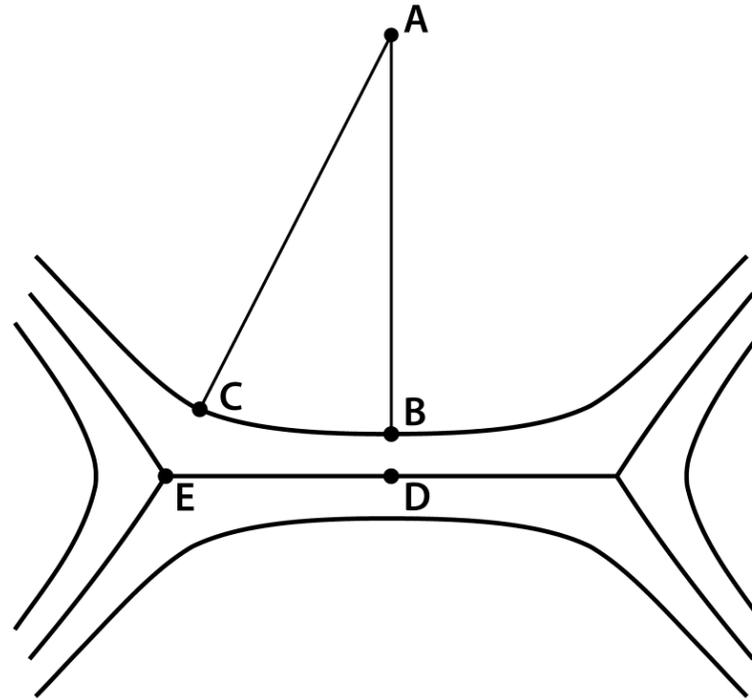
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Replace plasma with a straight guide field.

$P =$ guide field pressure, $B_{guide}^2/8\pi$

Can add a uniform guide field, B_{g0} .

Field rotates by $< 180^\circ$.

$$P_E < B_C^2/8\pi$$

$$B_C^2/8\pi < B_B^2/8\pi$$

$$P_D \approx B_B^2/8\pi$$

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